

Mesoscale Ocean Large Eddy Simulations (MOLES)

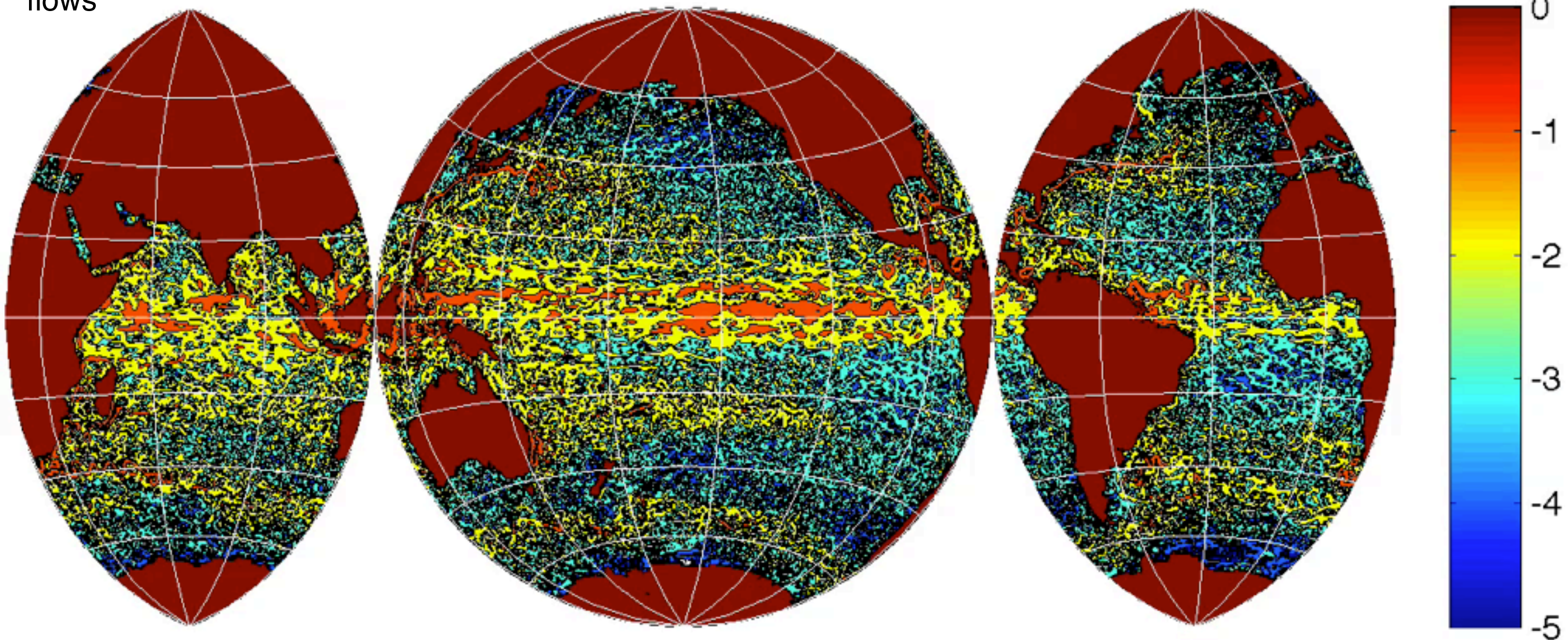
Baylor Fox-Kemper (Brown DEEP Sciences)

with Brodie Pearson (Brown), Frank O. Bryan (NCAR), and S. Bachman (DAMTP)

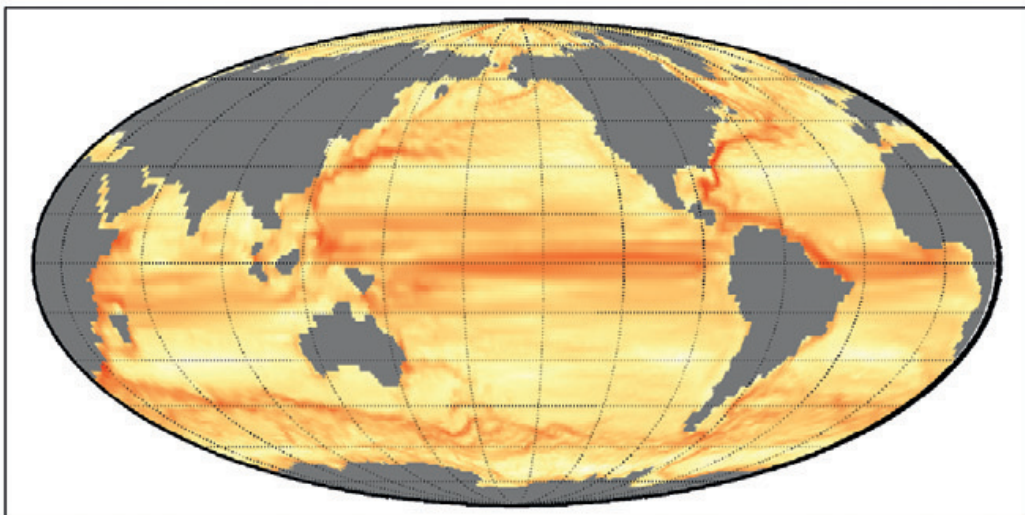
UKMO Hadley Centre 4/15/16, Sponsor: NSF 1350795

Satellite altimetry
view of mesoscale
flows

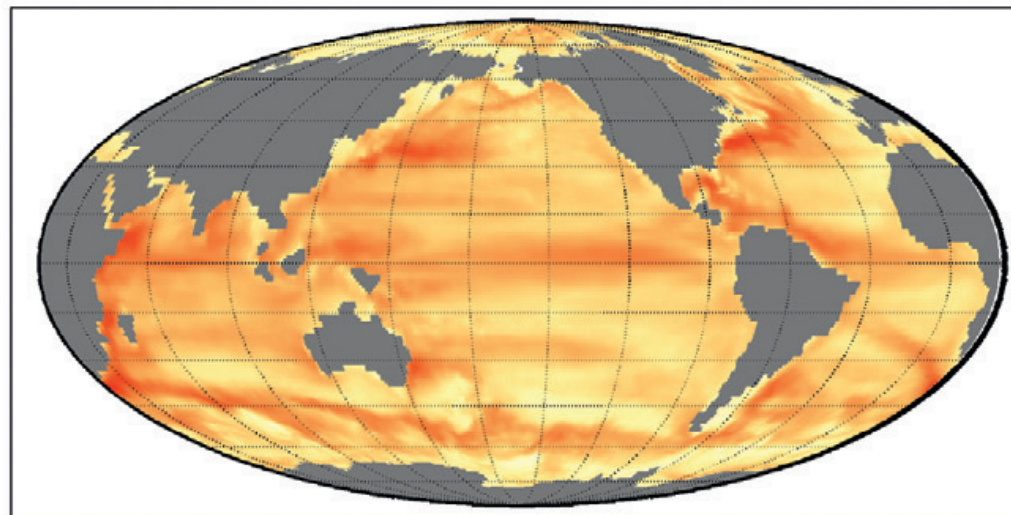
AVISO: $\log_{10}(0.5(u^2+v^2))$ on 19940101



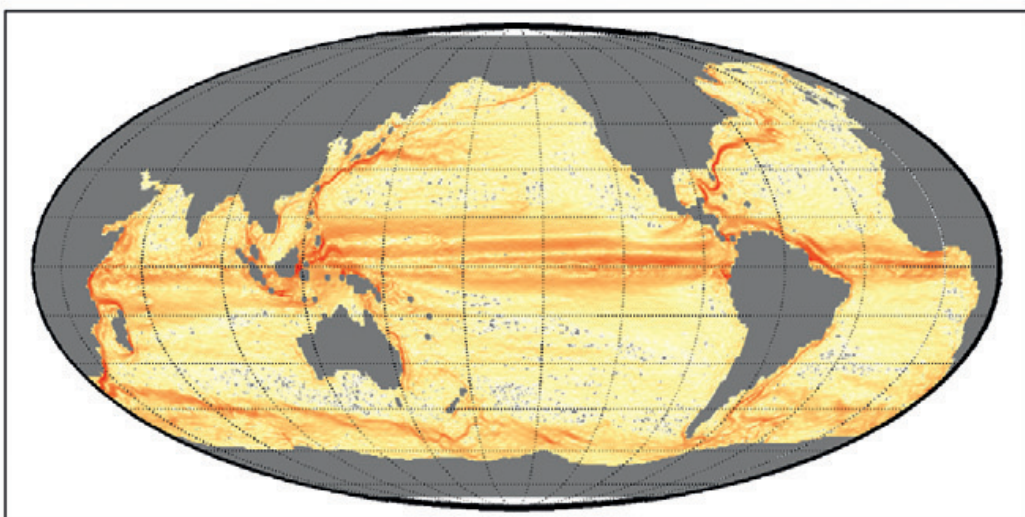
(a) $\log_{10}(\text{Mean kinetic energy from model (cm}^2/\text{s}^2))$



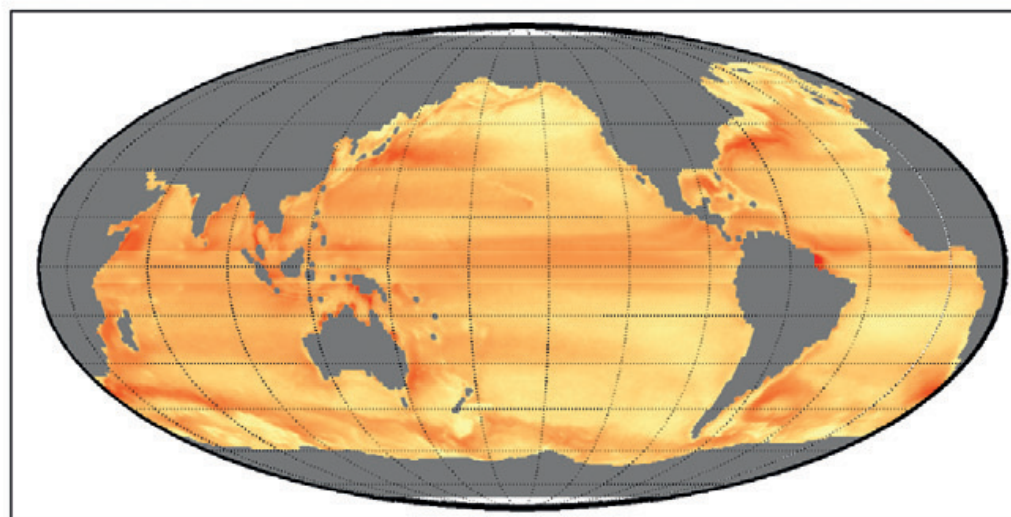
(b) $\log_{10}(\text{Eddy kinetic energy from model (cm}^2/\text{s}^2))$



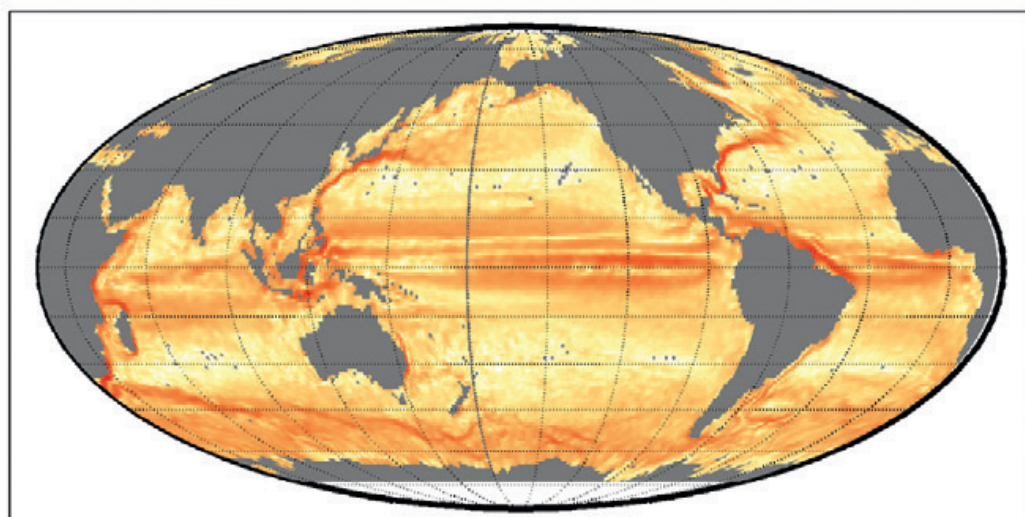
(c) $\log_{10}(\text{Mean kinetic energy from AVISO 1993-2010 (cm}^2/\text{s}^2))$



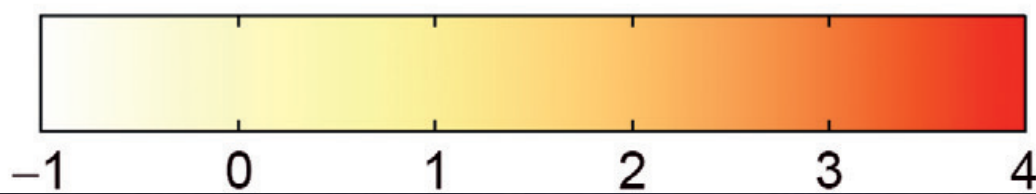
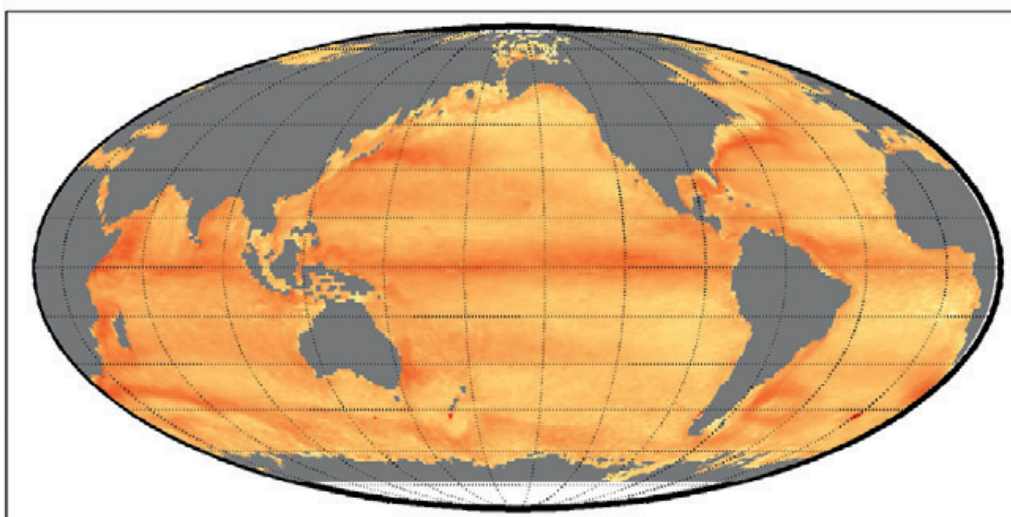
(d) $\log_{10}(\text{Eddy kinetic energy from AVISO 1993-2010 (cm}^2/\text{s}^2))$



(e) $\log_{10}(\text{Mean kinetic energy from drifters (cm}^2/\text{s}^2))$



(f) $\log_{10}(\text{Eddy kinetic energy from drifters (cm}^2/\text{s}^2))$



On cursory
analysis,
0.1 degree
models do
well vs.
Satellites
and
Drifters

B. Fox-Kemper, R. Lumpkin, and F. O. Bryan. Lateral transport in the ocean interior. In G. Siedler, S. M. Griffies, J. Gould, and J. A. Church, editors, *Ocean Circulation and Climate: A 21st century perspective*, volume 103 of *International Geophysics Series*, chapter 8, pages 185-209. Academic Press (Elsevier Online), 2013.

But, we know choices
are made in models...

- Subgrid parameterizations
 - "Do no harm" vs. "approximate unresolved scales"
- Resolution
 - "Permitting", "Resolving", Etc.



Viscosity Scheme: BFK and D. Menemenlis. Can large eddy simulation techniques improve mesoscale-rich ocean models? In M. Hecht and H. Hasumi, editors, *Ocean Modeling in an Eddying Regime*, volume 177, pages 319-338. AGU Geophysical Monograph Series, 2008.

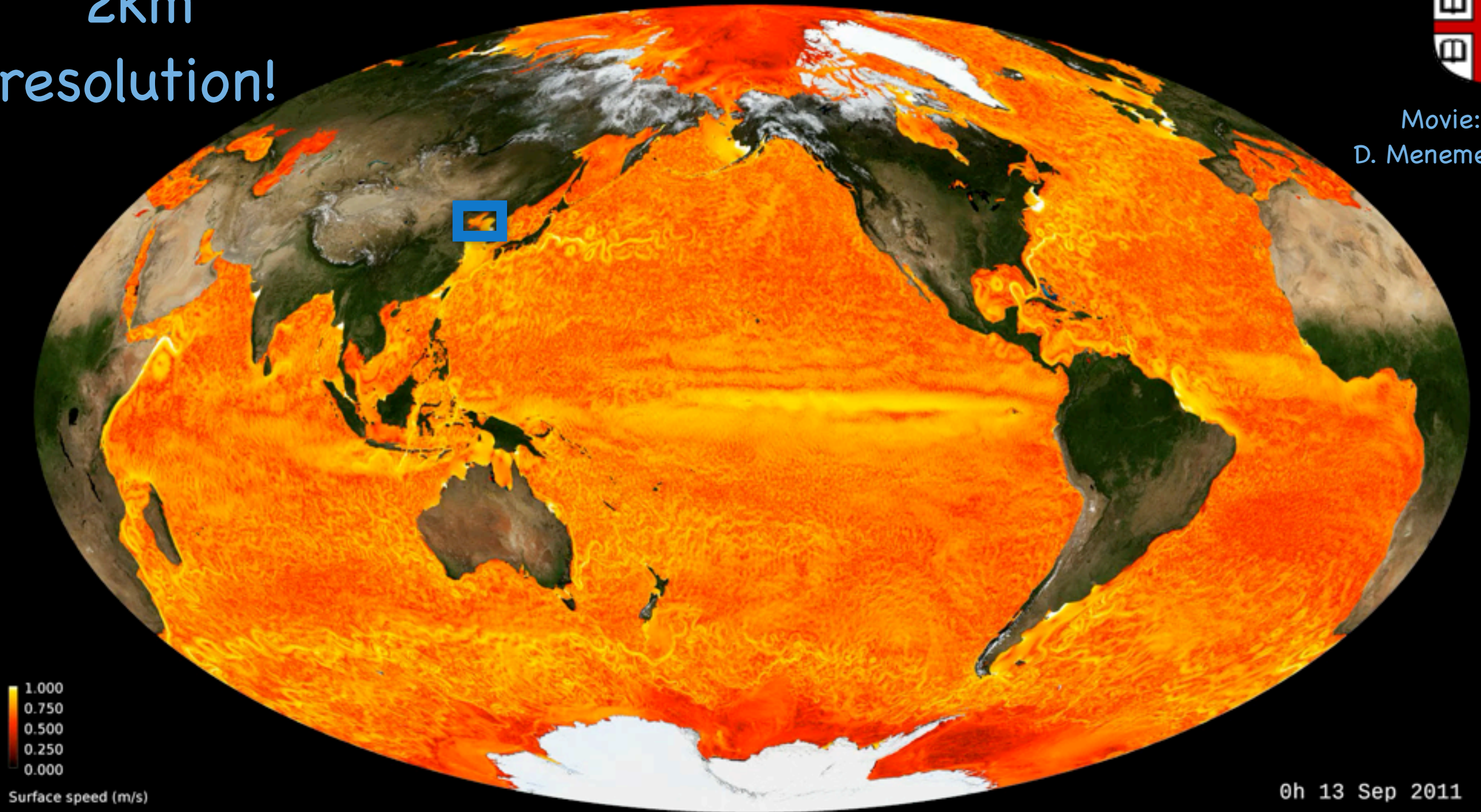
18km resolution



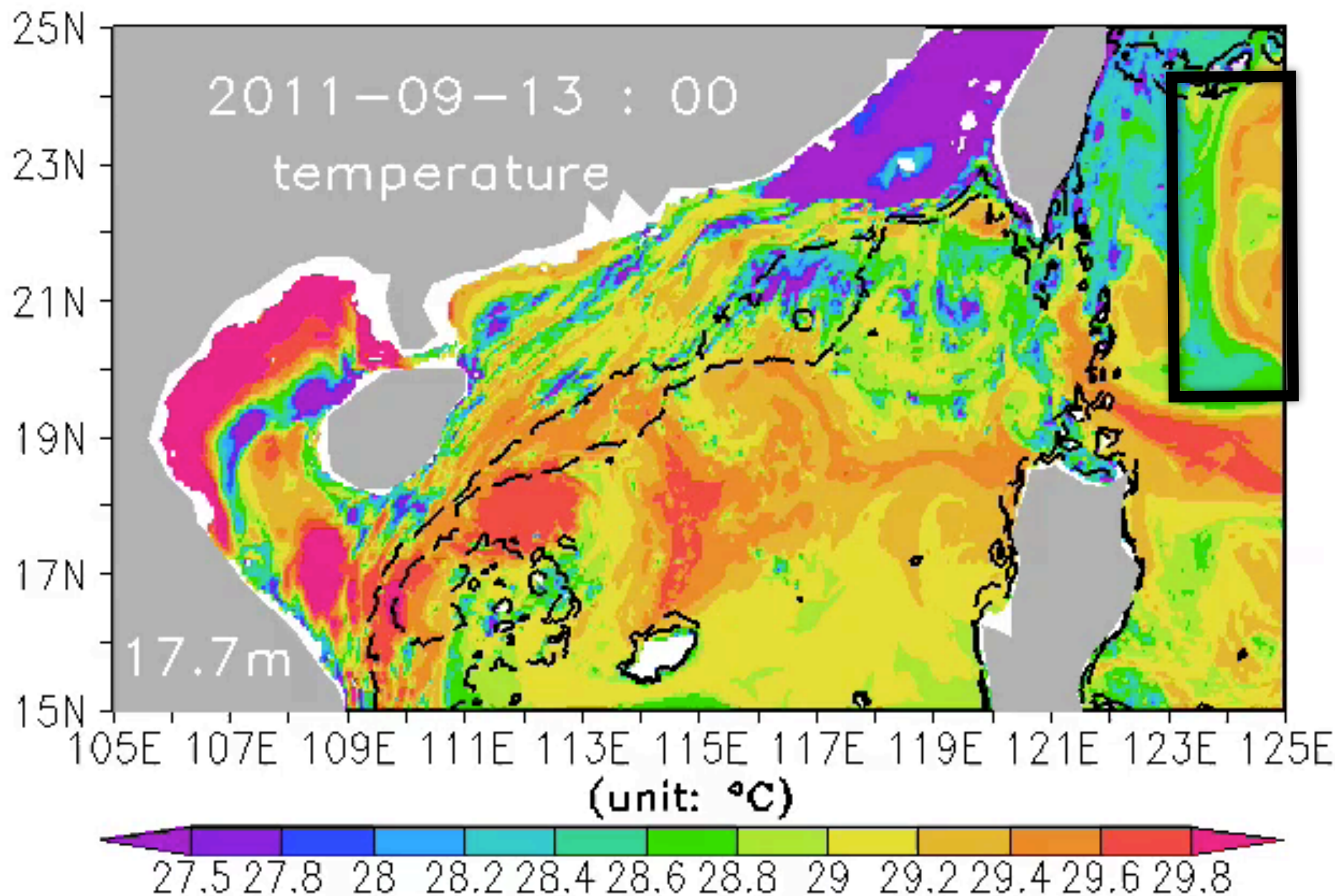
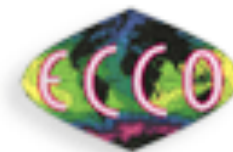
2km
resolution!



Movie:
D. Menemenlis



B. Fox-Kemper, S. Bachman, B. Pearson, and S. Reckinger. Principles and advances in subgrid modeling for eddy-rich simulations. CLIVAR Exchanges, 19(2):42-46, July 2014.



Movie:
Z. Jing



Brown Visitor
from
S. China Sea
Institute of Ocean.

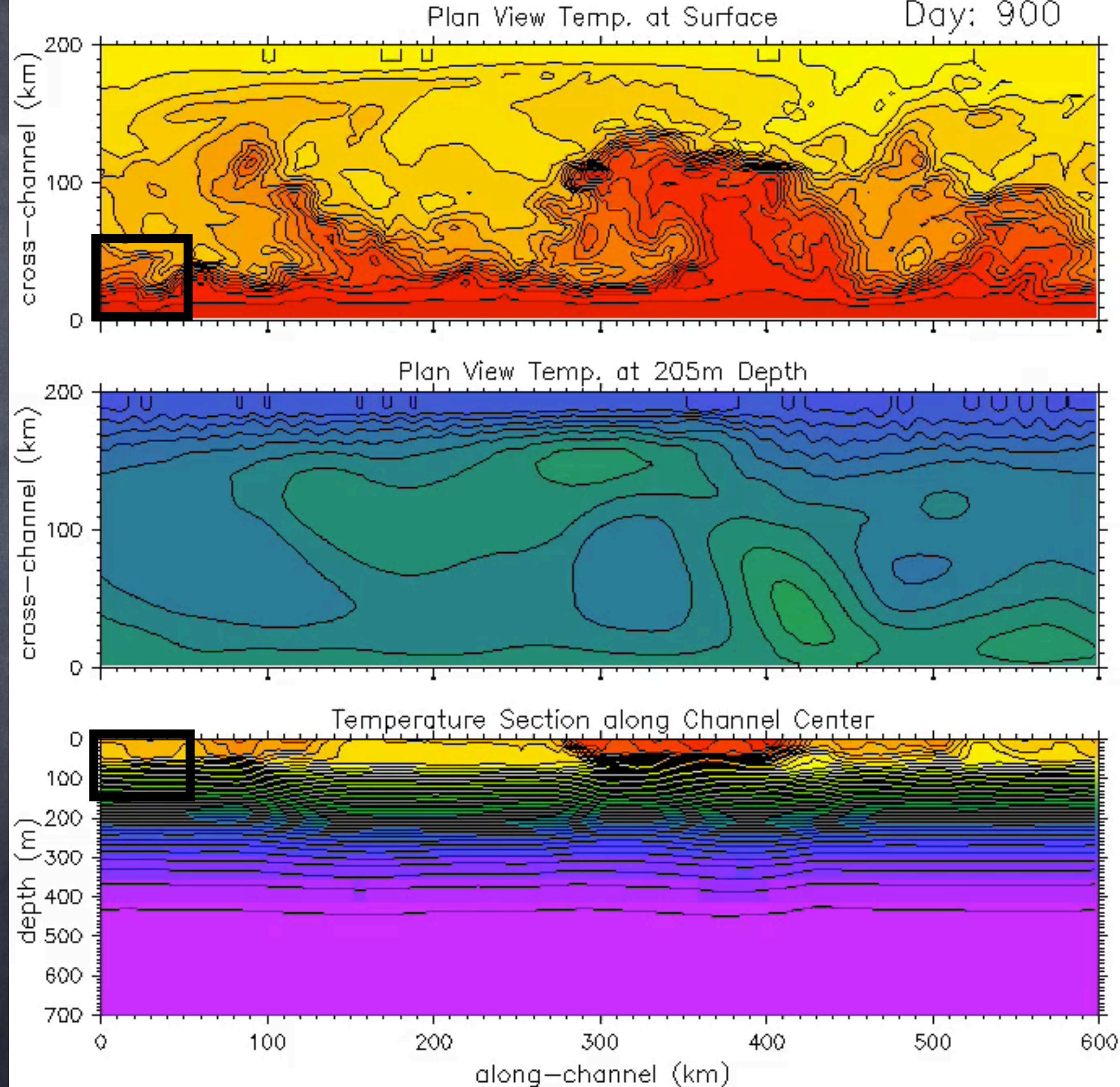
Local Analysis: Z. Jing, Y. Qi, BFK, Y. Du, and S. Lian. Seasonal thermal fronts and their associations with monsoon forcing on the continental shelf of northern South China Sea: Satellite measurements and three repeated field surveys in winter, spring and summer. *Journal of Geophysical Research-Oceans*, August 2015. In press.

200km x 600km
x 700m
domain

1000 Day
Simulation

If we lose
the globe,
much higher
resolution!

G. Boccaletti, R. Ferrari, and BFK.
Mixed layer instabilities and
restratification. *Journal of Physical
Oceanography*, 37(9):2228-2250,
2007.



20km x 20km x 150m
domain

10 Day Simulation

4m x 4m x 1m
Resolution

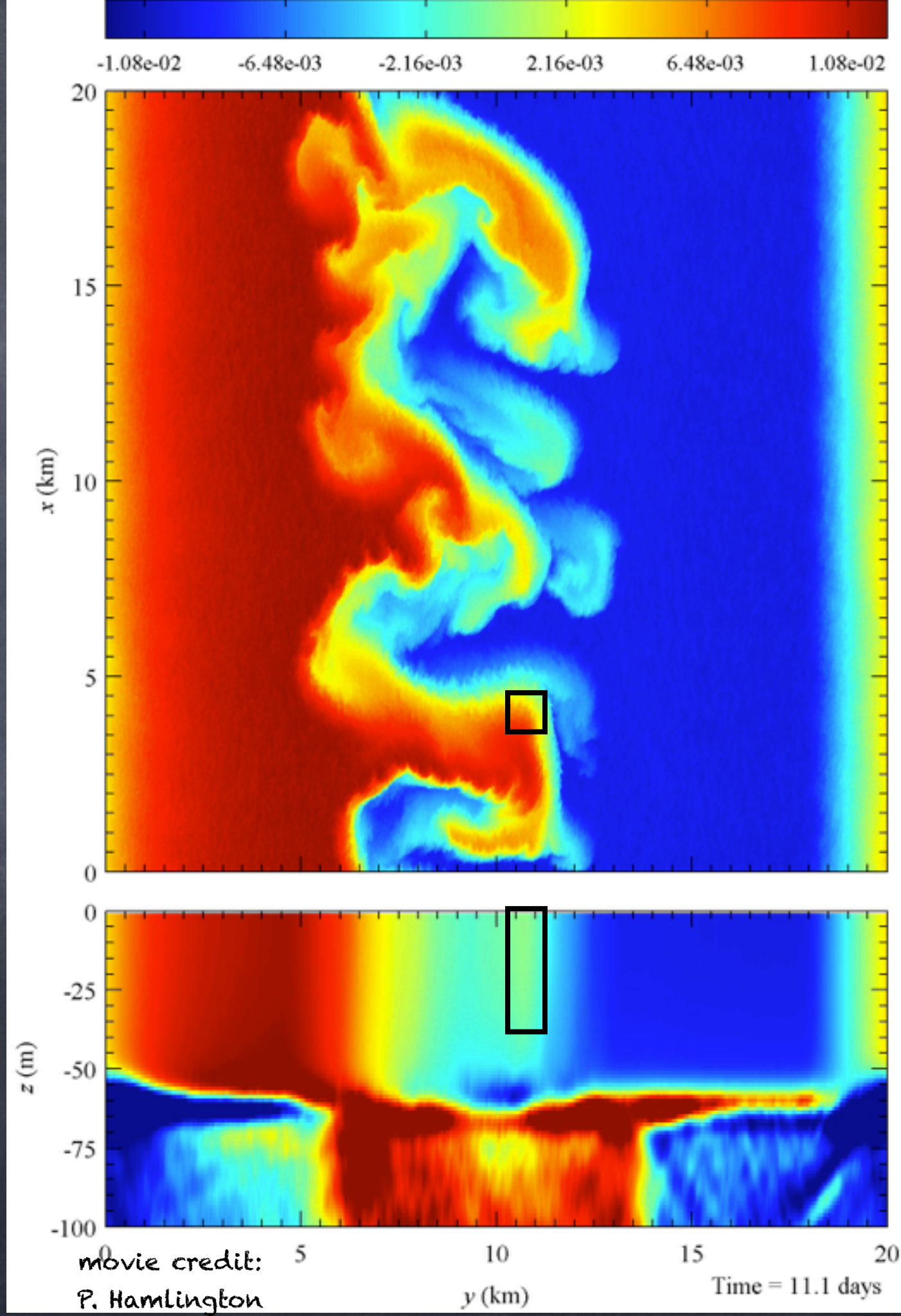


CU, now CU



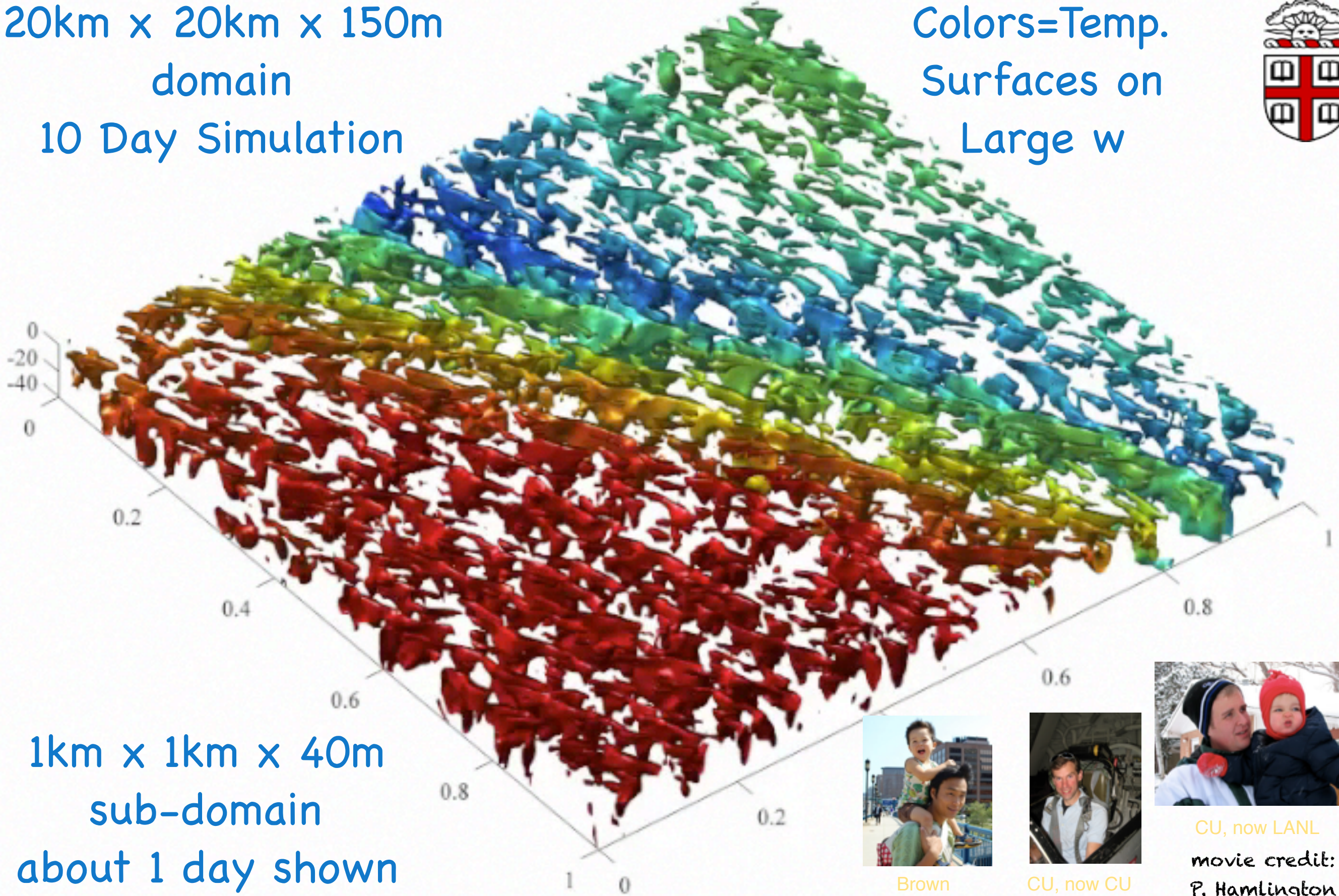
CU, now LANL

P. E. Hamlington, L. P. Van Roekel, BFK, K. Julien, and G. P. Chini. Langmuir-submesoscale interactions: Descriptive analysis of multiscale frontal spin-down simulations. *Journal of Physical Oceanography*, 44(9):2249-2272, September 2014.



20km x 20km x 150m
domain
10 Day Simulation

Colors=Temp.
Surfaces on
Large w



1km x 1km x 40m
sub-domain
about 1 day shown



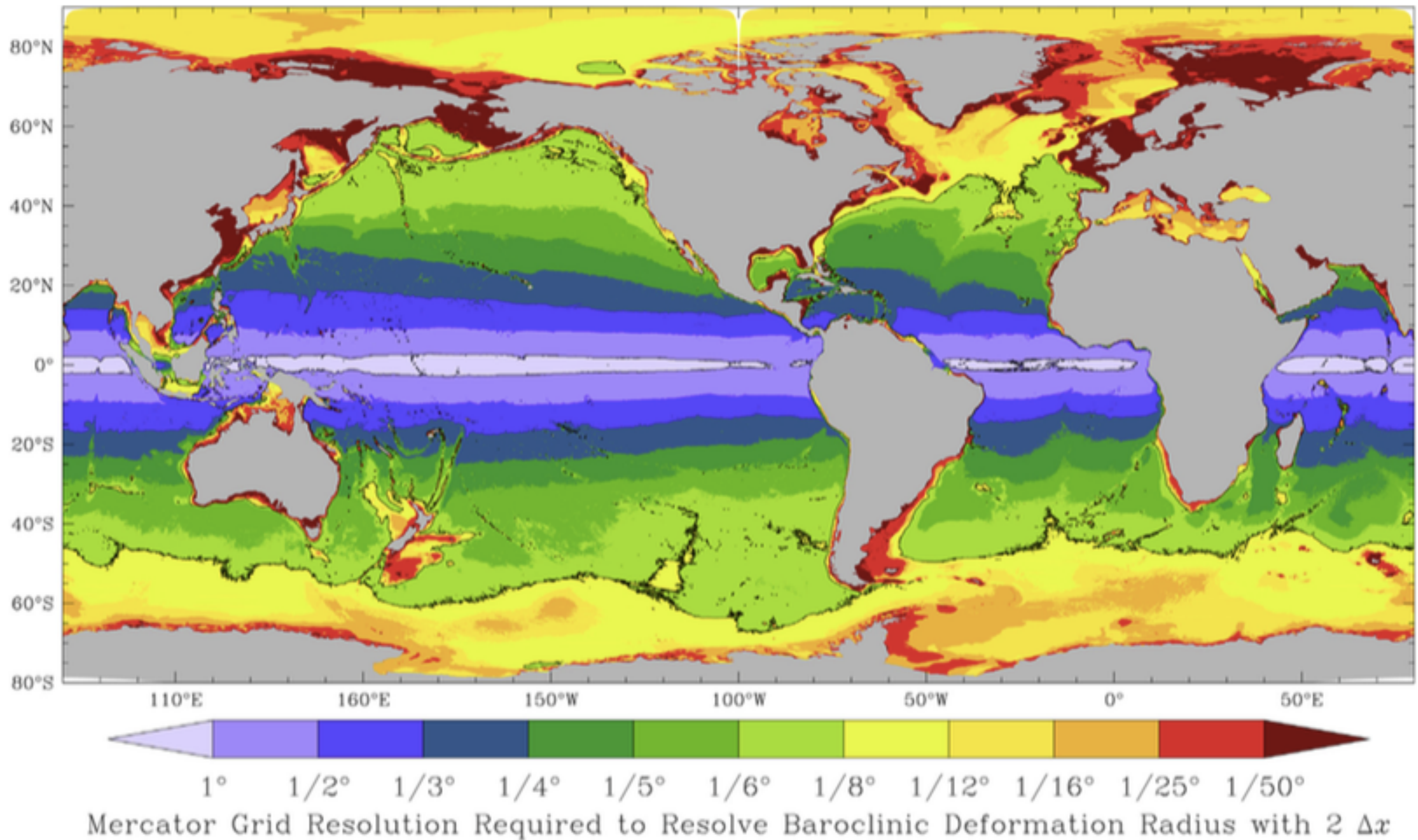
Brown



CU, now CU



CU, now LANL
movie credit:
P. Hamlington



- In most places, 0.1 degree resolves the largest deformation radius, plus a bit: Mesoscale Ocean Large Eddy Simulation



Key Concept for

Mesoscale Ocean Large Eddy Simulations (MOLES): Gridscale* Nondimensional Parameters

Gridscale Reynolds¹:

$$Re^* = \frac{U^* \Delta x}{\nu^*}$$

Gridscale Péclet¹:

$$Pe^* = \frac{U^* \Delta x}{\kappa^*}$$

Gridscale Rossby:

$$Ro^* = \frac{U^*}{f \Delta x}$$

Gridscale Richardson:

$$Ri^* = \frac{\Delta b^{*2} \Delta z}{\Delta U^{*2}}$$

Gridscale Burger:

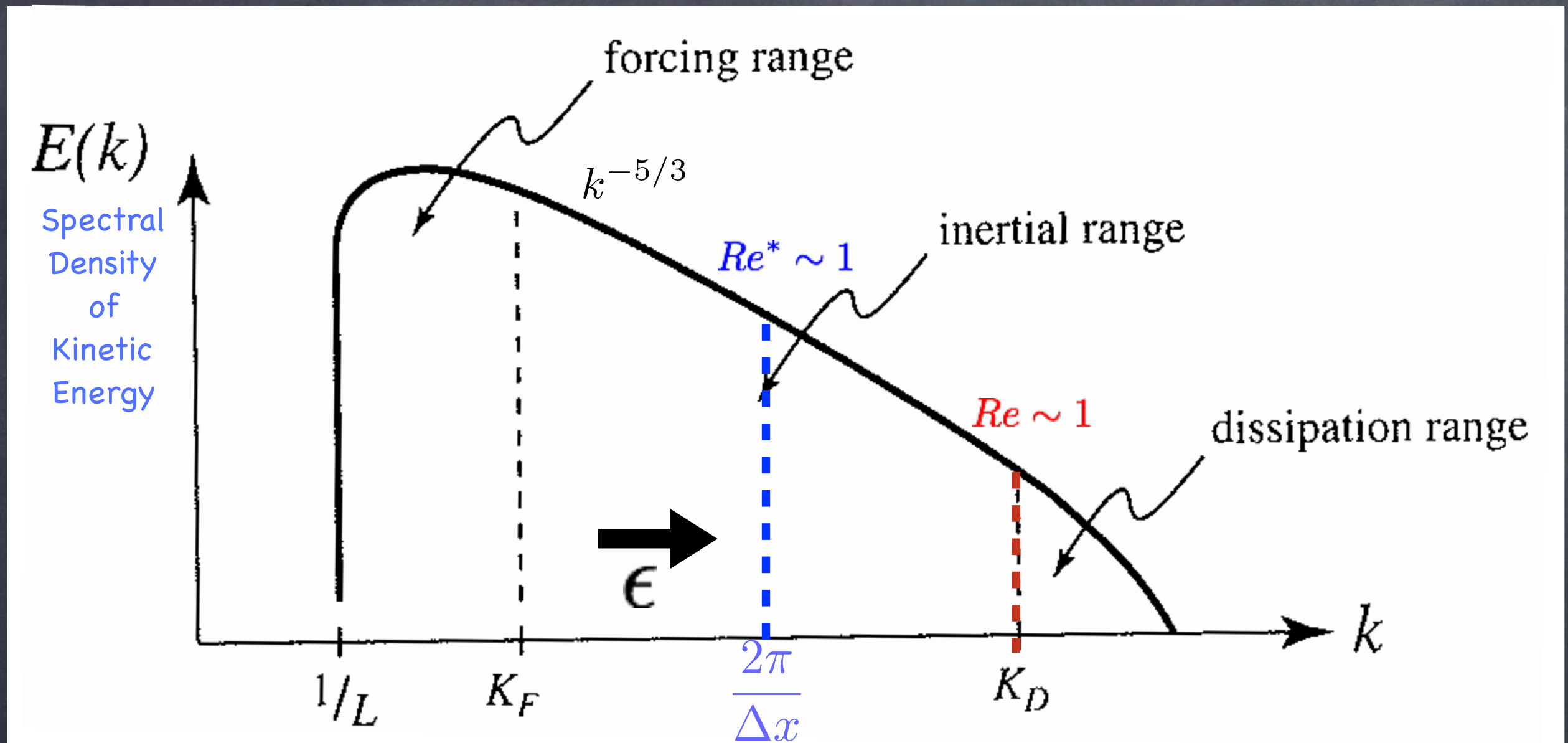
$$Bu^* = \frac{N^{*2} \Delta z^2}{f^2 \Delta x^2} \sim Ro^{*2} Ri^*$$

Asterisks denote *resolved* quantities, rather than true values

¹ Gridscale Reynolds and Péclet numbers MUST be $O(1)$ for numerical stability

3D Turbulence Cascade

Kolmogorov, 1941



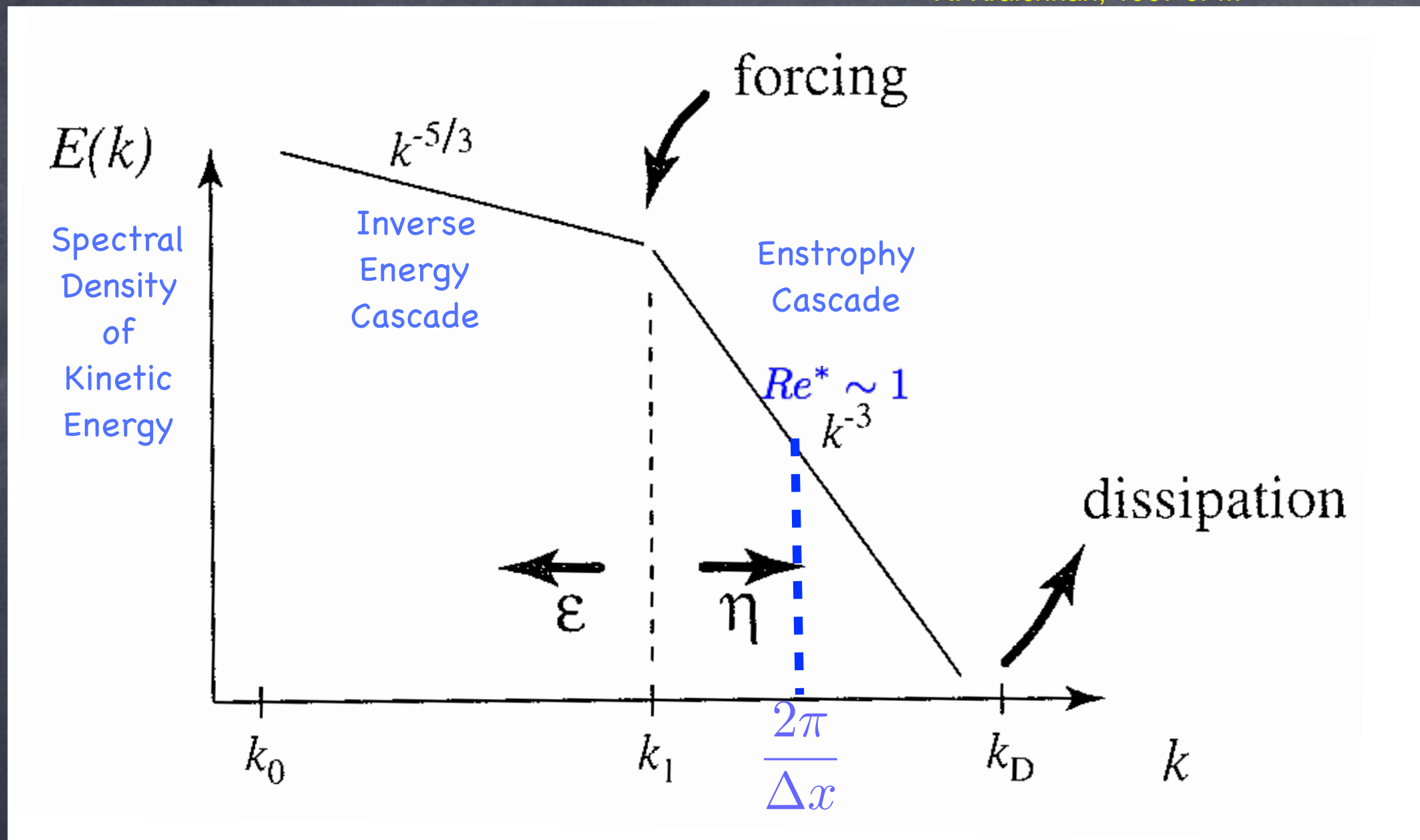
Smagorinsky (1963) Scale & Flow Aware Viscosity Scaling,
 So the Energy Cascade is Preserved, and $Re^* = \frac{U^* \Delta x}{\nu^*} = O(1)$

$$\mathbf{v}_{*h} = \left(\frac{\Upsilon_h \Delta x}{\pi} \right)^2 \sqrt{\left(\frac{\partial u_*}{\partial x} - \frac{\partial v_*}{\partial y} \right)^2 + \left(\frac{\partial u_*}{\partial y} + \frac{\partial v_*}{\partial x} \right)^2}$$

$$\Upsilon_h \approx 1$$

2D Turbulence Differs

R. Kraichnan, 1967 JFM



Leith (1996) Devises Viscosity Scaling,
 So that the Enstrophy Cascade is preserved, and $Re^* = \frac{U^* \Delta x}{\nu^*} = O(1)$

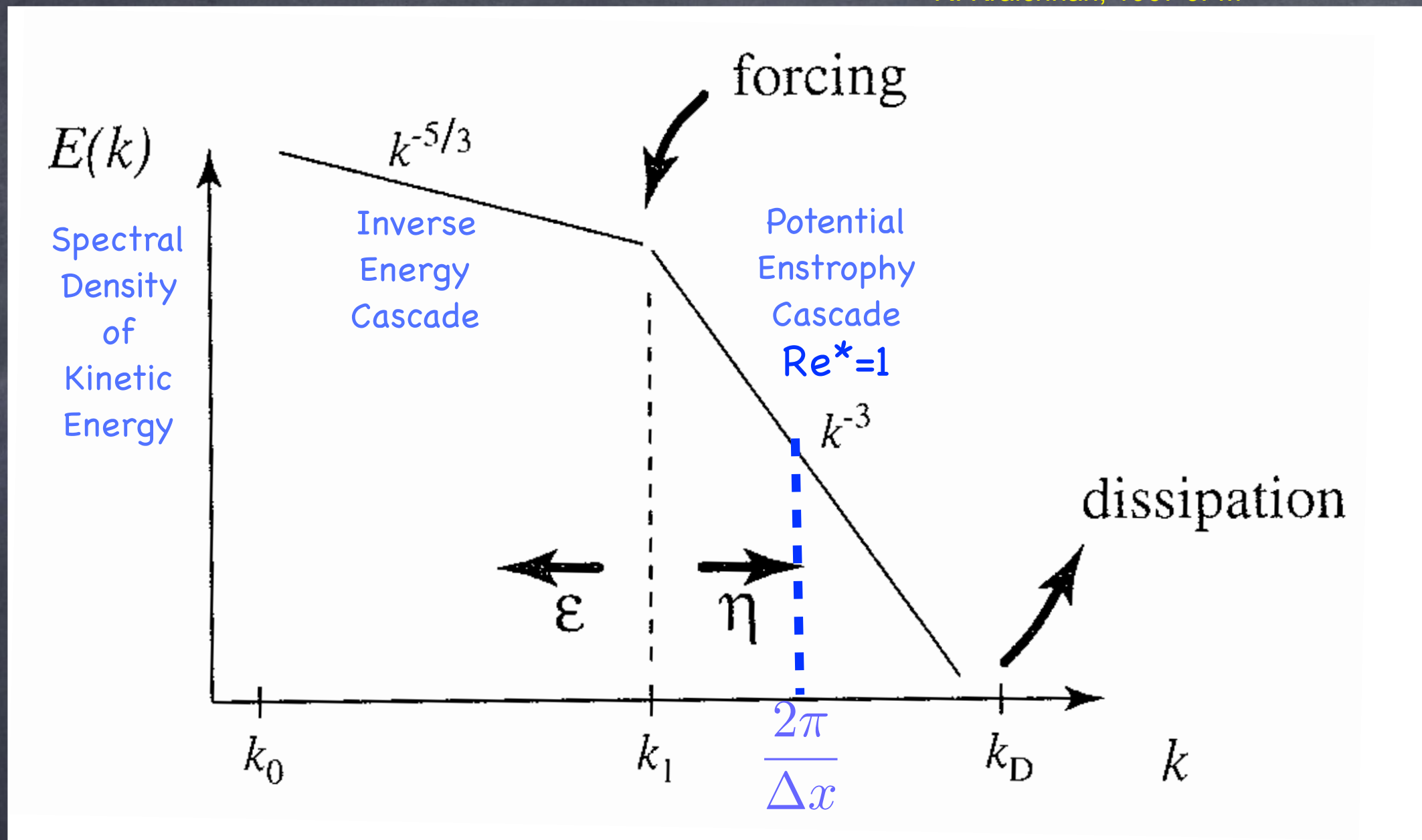
2D Leith:

$$\nu_{2d}^* = \left(\frac{\Lambda_{2d} \Delta x}{\pi} \right)^3 |\nabla_h q_{2d}^*|$$

$$q_{2d}^* = \frac{\partial u^*}{\partial y} - \frac{\partial v^*}{\partial x}$$

2D Turbulence: (enstrophy=vorticity²)

R. Kraichnan, 1967 JFM



F-K & Menemenlis '08: Revise Leith Viscosity Scaling, So that diverging, vorticity-free, modes are also damped

$$\mathbf{v}_* = \left(\frac{\Delta x}{\pi} \right)^3 \sqrt{\Lambda_d^6 |\nabla_h q_{2d}^*|^2 + \Lambda_d^6 |\nabla_h (\nabla_h \cdot \mathbf{u}_*)|^2}$$

B. Fox-Kemper and D. Menemenlis. Can large eddy simulation techniques improve mesoscale-rich ocean models? In M. Hecht and H. Hasumi, editors, Ocean Modeling in an Eddy Regime, volume 177, pages 319-338. AGU Geophysical Monograph Series, 2008.



Viscosity Scheme: BFK and D. Menemenlis. Can large eddy simulation techniques improve mesoscale-rich ocean models? In M. Hecht and H. Hasumi, editors, *Ocean Modeling in an Eddying Regime*, volume 177, pages 319-338. AGU Geophysical Monograph Series, 2008.

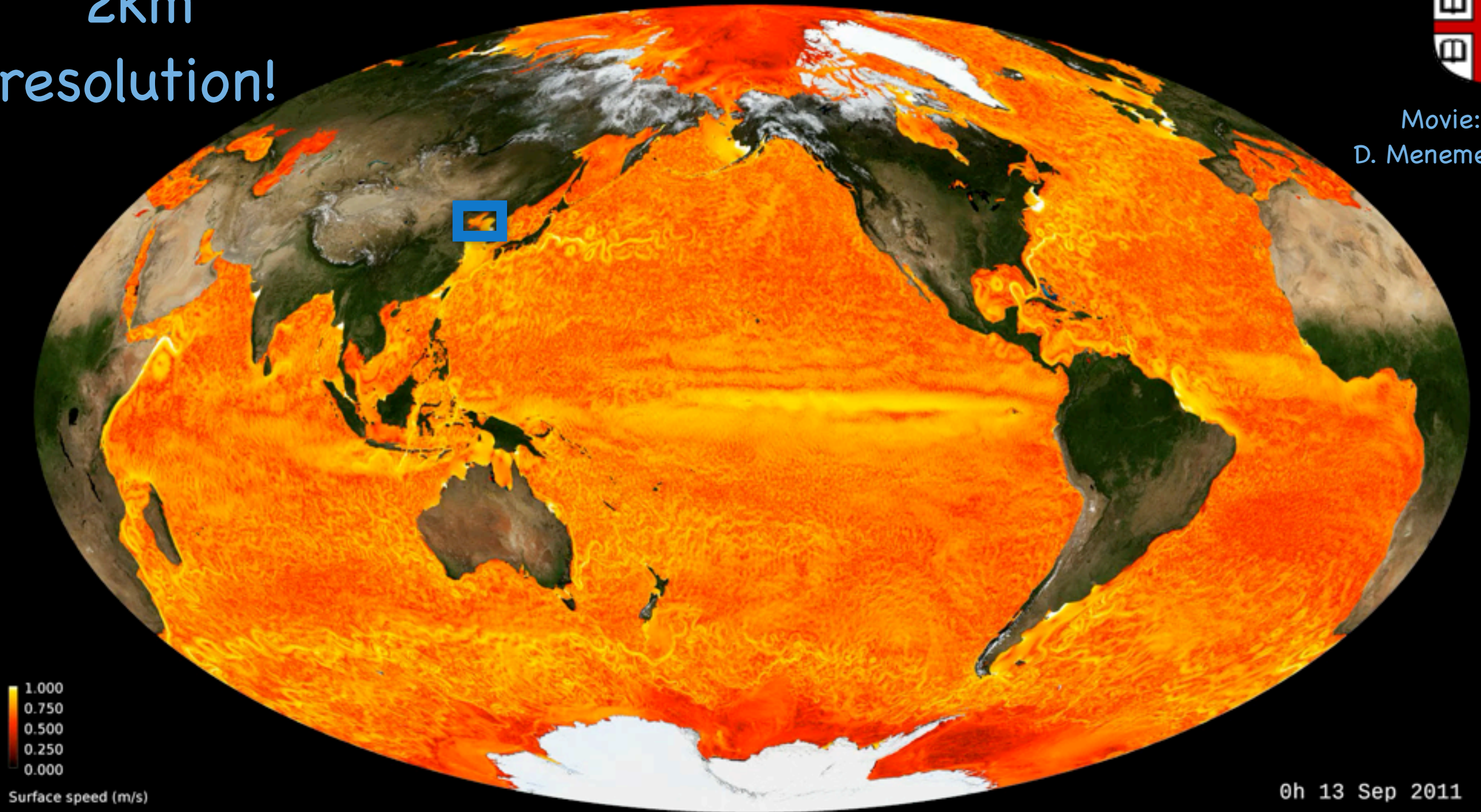
18km resolution



2km
resolution!



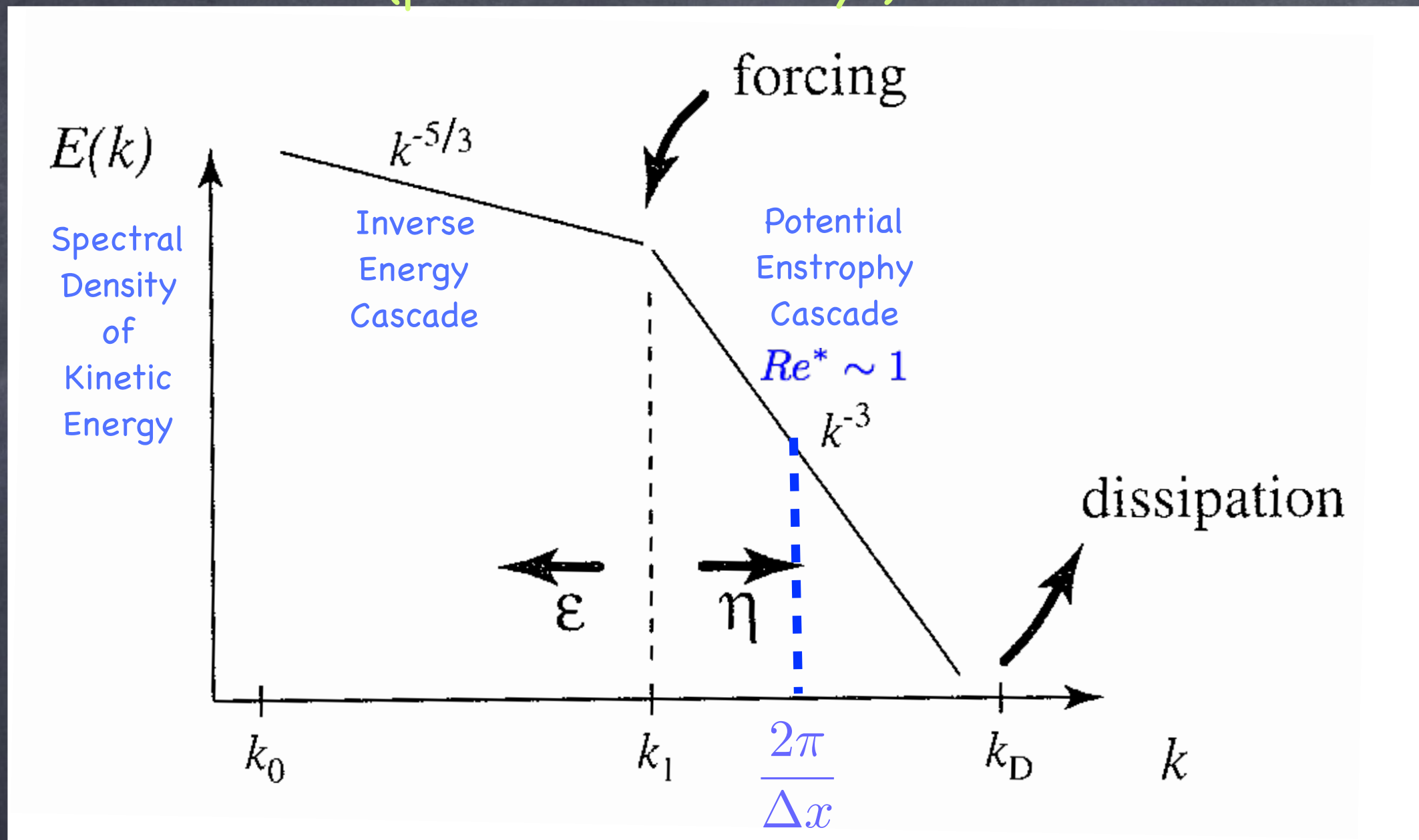
Movie:
D. Menemenlis



QG Turbulence: Pot'l Enstrophy cascade

(potential vorticity²)

J. Charney, 1971 JAS



QG Leith:

$$v_{qg}^* = \left(\frac{\Lambda_{qg} \Delta x}{\pi} \right)^3 |\nabla q_{qg}^*|$$

$$q_{2d}^* = f + \hat{k} \cdot \nabla \times u^*$$

$$q_{qg}^* = f + \hat{k} \cdot \nabla \times u^* + \frac{\partial}{\partial z} \frac{f^2}{N^2} b^*$$

B. Fox-Kemper and D. Menemenlis. Can large eddy simulation techniques improve mesoscale-rich ocean models? In M. Hecht and H. Hasumi, editors, Ocean Modeling in an Eddy Regime, volume 177, pages 319-338. AGU Geophysical Monograph Series, 2008.

QG vs. 2D

$$\nu_{qg}^* = \left(\frac{\Lambda_{qg} \Delta x}{\pi} \right)^3 |\nabla q_{qg}^*|$$

Different (Pot'l) Vorticity Gradients:

$$q_{2d}^* = f + \hat{k} \cdot \nabla \times u^*$$

$$q_{qg}^* = f + \hat{k} \cdot \nabla \times u^* + \frac{\partial}{\partial z} \frac{f^2}{N^2} b^*$$

Also, different implications, because relative vorticity, buoyancy, T, S dissipation now must be consistent with PV:

$$\frac{Dq_{qg}^*}{Dt} = -\nabla \cdot \overline{u'q'_{qg}} \approx \nabla \cdot [\nu^* \nabla q_{2d} + \kappa_{gm}^* \nabla (q_{qg} - q_{2d})] \rightarrow \kappa_{gm}^* = \nu^* = \kappa_i^*$$

QG vs. 2D

$$v_{qg}^* = \left(\frac{\Lambda_{qg} \Delta x}{\pi} \right)^3 |\nabla q_{qg}^*|$$

Different Vorticity Gradients

$$q_{2d}^* = f + \hat{k} \cdot \nabla \times u^*$$

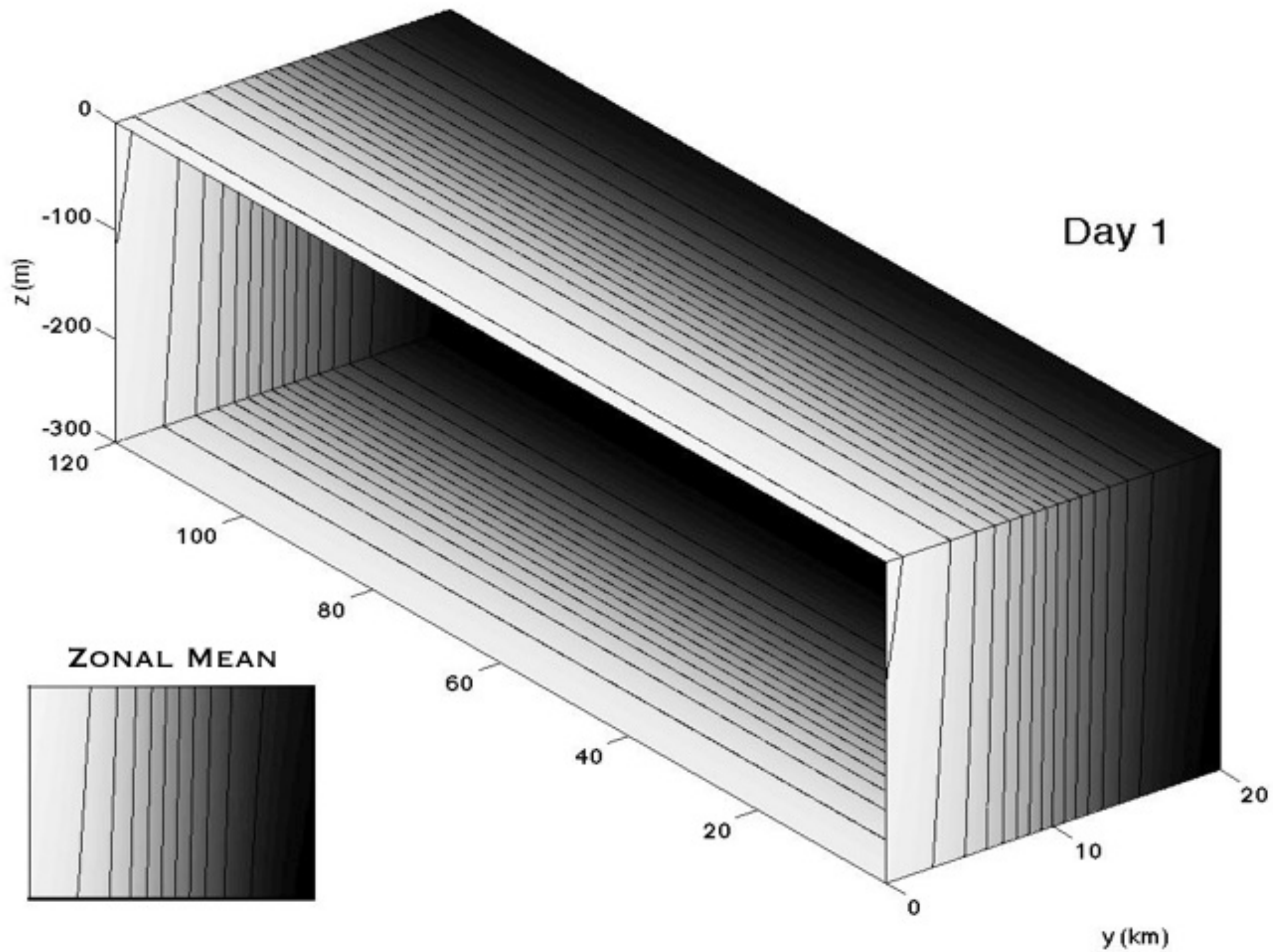
$$q_{qg}^* = f + \hat{k} \cdot \nabla \times u^* + \frac{\partial}{\partial z} \frac{f^2}{N^2} b^*$$

stretching—needs “taming” where QG is a bad approx (equator, boundary layers, etc.)

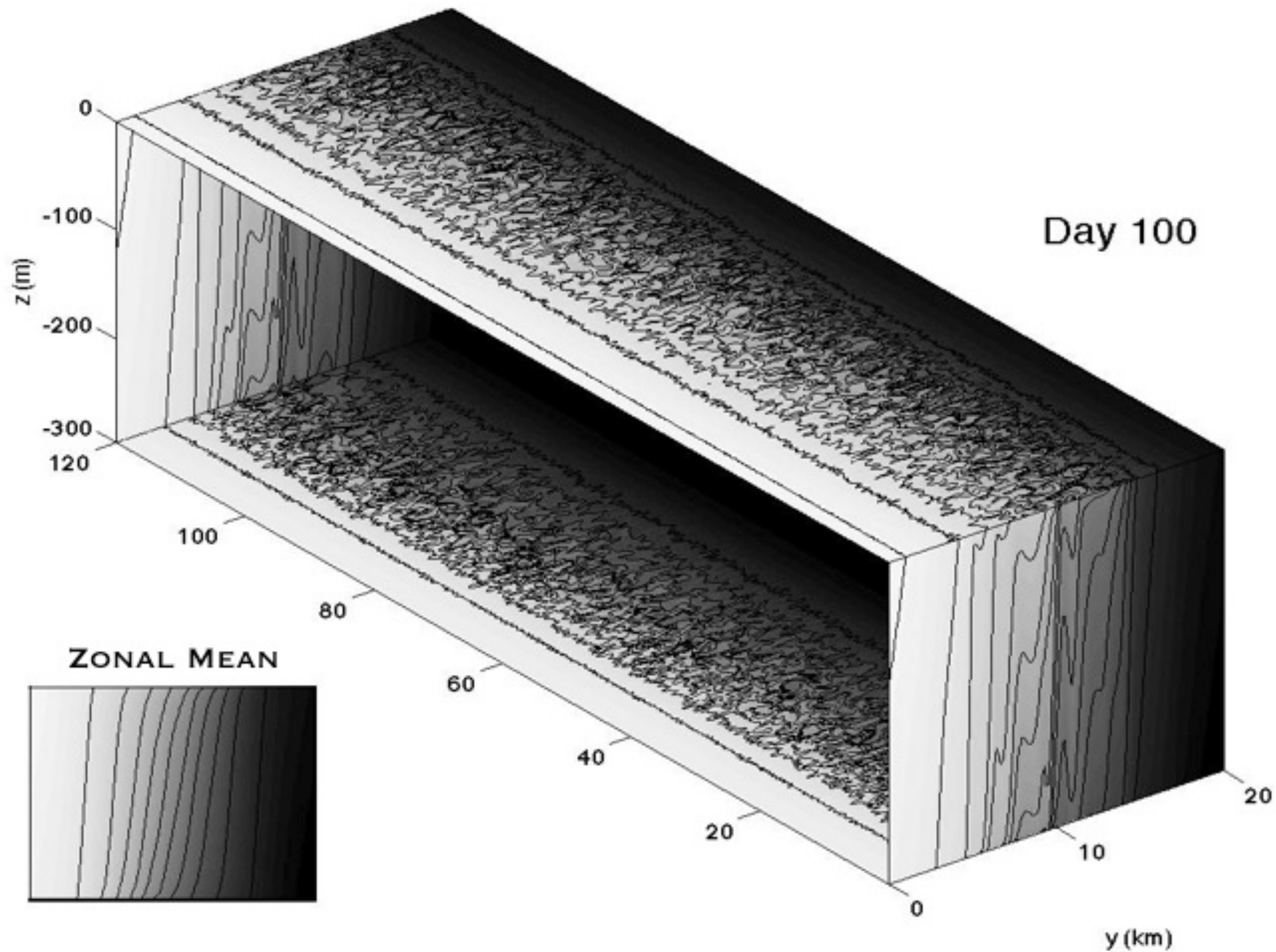
Use gridscale nondims to determine when on the fly

$$Ro^* = \frac{U^*}{f \Delta x} \quad Bu^* = \frac{N^{*2} \Delta z^2}{f^2 \Delta x^2} \sim Ro^{*2} Ri^*$$

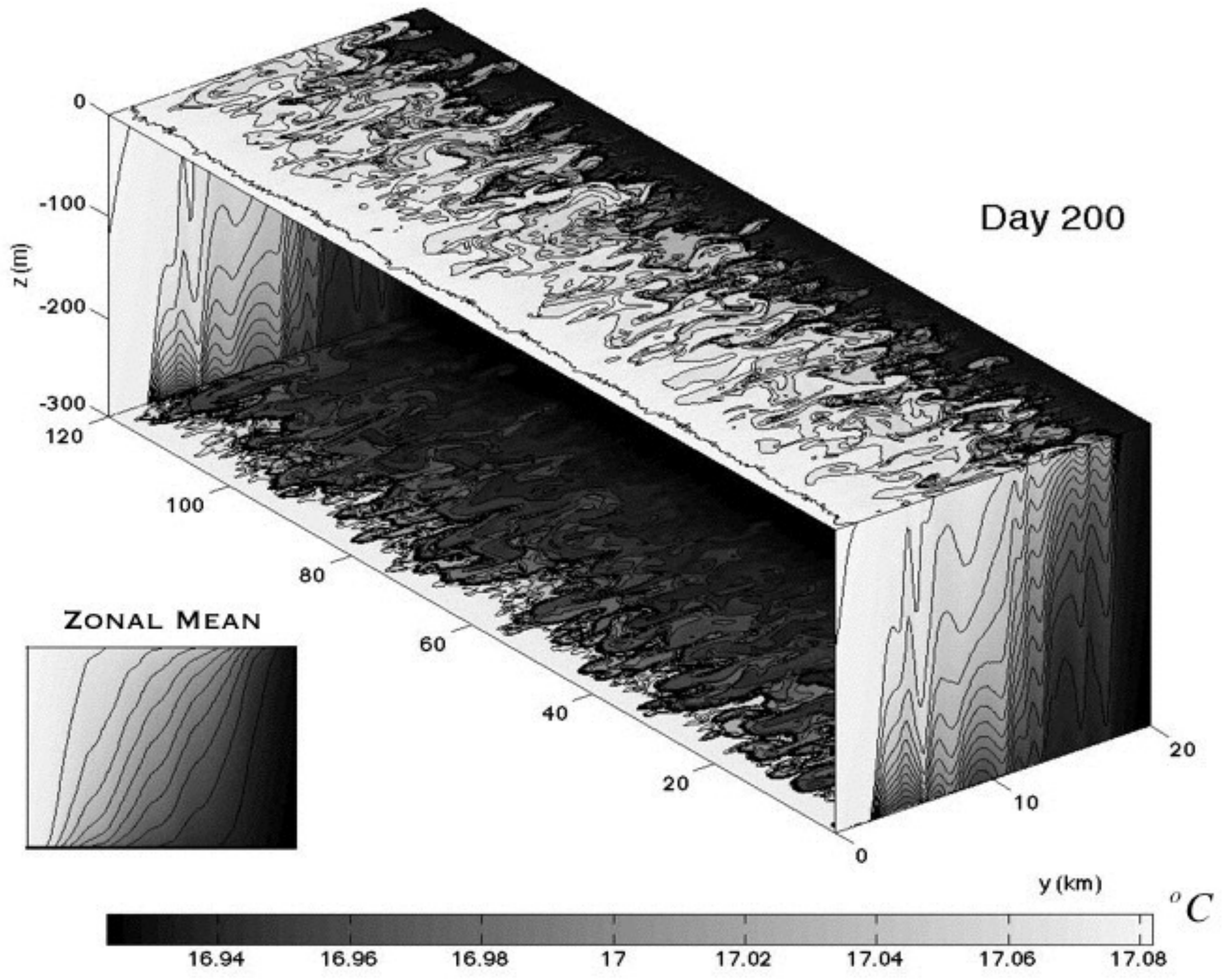
a)

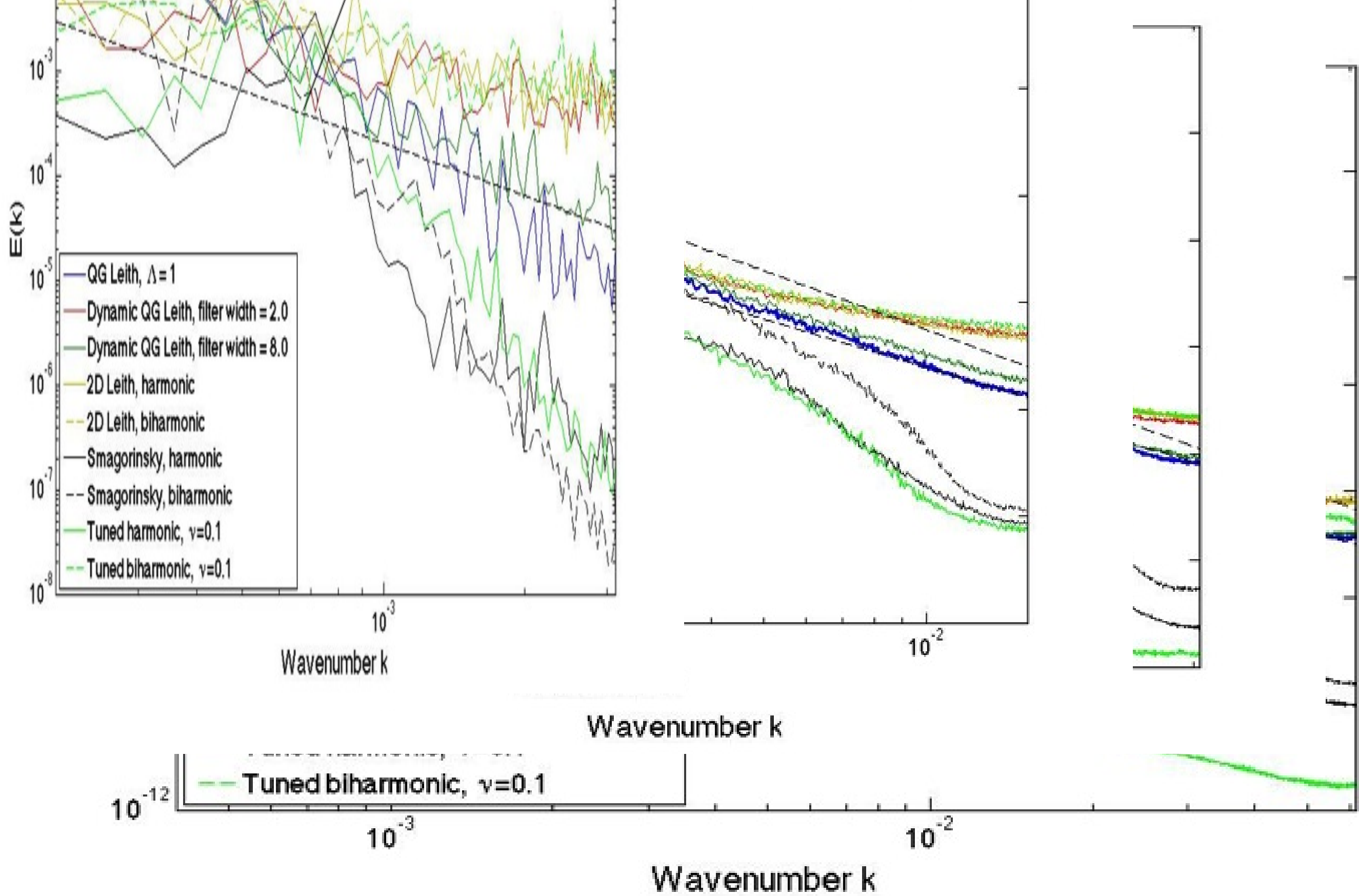


b)



c)





$$L_d = 0.4\Delta x$$

$$1/4^\circ$$

$$L_d = 2.5\Delta x$$

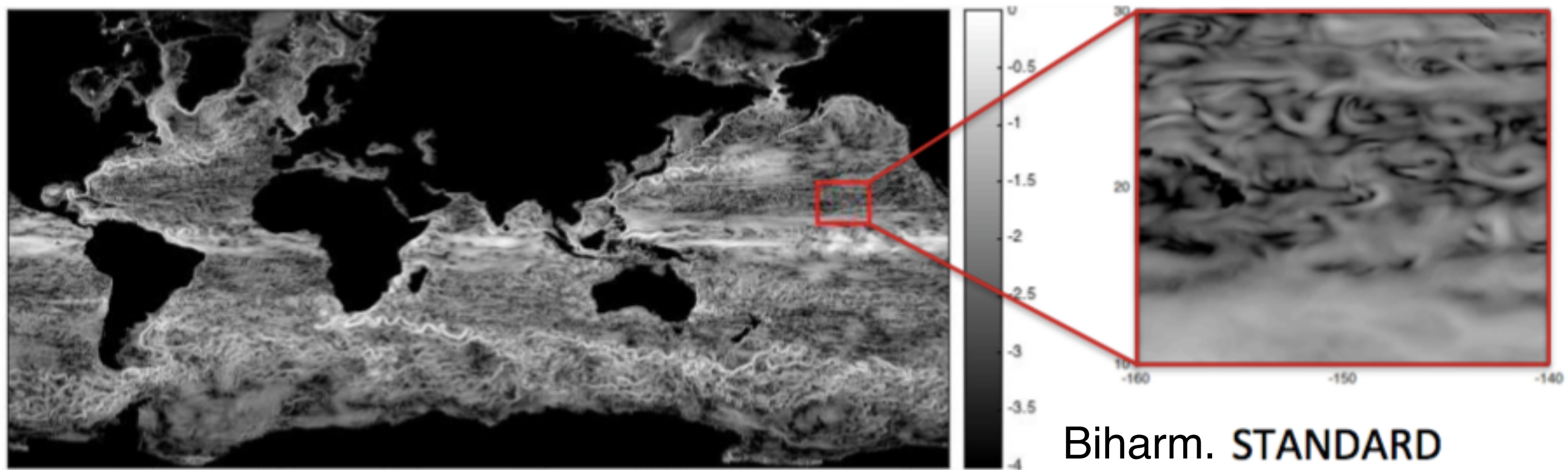
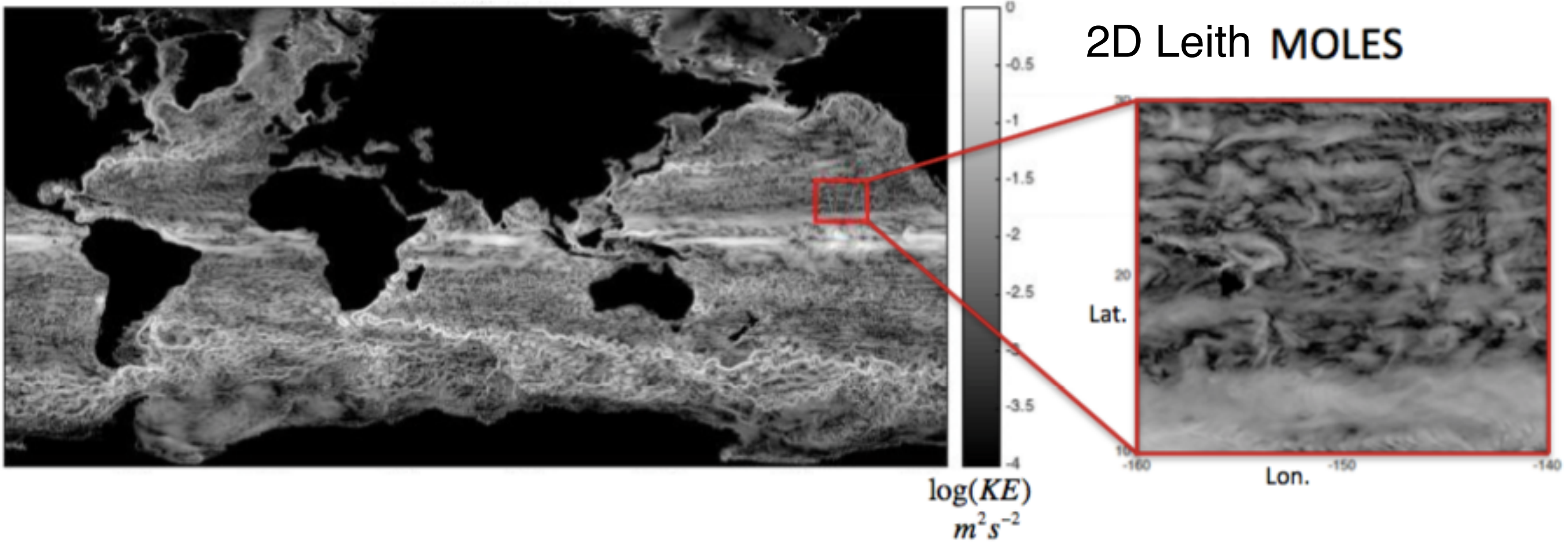
$$1/25^\circ$$

$$L_d = 5\Delta x$$

$$1/50^\circ$$

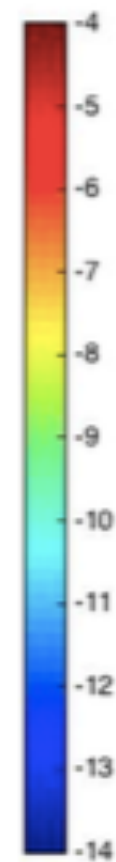
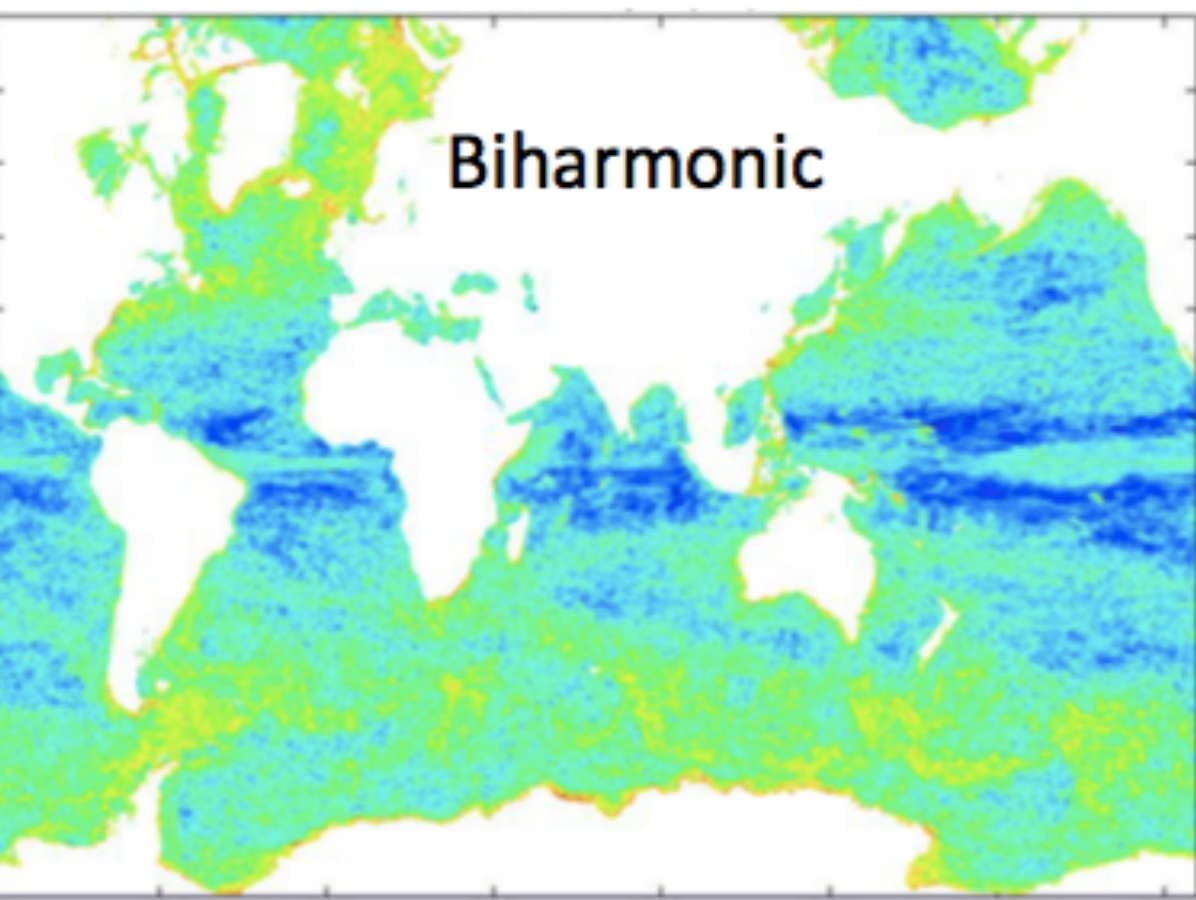
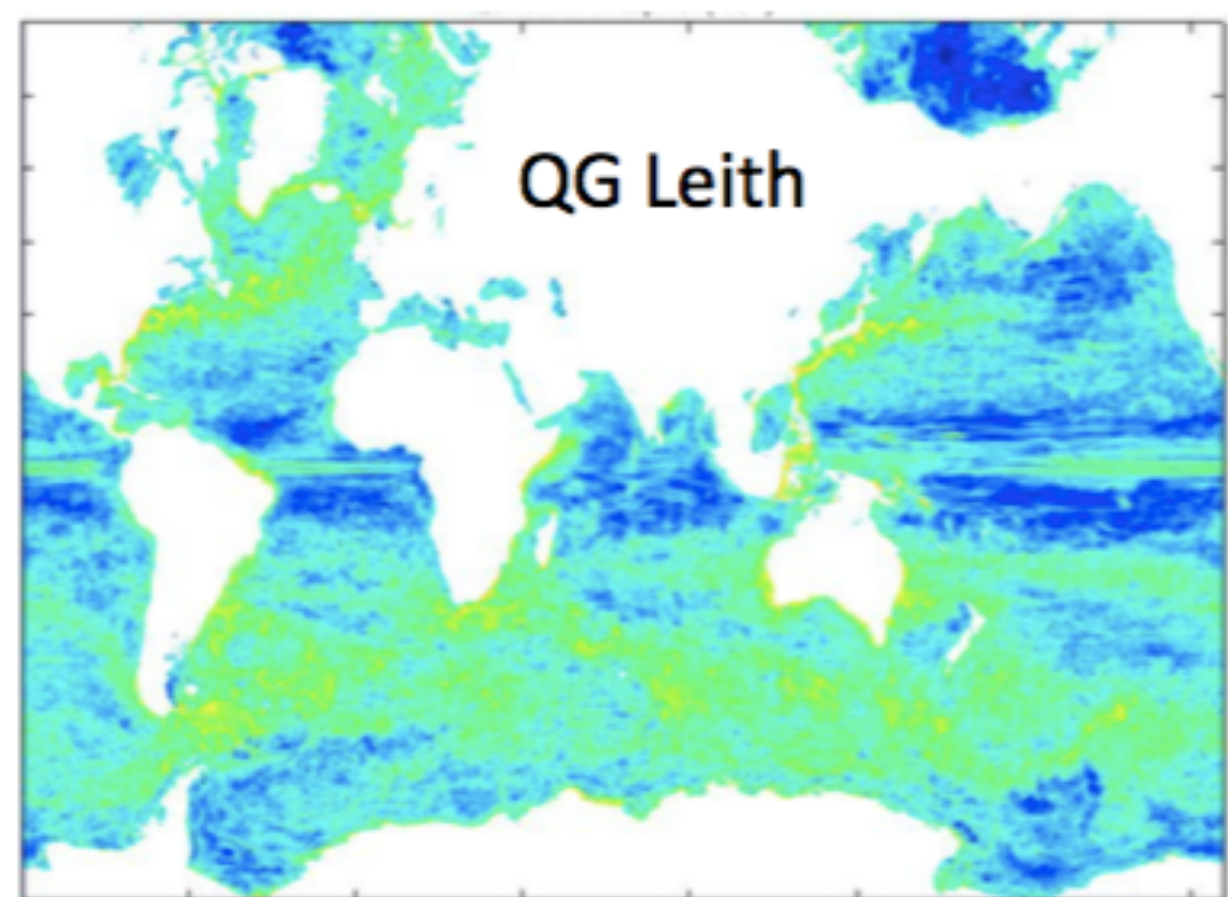
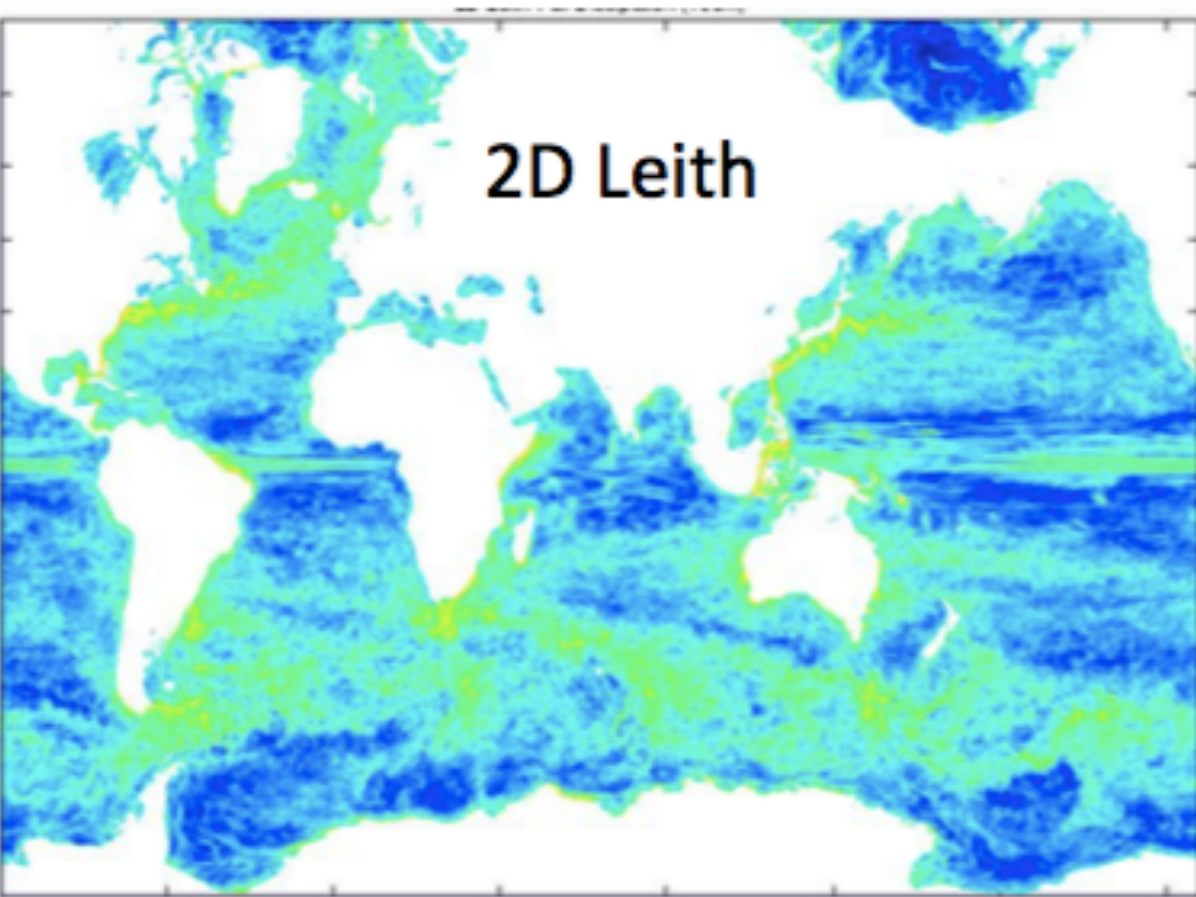
$$L_d = 10\Delta x$$

$$1/100^\circ$$



More EKE and Small Structures in MOLES

B. Pearson, S. Bachman, BFK. Global Application of a Scale-Aware Subgrid Model for Oceanic Quasigeostrophic Turbulence. In prep.



$\log_{10}(\epsilon [m^2 s^{-3}])$

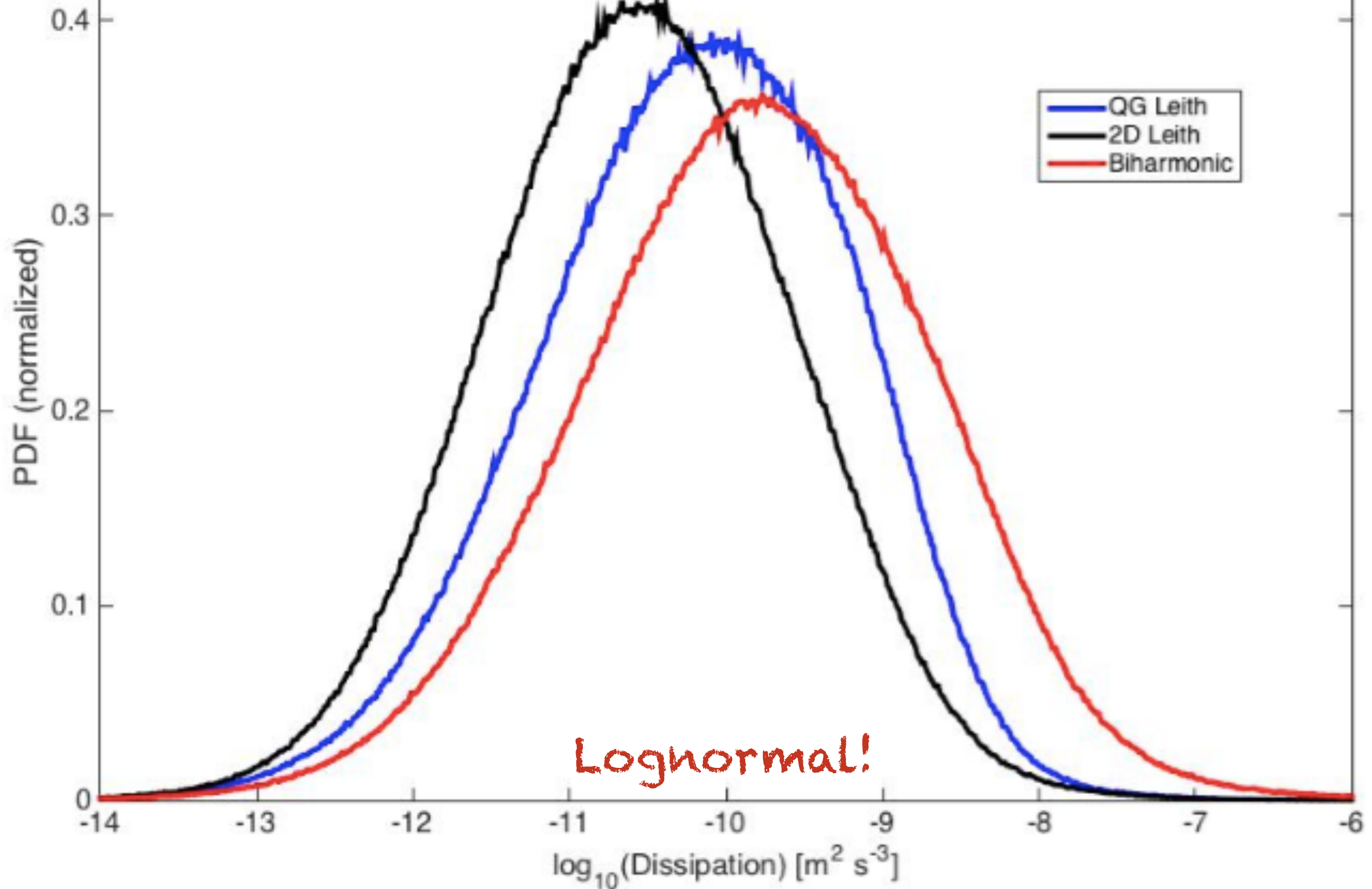
KE Dissipation

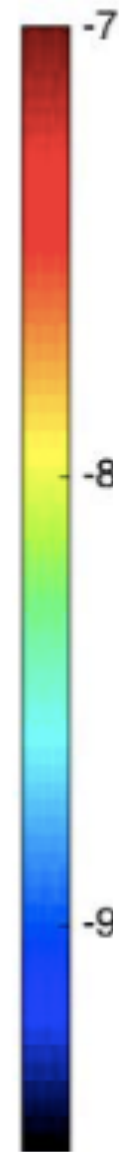
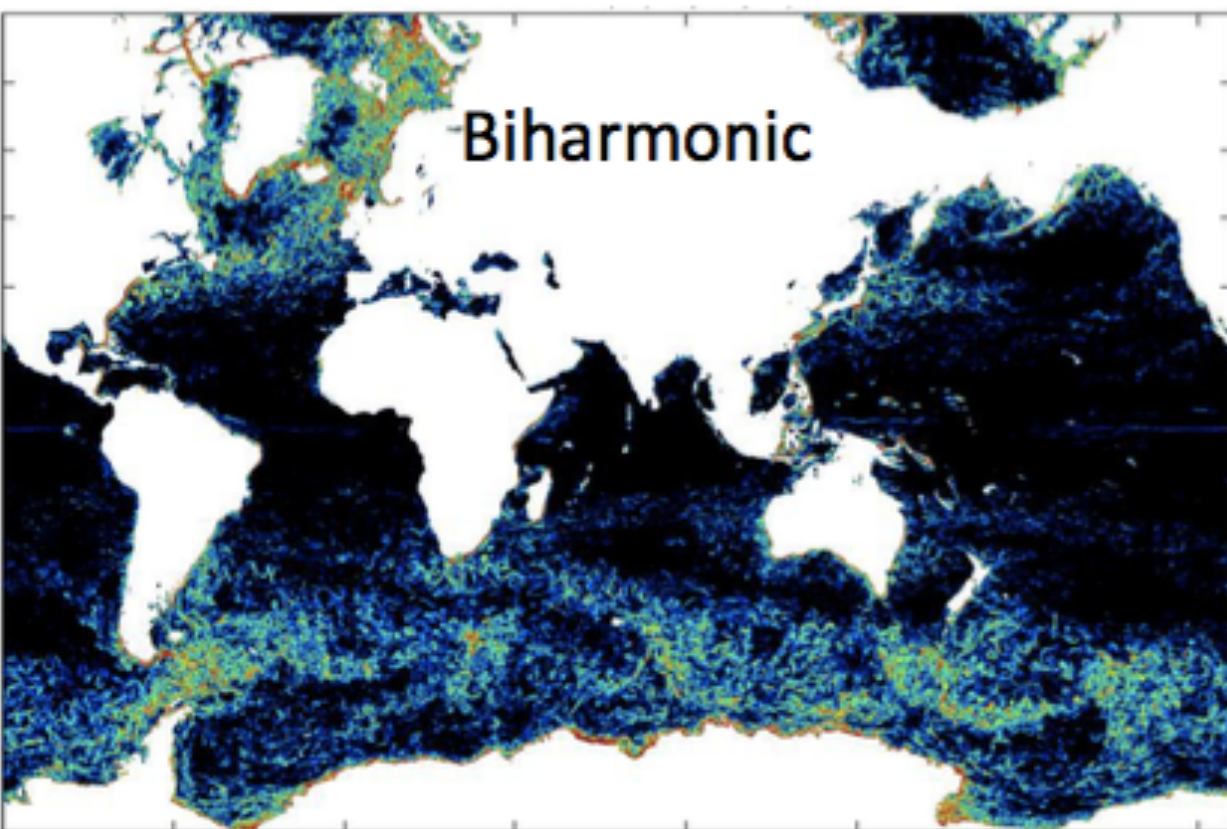
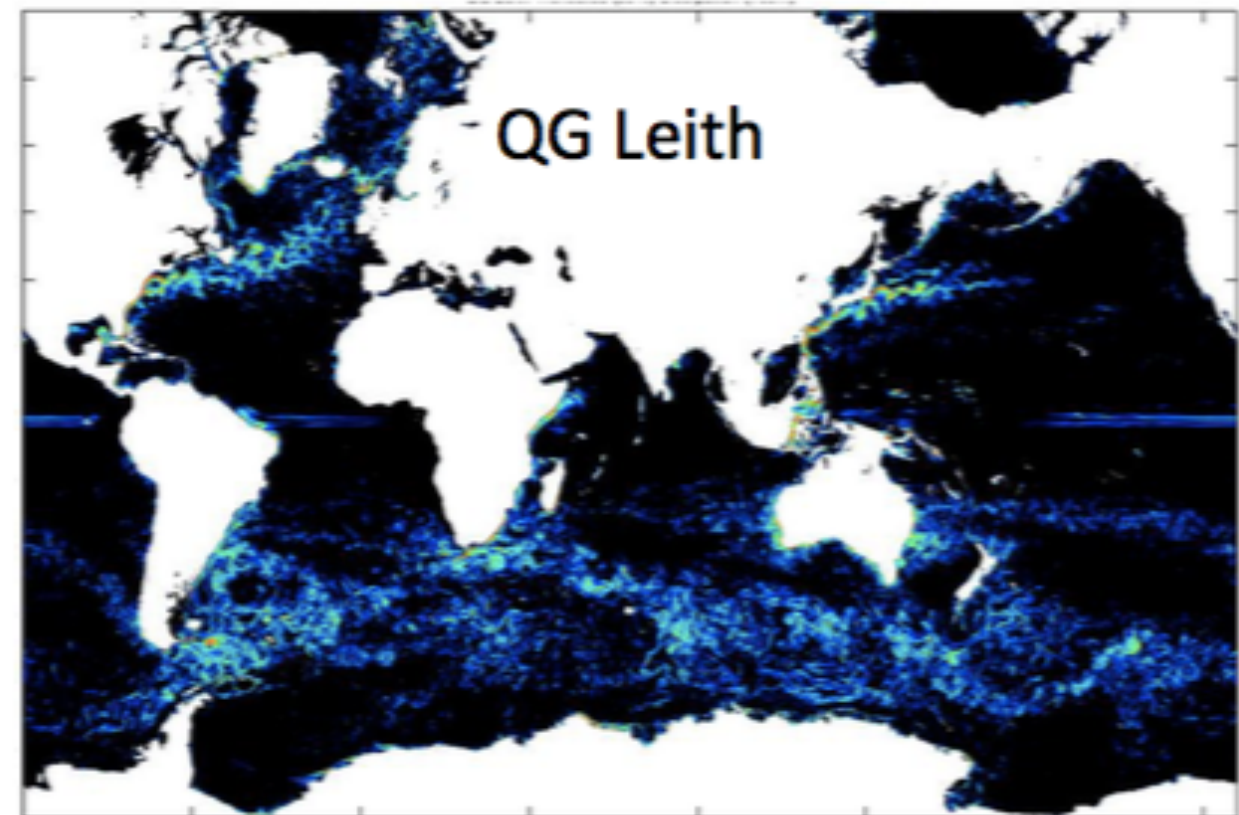
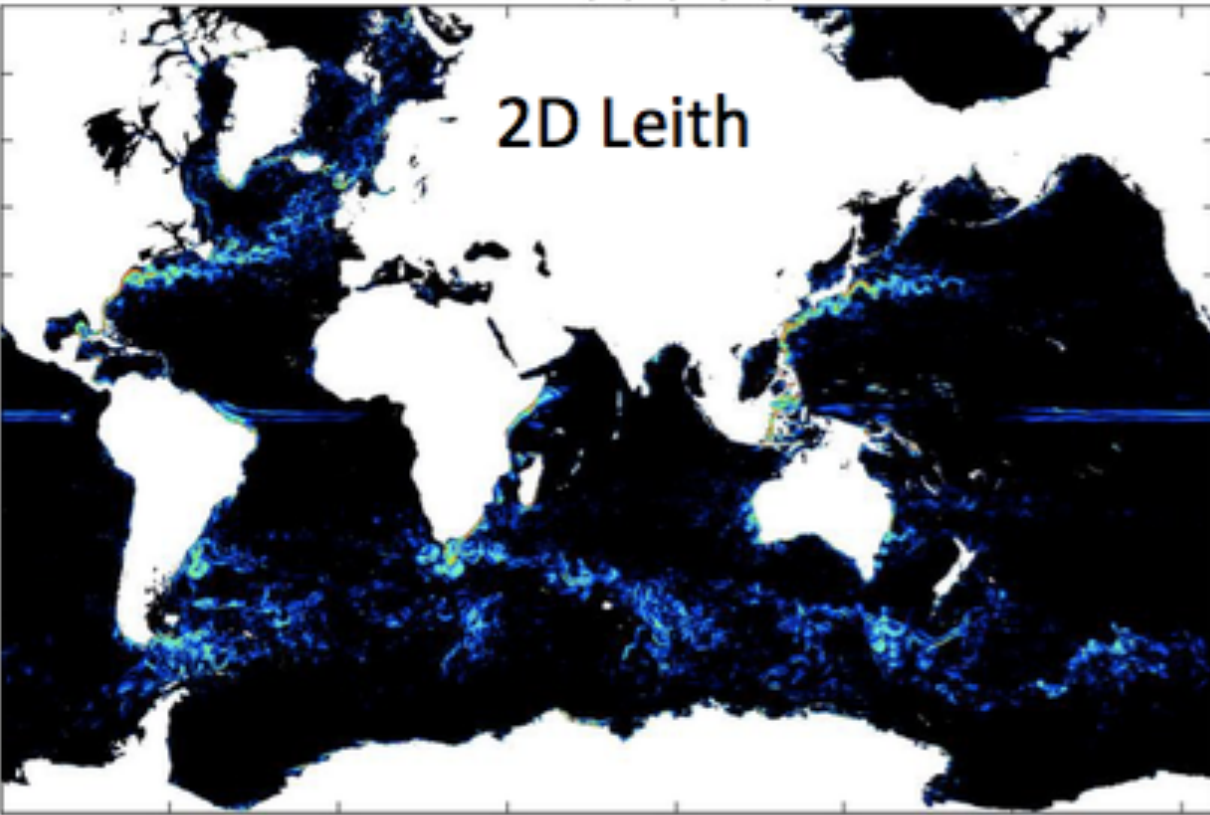
$$\epsilon_{2d} = \nu_{2d} \left(|\nabla_h u|^2 + |\nabla_h v|^2 \right)$$

$$\epsilon_{qg} = \nu_{qg} \left(|\nabla_h u|^2 + |\nabla_h v|^2 \right)$$

$$\epsilon_{bh} = \nu_{bh} \left(|\nabla_h^2 u|^2 + |\nabla_h^2 v|^2 \right)$$

Probability Distribution of KE Dissipation





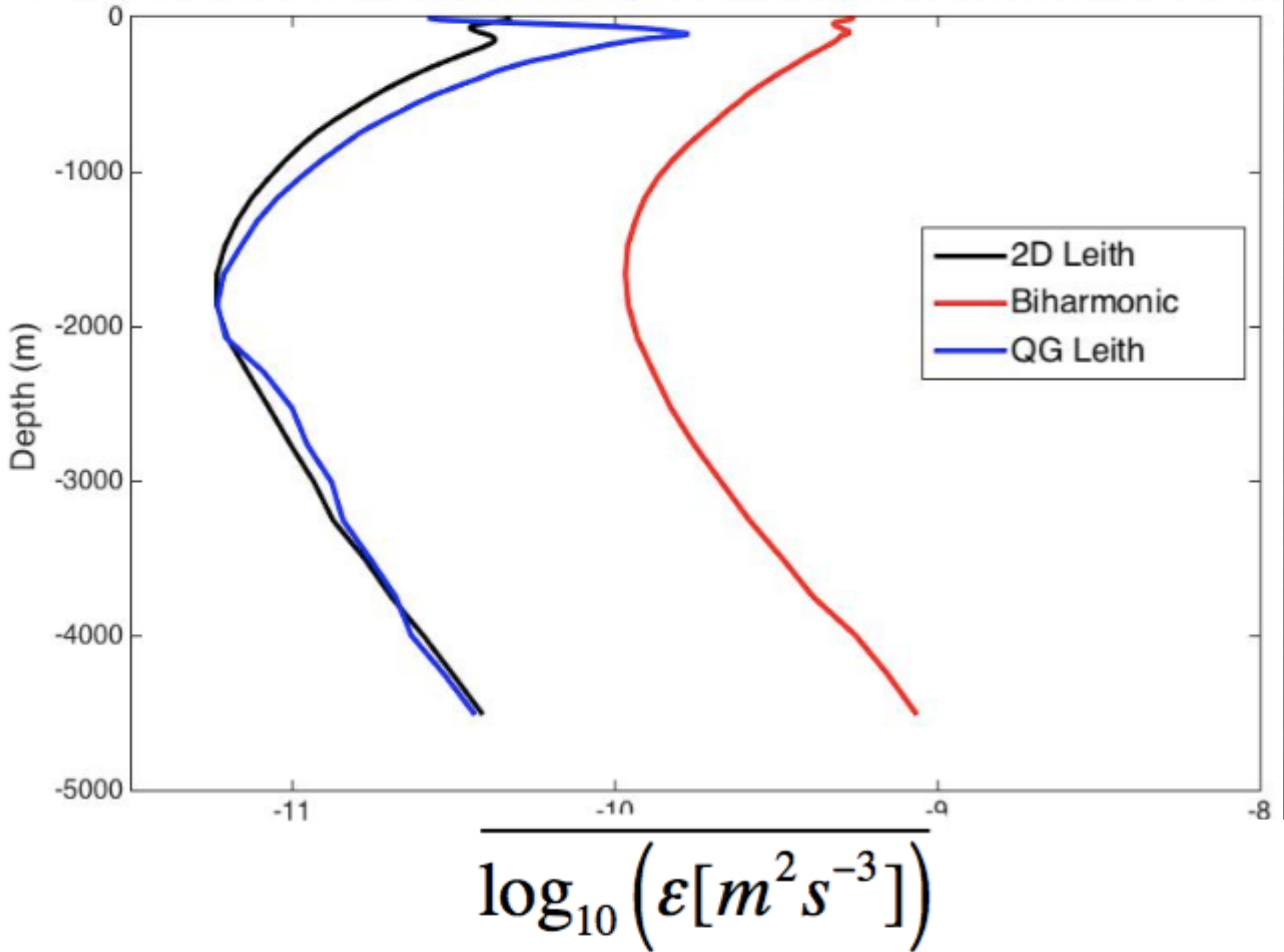
$\log_{10}(\epsilon[m^2s^{-3}])$

A Consequence of Lognormal Statistics—limited regions do most of the work!

Number of points providing 95% of dissipation at 100 m varies between models:
Biharmonic (13% of data points for 95% of dissipation),
2D Leith (23%)
QG Leith (31%)

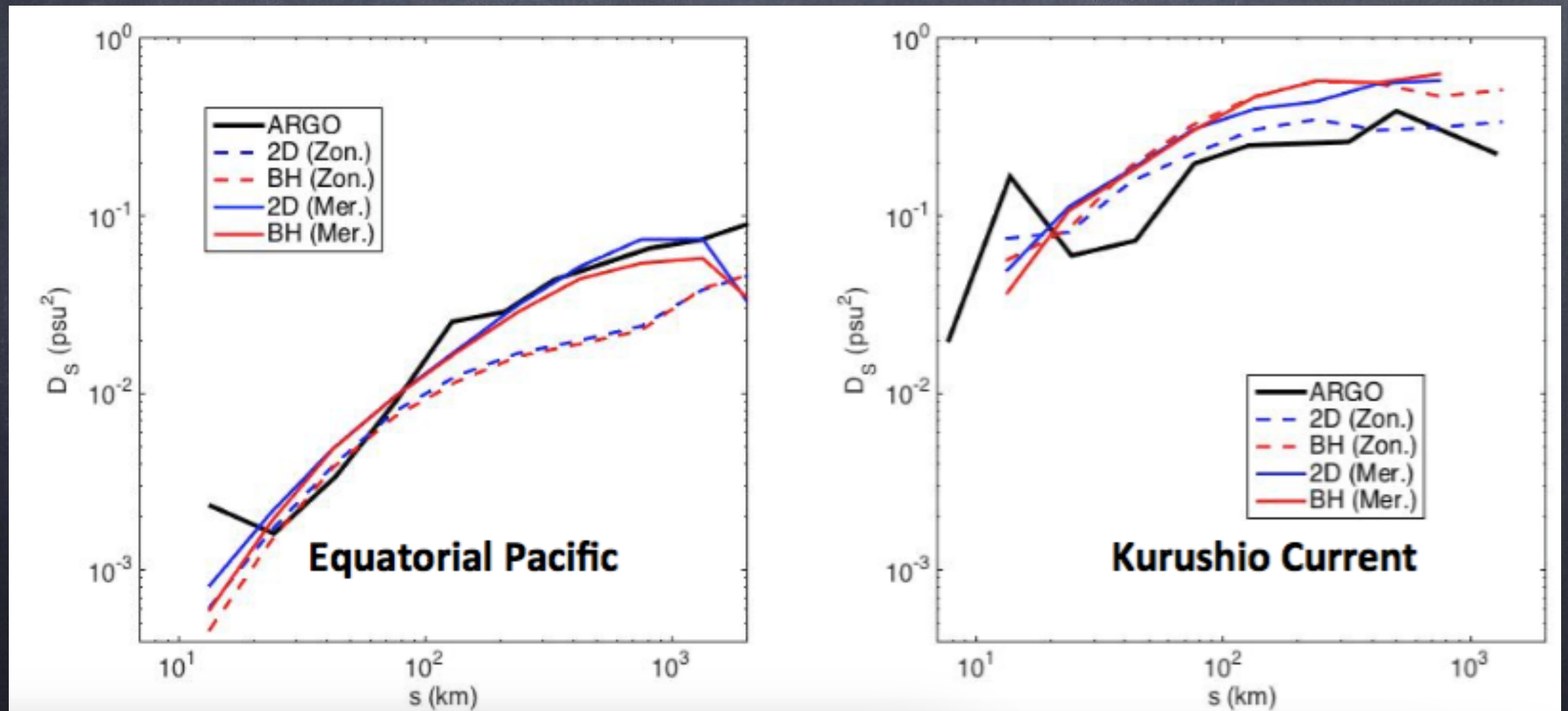
KE Dissipation in Vertical

Dissipation profiles averaged approximately meridionally around the Southern Ocean ($j=$



New Benchmark: Structure Functions

$$D_S(s) = \overline{[S'(x) - S'(x + s)]^2},$$

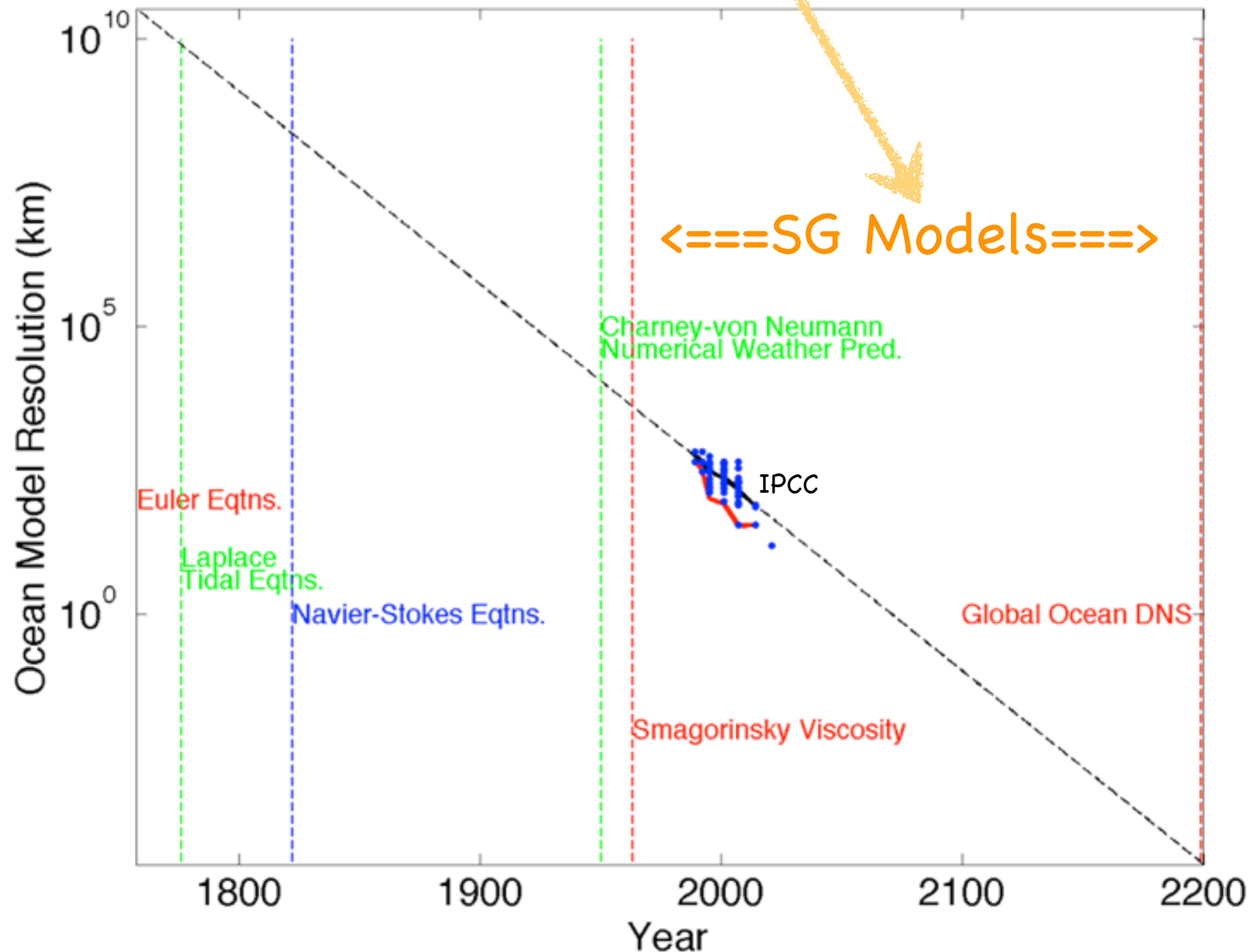


K. McCaffrey, B. Fox-Kemper, and G. Forget. Estimates of ocean macro-turbulence: Structure function and spectral slope from Argo profiling floats. *Journal of Physical Oceanography*, 45(7):1773-1793, July 2015.

Conclusions

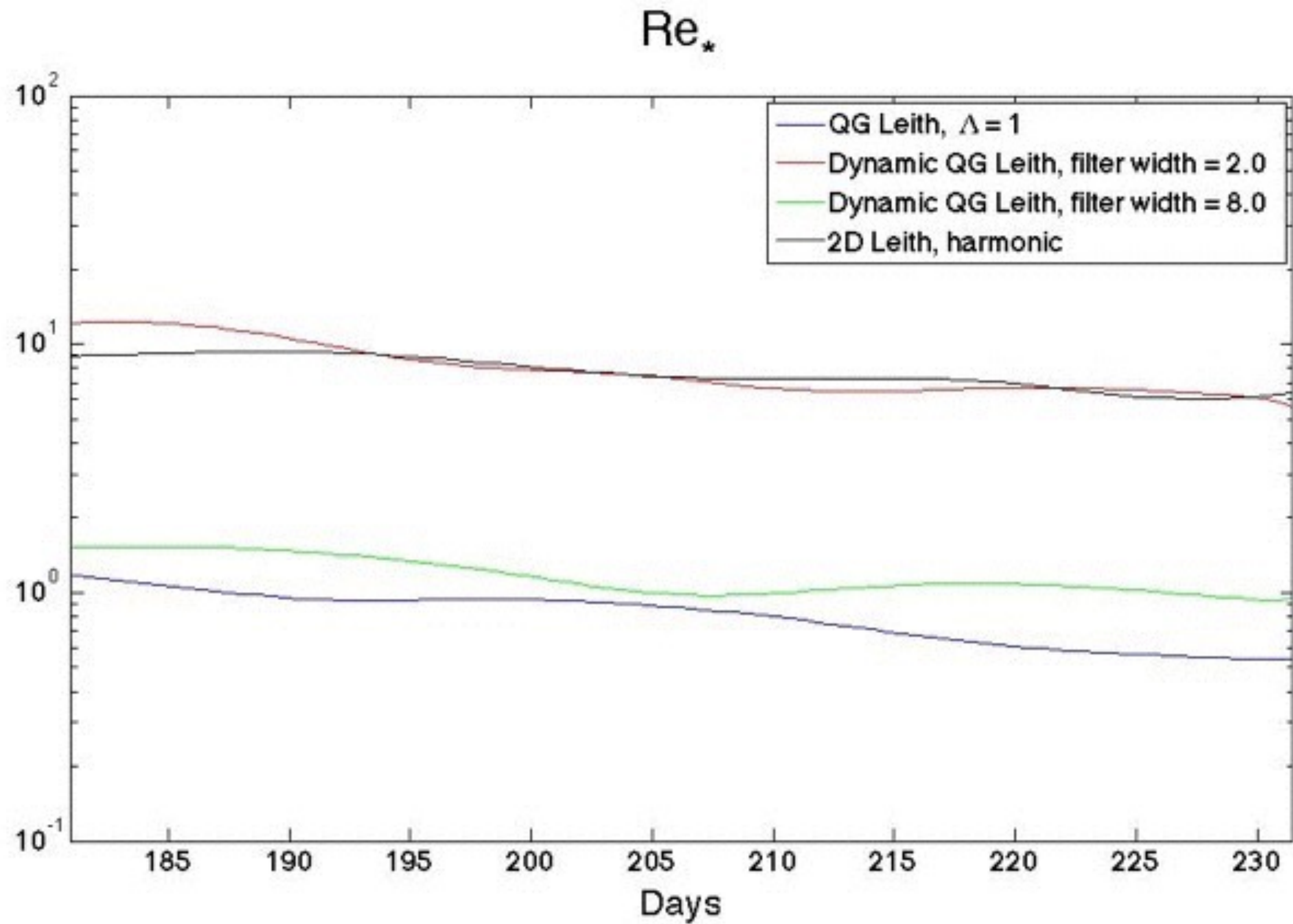
- It is best to think of high-res simulations as “large eddy simulations”.
- Then, take advantage of resolved flow and scaling for physically-based subgrid schemes.
- QG theory has provided such a scheme for mesoscale-permitting to resolving simulations.
- 10x less dissipative than biharmonic viscosity and dissipates where theory suggests it should do.
- Small scales are more energetic, salinity variance can be doubled, even at $O(1000\text{km})$ scales).
- Dynamic version more expensive, no better

Extrapolate for historical perspective: The Golden Era of Subgrid Modeling is Now!



All papers at: fox-kemper.com/research

b)

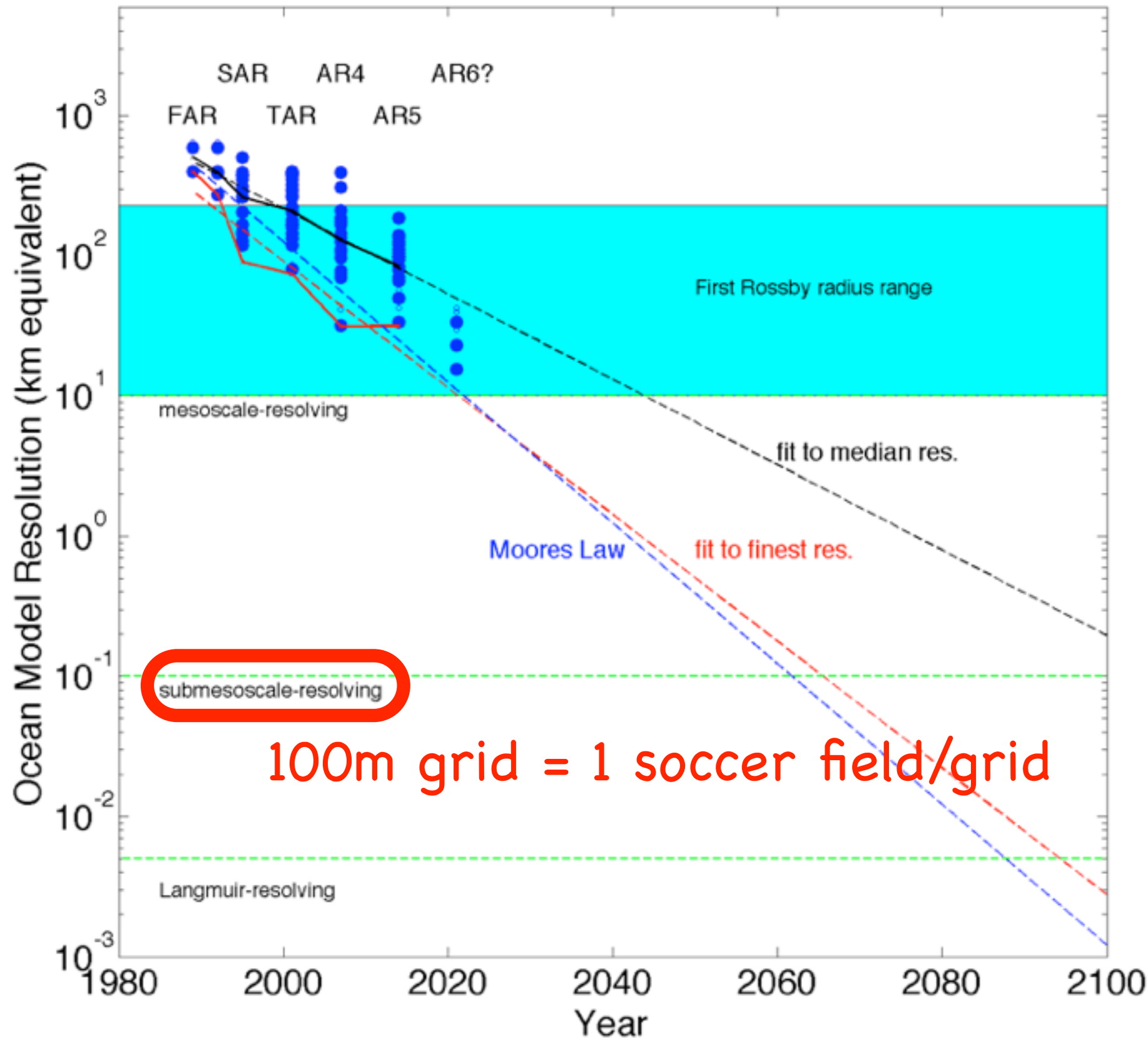


What about modeling important processes in climate models?

Don't we have big enough computers? or won't we soon?



Resolution of Ocean Component of Coupled IPCC models



Here are the collection of IPCC models...

If we can't resolve a process, we need to develop a parameterization or subgrid model of its effect