Demonstrated Sensitivity to Langmuir Mixing in a Global Climate Model (CCSM)

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NASA ROSES Physical Oceanography NNX09AF38G, NSF DMS, and OCE 0934737 "Essentially, all models are wrong, but some are useful" - George Box

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All *parametrizations* are wrong, but some are useful

Talking Points

- Performed a climate modeling version of "back of the envelope" calculation for adding Langmuir mixing to the Community Climate System Model (CCSM)
- CCSM is sensitive to our "back of the envelope" parametrization"
- CCSM is sensitive to the details of the parametrization
- Essential to implement correctly and suggestions welcome

Langmuir Mixing and GCMs

Langmuir mixing is not directly included in any global climate model (2/2010)

- Mixing models (KPP) are trained against data which contains Langmuir mixing
- Is this sufficient? Why important for GCMs?
- Expect large areas of Langmuir mixing in the Southern Ocean





Inverse Turbulent Langmuir Mixing Number The inverse turbulent Langmuir mixing number (La_i) accounts for nonaligned wind and wave fields.

• Mixing can be characterized by the turbulent Langmuir number:

$$La\equiv\sqrt{u^*/u_{stokes}}$$

• A new mixing number *La_i* is defined as

$$La_{i} = \begin{cases} \left(\frac{\mathbf{u}_{stokes} \cdot \mathbf{u}^{*}}{|\mathbf{u}^{*}|^{2}}\right)^{1/2}, & |\theta| < \pi/2; \\ 0, & |\theta| \ge \pi/2. \end{cases}$$

where θ is the difference in wind and wave directions

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Snapshot of La_i^2



 $\mathbf{u}^* = \sqrt{|\tau|/\rho} \frac{\tau}{|\tau|}, \qquad |\mathbf{u}_{stokes}| = \frac{16\pi^3}{g} \int_0^{2\pi} \int_0^{\infty} f^3 S(f,\theta) \, df \, d\theta, \qquad La_i^2 = \frac{\mathbf{u}_{stokes} \cdot \mathbf{u}^*}{|\mathbf{u}^*|^2}$

A Simple Climatology

A preliminary La_i climatology was developed using NOAA WaveWatch III (NWW3) output.



- *Left bottom:* NWW3 0hr forecast significant wave heights and wind speeds (2008/5/21 00:00)
- *Left top:* Estimated regions of Langmuir mixing (La_i^2)

Below: Zonal mean inverse Langmuir (*La_i*) climatologies



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A Simple Scaling for Langmuir Depth/Entrainment



- The Algorithm uses *Fr* to determine *H*. If *H* is deeper than the KPP boundary layer depth (BLD), *H* is used.
- Large came up with clever choices for N, H that lead to a robust implementation in KPP. With these choices, H and BLD converge over time.

Sensitivity to Inclusion

There is a persistent shallow mixed-layer bias in the Southern Ocean in all GCMs.

• Addition of Langmuir parametrization helps correct this bias





Motivation and Applications of Coupling a Wave Model

- Calculate Langmuir Mixing forcing prognostically
 - A coupled wave model will allow use of more sophisticated and validated parametrizations (e.g., Smyth et al, 04; Harcourt & D'Asaro, 08; Grant & Belcher, 09)
- Improve the air-sea momentum flux
- Improve the air-sea tracer flux
- Conduct climate change studies like coastal erosion
- Others suggestions?

Conclusion

- Added a simple parametrization for Langmuir mixing to NCAR's CCSM
- CCSM is sensitive to our parametrization and has potential to correct long-standing biases
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Please check out poster PO35F-01 by Erik Baldwin-Stevens for more details about a newer Langmuir climatology

Calculating Stokes drift at the surface

From previous work by Kenyon (1969) and McWilliams & Restrepo (1999), we can calculate Stokes drift using the full 2-D spectrum as

$$|\mathbf{u}_{stokes}| = \frac{16\pi^3}{g} \int_0^{2\pi} \int_0^{\infty} f^3 S(f,\theta) \, df \, d\theta$$

For monochromatic waves, this simplifies to

$$|\mathbf{u}_{stokes}| = \frac{\pi^3 H s^2}{g T m^3}$$

where $Hs = 4\sqrt{m_0}$ and m_0 is the zeroth moment of variance.

Refining our Stokes Drift Approximation

We would like to approximate Stokes drift for different products without using the wave spectrum.

• Currently examining if there is an empirical or mathematical relationship similar to the monochromatic form

$$\mathbf{u}_{stokes} \approx a(f) \frac{\pi^3 H s^2}{g T m^3} \, \hat{\mathbf{e}}_d$$



Different Definitions of Mean Wave Period

WaveWatch:
$$Tm_0 = \overline{(f^{-1})}$$
ERA40: $Tm_1 = 1/(\overline{f})$ TOPEX: $Tm_2 = 1/\sqrt{(\overline{f^2})}$

$$m_n = \int_0^{2\pi} \int_0^\infty f^n S(f,\theta) \, df \, d\theta$$
$$Tm_0 = \frac{m_{-1}}{m_0}, \quad Tm_1 = \frac{m_0}{m_1}, \quad Tm_2 = \left(\frac{m_0}{m_2}\right)^{1/2}$$

Using Assumed Forms of the Spectrum

Pierson-Moskowitz Spectrum

$$S(f,\theta) = S(f) = rac{lpha g^2}{(2\pi)^4} f^{-5} \operatorname{Exp}\left[-rac{5}{4}\left(rac{f_p}{f}
ight)^4
ight]$$

where α is the Phillips constant and f_p the peak frequency



1994-7 January zonal mean data for WW3 (black), ERA40 (red), and TOPEX (blue)

Wave Watch Details



Blank Spectral Polar Plot

In WAVEWATCH III, every bin holds a value for wave energy. The model tracks this energy at every time step for every grid point across the global oceans, accounting for wave generation and swell propagation. After the model run is complete, the spectrum for selected grid points is downloaded and contours are drawn for wave energy. The resulting plois <u>Look like this</u>.

- 3rd generation wave model
- Solves the spectral action density balance equation
- 15-20 sec per time step (1 hr) for one processor (\approx 35-50 hr/yr)
- Plan on scaling back the number of bins significantly and turning off some interactions
- Aternative 2nd generation model (George Mellor, Princeton) worth exploring