

Abstract

We demonstrate the use of information theory metrics, Shannon entropy and mutual information, for measuring internal and forced variability in general circulation coastal and global ocean models. These metrics have been applied on spatially and temporally averaged data. A combined metric reliably delineates intrinsic and extrinsic variability in a wider range of circumstances than previous approaches based on variance ratios that therefore assume Gaussian distributions. Shannon entropy and mutual information manage correlated fields, apply to any distribution, and are insensitive to outliers and a change of units or scale. Different metrics are used to quantify internal vs forced variability in (1) idealized Gaussian and uniformly distributed data, (2) an initial condition ensemble of a realistic coastal ocean model (OSOM), (3) the GFDL-ESM2M climate model large ensemble. A metric based on information theory partly agrees with the traditional variance-based metric and identifies regions where non-linear correlations might exist. Mutual information and Shannon entropy are used to quantify the impact of different boundary forcings in a coastal ocean model ensemble. Information theory enables ranking the potential impacts of improving boundary and forcing conditions across multiple predicted variables with different dimensions. The climate model ensemble application shows how information theory metrics are robust even in a highly skewed probability distribution (Arctic sea surface temperature) resulting from sharply non-linear behavior (freezing point).

Plain Language Summary

It is important in climate and environmental modeling to distinguish variability caused by external forces versus variability that arises within the system being modeled itself. In this paper, we study multiple runs of a coastal ocean model that are forced by tides, winds, and offshore and atmospheric conditions and multiple runs of climate model simulations that are forced by greenhouse gases and solar warming. We use information theory—a way to count the number of physical states visited by a system under study—to quantify the amount of variability in these models that results from the external forcing versus the amount from the internal chaotic variability. In this way, we can prioritize improvements or inclusion of the different forcings based on how large the model response to them is.

1 Introduction

In an ocean or climate model, it is pertinent to understand the cause of variability, as it leads to implications for predictability, prioritization of data collections for assimilation, and provides an understanding of the dynamics at play in different regions. In a coastal model, variability can arise from extrinsic factors such as wind forcing, solar and thermal forcing, tides, rivers, evaporation, and precipitation, or it can be due to internal chaos inherent to the governing fluid equations (Sane et al., 2021). In a climate model, modes of variability such as El Niño, the North Atlantic Oscillation, or the Southern Annular Mode can conceal or delay the emergence of attributable anthropogenic climate change signals (Milinski et al., 2019). In high-resolution ocean models, internal chaos or intrinsic variability can also be due to eddies (Leroux et al., 2018; Llovel et al., 2018). Accurately quantifying the relative contribution of external and internal factors can help to elucidate the causes responsible for observed variability in models, help to identify key observable metrics, and help quantify concepts such as the time of emergence of climate signals (Hawkins & Sutton, 2012).

Numerous methods exist in the literature to quantify intrinsic and extrinsic variability using models or observations (e.g., Frankcombe et al. (2015); Schurer et al. (2013); Y.-c. Liang et al. (2020)). Two types of model ensembles are common: initial condition ensembles (where the same model is used repeatedly with perturbed initial conditions and intrinsic variability occurs via chaos), and multi-model ensembles (where a variety of models differing in numerics and parameterizations are used to simulate change under the same forcing—in this case “intrinsic” variability also includes aspects of model formulations). Initial condition ensembles are a set of simulations sharing the same forcing and the same governing equations and identical parameterizations, but they still diverge from one another because slightly different initial conditions evolve into substantially different conditions later in the simulations due to intrinsic chaos—most geophysical fluid dynamics models and climate models are intrinsically chaotic. Most of the discussion here will focus on initial condition ensembles, but the metrics proposed can be adapted to both types of ensembles.

To help visualize variability, a generic idealized output from an ocean or atmospheric model is shown in Figure 1. Each color represents a different ensemble member, and the black solid line is the mean of those members. The solid black line is the signal due mainly

71 to extrinsic factors (aside from the limits of the finite ensemble size) and the spread of
 72 the model (schematized by the double-headed magenta arrow in Figure 1) can be con-
 73 sidered due to intrinsic variability or internal chaos.

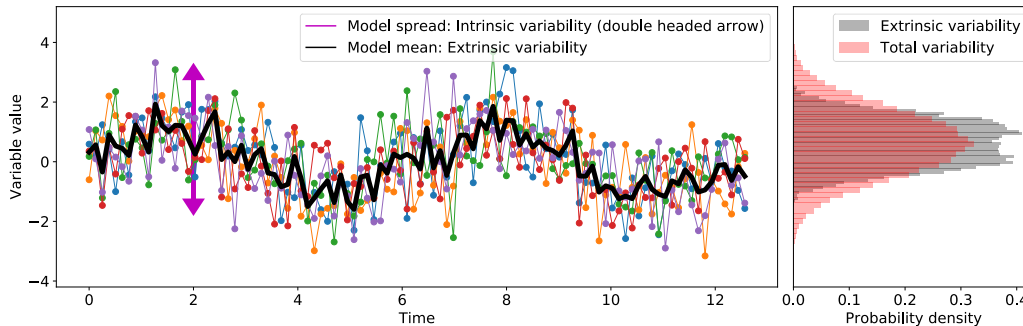


Figure 1. A sketch of a typical ocean or climate model output for an arbitrary variable. Each ensemble is shown in a different color, and the mean of the ensemble is shown as a black line. The ensemble mean can be considered to be the trend set by external forcings. The model spread shown by the double-headed magenta arrow indicates the chaos of the model.

74 One method of quantifying intrinsic and extrinsic variability is to look at variances
 75 (second central statistical moment) of the model spread and the mean of the model (Leroux
 76 et al., 2018; Llovel et al., 2018; Waldman et al., 2018; Yettella et al., 2018). Variance is
 77 sufficient to constrain all metrics of variability about the mean when distributions are
 78 Gaussian and uncorrelated, but a single statistical moment usually measures only part
 79 of a more complex variability distribution. Many climatological variables show non-Gaussian
 80 distributions (e.g., Franzke et al. (2020)). In fact, generalized variance might be mislead-
 81 ing (e.g., Kowal (1971)). Quantification of variability should be robust to or have a known
 82 dependence on changes in the units of the quantity or the scale (e.g., changing temper-
 83 ature from Celsius to Fahrenheit or Kelvin). Comparative metrics, such as intrinsic vs.
 84 extrinsic variability, should not depend on these arbitrary choices of units at all.

85 Variability, in essence, is a function of the number of occurrences or frequency of
 86 occurrence, often estimated by a histogram formed after appropriately binning the data,
 87 which then approximates a distribution with a discrete probability p_i as a fraction over
 88 all states of the visited system. A histogram thus makes the estimated and visited num-
 89 ber of states discrete rather than continuous. Information entropy metrics measure vari-
 90 ability by taking into account the probability distribution of the binned data, drawing

91 on the concept from statistical mechanics of entropy in quantifying the number of mi-
92 crostates that a variable can occupy. The fundamental measure in information theory
93 is the Shannon entropy (Shannon, 1948) (a.k.a. the information entropy) that charac-
94 terizes the amount of variability in a variable (Carcassi et al., 2021). Mutual informa-
95 tion, another metric introduced by Shannon (1948), measures how much information a
96 variable contains about another variable.

97 Information theory is applied in signal processing, computer science, statistical me-
98 chanics, quantum mechanics, etc. It is used to quantify the amount of information, dis-
99 order, freedom, or lack of freedom (Brissaud, 2005). The application of these abstract
100 notions to geophysical flows can have immense practical benefit when information en-
101 tropy is interpreted as a measure of variability, as entropy does not rely on any partic-
102 ular parametric probability distribution. Information theory metrics are not new to cli-
103 mate sciences. They have been introduced in predictability studies, evaluating the skill
104 of statistical models, as well as uncertainty studies (Leung & North, 1990; Schneider &
105 Griffies, 1999; Kleeman, 2002; DelSole & Tippett, 2007; Majda & Gershgorin, 2010; Steven-
106 son et al., 2013) and recently in studying variability (Gomez, 2020), coastal predictabil-
107 ity (Sane et al., 2021) and drivers of drought (Shin et al., 2023).

108 In the two parts of this article, we bring well-established concepts of information
109 theory to the particular application of measuring intrinsic and extrinsic variability for
110 ensemble model runs within atmospheric and oceanographic modeling. We use Shannon
111 entropy and mutual information and a particularly useful combination of the two. We
112 indirectly employ conditional entropy, which depends on Shannon entropy and mutual
113 information but is less intuitive so is not discussed in detail. Recent theoretical advances
114 in understanding dynamical systems through the lens of information theory relate causal-
115 ity analysis and information transfer (e.g., X. S. Liang (2014)). Although important, this
116 theory has had few concrete applications. Even the basic information theory concepts
117 (Shannon entropy and mutual information) have enjoyed only limited adoption by the
118 oceanic and atmospheric community, primarily arising in predictability quantification
119 (e.g., Sane et al. (2021)). We begin to bridge the gap with a pragmatic framework which
120 can be easily replicated and improved upon, including causality analysis and the evo-
121 lution of entropy within modeling systems like those studied here.

122 In Part 1, we apply this intrinsic vs. extrinsic metric to three sets of data: 1) Ide-
 123 alized Gaussian and uniformly distributed arrays with specified correlation, 2) Initial con-
 124 dition ensemble output of a regional coastal model (OSOM) (Sane et al., 2021) over July-
 125 August 2006 where most variables are not Gaussian, and 3) The GFDL-ESM2M Large
 126 Ensemble (Rodgers et al., 2015; Deser et al., 2020), an climate model initial condition
 127 ensemble hereby referred to as GFDL-LE. This large ensemble dataset contains histor-
 128 ical and future projection data following the RCP 8.5 scenario. All the GFDL-LE monthly
 129 mean data from 1950 to 2100 were used in the analysis.

130 In Part 2, we use OSOM to demonstrate the use of Shannon entropy and mutual
 131 information to quantify the extrinsic forcing effects of altered boundary forcing types.
 132 For example, is wind forcing dominant over river forcing, does using temporal averaged
 133 river runoff cause any appreciable changes in estuarine circulation, or does change in the
 134 wind product alter circulation? In coastal and estuarine systems, knowledge of which
 135 forcings are dominant helps prioritize data collection and refinement of the most impact-
 136 ful forcings.

137 1.1 Information theory

138 We introduce information theory concisely assuming the reader has no background
 139 knowledge—this section contains standard definitions. Consider a probability distribu-
 140 tion p_i obtained after binning data into N bins. The user chooses the appropriate num-
 141 ber of bins or bin widths for the range of data. Shannon (1948) identified the average
 142 information content in N possible outcomes, equally or not equally likely, as given by:

$$H = \sum_{i=1}^N p_i \log_2(1/p_i), \quad (1)$$

143 where H is the Shannon entropy with unit of bits when log is base 2 and p_i is the prob-
 144 ability of the i^{th} outcome. The factor $\log_2(1/p_i)$ measures the information of the i^{th} out-
 145 come as proposed by Hartley (1928) and is also a measure of uncertainty (Cover, 1999),
 146 as it measures the information gained by knowing that the i^{th} outcome has happened
 147 or equivalently that the variable falls in the i^{th} bin. The term information does not mean
 148 knowledge, but it means the amount of uncertainty shown by a variable or the freedom
 149 that a variable has when visiting different combinations of the N bins. Shannon (1948)
 150 found Equation 1 to provide the average information (or uncertainty) for all events in

151 a record. For the entire set of elements, a highly probable event has less uncertainty as-
 152 sociated with it, and a low probability event has high uncertainty associated with it. Thus,
 153 the prefactor p_i is used to weight the information over all possibilities. One way to in-
 154 terpret the need for the prefactor p_i is that in repeated experiments the events with higher
 155 probability will occur more often; hence they should contribute more to the quantifica-
 156 tion of variability than infrequent events.

157 Stone (2015) gives an intuitive way to understand Shannon entropy using a binary
 158 tree. A binary tree is a tree chart which starts with one node and splits into two branches
 159 at each node. At each node you can take a left or right turn to proceed, and if there are,
 160 say, 3 levels in the tree, then 8 (i.e. 2^3) outcomes or possible destinations exist. If a bi-
 161 nary tree has N equally probable outcomes, then the set of instructions required to reach
 162 the correct destination is given by $h = (N)(1/N) \log_2(N) = \log_2(N)$. The *uncertainty*
 163 about reaching the correct destination will be removed by providing $\log_2(N)$ *bits* of in-
 164 formation. In other words, if the entropy is h then 2^h states are possible.

A second metric from Shannon (1948) which is also widely used is *mutual infor-*
mation. The mutual information between two signals x and y denoted by $I(X; Y)$ is (Cover,
 1999)

$$I = \sum_{j=1}^N \sum_{i=1}^N p_{ij} \log_2 \left(\frac{p_{ij}}{p_i p_j} \right), \quad (2)$$

165 where p_{ij} is the joint probability of i^{th} outcome of x and j^{th} outcome of y . The marginal
 166 probability of i^{th} and j^{th} outcomes of x and y respectively are p_i and p_j . The addend
 167 within the summations can be expanded to $p_{ij} (\log_2(p_{ij}) - \log_2(p_i) - \log_2(p_j))$. I can
 168 be interpreted as the extra information in entropy of marginal distributions of x and y
 169 over the joint distribution. Mutual information is symmetric between x and y and is the
 170 measure of the amount of information they share. For example, if the distributions are
 171 statistically independent, then $p_{ij} = p_i p_j$ and thus $I = 0$. If the two records x and
 172 y are identical, then $p_{ij} = p_i = p_j$ and $I = H$. I is the average reduction in uncer-
 173 tainty in x due to knowing y or vice versa and denotes how much information is trans-
 174 mitted between the two variables.

175 In the context of ocean or climate modeling, entropy can be used to measure vari-
 176 ability in a model output or available data. This is in tandem with the interpretation
 177 of Shannon entropy in physical sciences as given in Carcassi et al. (2021). When calcu-
 178 lating the Shannon entropy, the primary concern is counting the possible states, e.g. the

179 various bins in a histogram, where the variable can go into while any assigned bin value
 180 or its dimensions are of lesser importance. Entropy metrics measure variability in *bits*
 181 (when the logarithm is of base 2), and hence changing the scale, e.g. switching from Cel-
 182 sium to Fahrenheit for temperature, does not change the value of variability (under equiv-
 183 alent binning). Mutual information and entropy are both dimensionally agnostic. They
 184 are also not sensitive to outliers due to the weighting prefactor and can capture nonlin-
 185 ear interactions (Watanabe, 1960; Correa & Lindstrom, 2013) and discontinuous distri-
 186 butions, including states visited intermittently. We will present the effect of correlation
 187 and outliers by examples of idealized random vectors.

188 The following methods and results sections are divided into the two parts of the
 189 overall objective of the paper. Parts A of both sections relate to evaluating intrinsic and
 190 extrinsic variability in ensemble models. Parts B describe the usage of Shannon entropy
 191 and mutual information on coastal regional modeling data to understand and compare
 192 the effects of using different boundary conditions.

193 2 Methods

194 2.1 Part A: Intrinsic and Extrinsic Variability for Ensemble Data

Analysis begins on each grid point on the ocean surface or ocean bottom. Let a vari-
 able in the ensemble be given by $f(n, t, x, y)$ where f is the variable, n denotes the in-
 dex of the ensemble member and goes from 1 to N , t is the time index and goes from
 t_1 to t_M , x, y represents the spatial grid point at the surface or bottom. The total num-
 ber of members of the ensemble is N and each member has M time steps. At a partic-
 ular grid point $f(n, t, x, y)$ is $f(n, t)$. To obtain the signal due to extrinsic forcings, the
 “differencing” approach (Frankcombe et al., 2015) has been followed to estimate the forced
 response. This approach involves averaging the members of the ensemble to derive *en-*
semble mean. The ensemble mean is given by the following:

$$g(t) = \frac{1}{N} \sum_{n=1}^{n=N} f(n, t) \quad (3)$$

$g(t)$ is a single time-varying signal for each grid point obtained by averaging across the
 ensemble members. There are potential problems with assuming that the ensemble mean
 represents extrinsic variability only, such as if models are differently sensitive to the forc-
 ing signal based on the model’s equilibrium sensitivity, as elaborated in Frankcombe et
 al. (2015) and Johnson et al. (2023). For a first-order approximation, we will assume the

ensemble mean is the best estimate of the forced response. Once $g(t)$ is obtained, the intrinsic variability can be estimated by subtracting the ensemble mean $g(t)$ from each ensemble member. The ensemble signal, forced response, and intrinsic variability are then related by:

$$f(n, t) = g(t) + \eta(n, t), \quad (4)$$

195 where $\eta(n, t)$ is the intrinsic variability or noise that differs from one ensemble member
 196 to another. Note that the above decomposition takes place at each grid point. In Fig-
 197 ure 1a, $f(n, t)$ are shown by multi-colored ensemble members. $g(t)$ is shown by a thick
 198 black line. As seen in Figure 1b, $g(t)$ has a probability distribution shown in gray and
 199 subsequently has the first, second and possibly important higher statistical moments. The
 200 gray density histogram shows variability due to extrinsic factors, and the pink density
 201 histogram shows total variability given by extrinsic and intrinsic factors.

2.1.1 Evaluating entropies

202
 203 The ensemble simulation data has been used without detrending to evaluate $g(t)$
 204 and $\eta(n, t)$. Detrending will remove some nonstationarity from the data, but will also
 205 remove some part of the extrinsic variability. Our aim is not to determine the forced re-
 206 sponse but to estimate the degree of *variability* contributed by the forced response (ex-
 207 trinsic response) and the intrinsic variability originating from the intrinsic chaos. Met-
 208 rics have been calculated at each grid point by treating them independently.

209 Usually we are limited in the number of ensemble members due to computational
 210 costs, so we concatenate into a *jugaad* in order to use *all* the ensemble members at once
 211 to evaluate information entropies. All the ensemble members given by $f(n, t)$ are rear-
 212 ranged into a single row vector f as:

$$f = [f(1, t_1), f(1, t_2), \dots, f(1, t_M), f(2, t_1), f(2, t_2), \dots, f(N-1, t_M), f(N, t_1), \dots, f(N, t_M)], \quad (5)$$

213 and g is the row vector obtained by arranging N copies of $g(t)$ in the following fashion:

$$g = \underbrace{[g(t_1), g(t_2), \dots, g(t_M)]}_1, \underbrace{[g(t_1), g(t_2), \dots, g(t_M)]}_2, \dots, \underbrace{[g(t_1), g(t_2), \dots, g(t_M)]}_N \quad (6)$$

214 This enables wide sampling and obtains an accurate probability distribution for f (as-
215 suming approximate stationarity, or enforcing stationarity by detrending), and allows
216 g to be of the same size as f and having the same probability distribution as that of $g(t)$.
217 The information statistics we get at each grid point are time-invariant, since the com-
218 plete time series is considered. It is the user's choice to choose either the complete time
219 series or a section of it for analysis. We have chosen the whole time series because this
220 is a sufficient demonstration of the value of information theory metrics. A time-evolving
221 analysis raises additional issues about causality and the shifting probabilities distribu-
222 tions of climate states that are not the focus here (X. S. Liang, 2013; DelSole & Tippett,
223 2018). By using the whole time series, we treat all variability as drawn from the same
224 distribution and seek only to associate internal (associated with each ensemble member)
225 and external (associated with the ensemble mean) sources of variability following Leroux
226 et al. (2018). The time series f and g are both expressed as row vectors of the same size,
227 $N \times M$. This step is crucial, as vectors having the same number of elements are nec-
228 essary to evaluate joint probability distribution. This enables us to calculate the mutual
229 information between f and g .

230 Calculating the Shannon entropy of f and the mutual information between f and
231 g is a difficult task that necessitates careful consideration. Optimal binning for precise
232 measurement of information entropies is a research topic in itself, and various techniques
233 have been proposed, such as equidistant partitioning, equiprobable partitioning, k near-
234 est neighbor, usage of B-spline curves for binning to name a few (Hacine-Gharbi et al.,
235 2012; Kowalski et al., 2012; Knuth, 2019). A comprehensive review of these methods can
236 be found in Papan and Kugiumtzis (2008). Although the histogram binning technique
237 is one of the most commonly used techniques (for example Campuzano et al. (2018); Potha-
238 pakula et al. (2019); Shin et al. (2023)), it introduces uncertainty. There are several tech-
239 niques to estimate this uncertainty, such as the one proposed in Knuth et al. (2005). In
240 this article, we use histograms with equidistant partitioning where constant optimal bin
241 widths are determined using the Freedman-Diaconis rule (Freedman & Diaconis, 1981;
242 Knuth, 2019) at each grid point to get a discrete probability distribution. The same bin
243 width was used for the marginal and joint probability distributions. Two approaches were
244 used to estimate the sensitivity of the metric to binning: varying the bin width around
245 the optimal value and bootstrapping over the ensemble members. The metrics were found
246 to be more sensitive to changes in the bin widths than to bootstrapping. Therefore, to

247 estimate uncertainty, if the width of the bin was found to be δw , then it was varied from
 248 $0.5\delta w$ to $1.5\delta w$ to obtain a reasonable estimate of uncertainty. Sweeping across the num-
 249 ber of bins was performed also in (Sane et al., 2021) to get an estimate of predictabil-
 250 ity time-scale.

251 **2.1.2 Information theory based metric**

Using f and g , we propose the following metric γ , which has the same intent as met-
 rics in (Leroux et al., 2018) to quantify the fraction of variability that is intrinsic, i.e.,
 the typical amount that is unique to an ensemble member or statistical instance, but un-
 like (Leroux et al., 2018) this metric is built from standard information theory quanti-
 ties:

$$\gamma = 1 - \frac{I(f;g)}{H(f)}. \quad (7)$$

252 $H(f)$ is the Shannon entropy of f , and $I(f;g)$ is mutual information between f and g .
 253 $I(f;g)$ calculates the contribution of extrinsic signal g to the whole ensemble. $H(f)$ is
 254 the total variability in the ensemble output which is the result of extrinsic and intrin-
 255 sic factors. The metric γ gives *ratio of intrinsic variability to total variability*. When $f \rightarrow$
 256 g , then $I(f;g) \rightarrow H(f) = H(g)$ from (2). This makes $\gamma = 0$ when there is no intrin-
 257 sic variability or chaos. When intrinsic chaos fully dominates the ensemble output, i.e.
 258 f and g are fully decorrelated, then $I(f;g) = 0$ yielding $\gamma = 1$. We see that γ satis-
 259 fies the extremes of zero noise and total chaos.

260 Related quantities appear in other applications. The quantity $I(f;g)/H(f)$ is de-
 261 fined as ‘‘uncertainty coefficient’’ (Eshima et al., 2020). It is the ratio of entropy of f ex-
 262 plained by g . $H(f)$ and $I(f;g)$ are related through conditional entropy by $H(f) = I(f;g) +$
 263 $H(f|g)$ (Cover, 1999). $H(f|g)$ is the conditional entropy $H(X|Y) = \sum p(x|y) \log_2 p(x|y)$
 264 (Cover, 1999). It is not necessary to calculate the conditional entropy to arrive at γ . $H(X|Y)$
 265 gives the average uncertainty about the value of f after g is known, or just the uncer-
 266 tainty in f that is not attributed to g but is attributed to η . Hence $H(f) - I(f;g)$ es-
 267 timates variability due to intrinsic chaos and γ gives the fraction of the variability due
 268 to intrinsic chaos.

269 $I(f;g)$ takes into account any correlation or information shared between f and g .
 270 This is vital because even though the spread of the model η is treated similarly to the
 271 noise added to the mean signal, it might be that the spread of the model depends on the

272 mean signal. A simple example is that if the model spread is relative (e.g., 10% of the
 273 mean signal, or *multiplicative noise*), rather than absolute (e.g., 2 units, or *additive noise*),
 274 then there is information about the model spread contained in the ensemble mean sig-
 275 nal. The nonlinear and chaotic nature of fluids often leads the mean flow to amplify the
 276 chaotic signal (e.g., eddies) and thereby result in altered variability statistics that can
 277 be represented as multiplicative noise.

278 Returning to the binary tree analogy, $I(f;g)$ would be the set of instructions sent
 279 by a source to reach one among $2^{H(f)}$ possible destinations in the presence of noise hav-
 280 ing $H(f|g)$ entropy. To capture the entropy in the noisy binary tree, to each of the $2^{I(f;g)}$
 281 micro-state possibilities, noise ($2^{H(f|g)}$) gets multiplied and the relation becomes $2^{H(f)} =$
 282 $2^{I(f;g)}2^{H(f|g)}$. Another analog of a component of the climate system is a noisy commu-
 283 nication channel as given in Leung and North (1990), where the governing equations of
 284 ocean (atmosphere) modeling are taken to communicate from forcing to response. The
 285 extrinsic forcings are inputs to the channel, the intrinsic chaos is the noise created be-
 286 cause of channel’s inherent mechanisms, while the outputs are the ensemble members.
 287 A noiseless channel will give γ as zero, and a completely noisy channel where the out-
 288 put is independent of the input will give γ as 1.

289 A seemingly enticing and simpler alternative is $\gamma = 1 - \frac{H(g)}{H(f)}$, i.e. just the differ-
 290 ence between the entropy of the ensemble and the mean entropy as a ratio with the en-
 291 tropy of the ensemble. However, this formulation is incorrect because $H(g)$ does not quan-
 292 tify the contribution of extrinsic factors to the variability in the ensemble, it only quan-
 293 tifies the variability of the mean. Relatedly, $H(f) - H(g)$ does not correctly manage the
 294 mutual information between the ensemble members and their mean in estimating intrin-
 295 sic variability.

296 Another alternative was proposed by (Gomez, 2020): using Shannon entropy di-
 297 rectly as a measure of intrinsic variability. They propose using Shannon entropy of model
 298 spread $\eta(n, t)$ at each time step normalized by the logarithm of the number of bins uti-
 299 lized. Their metric has a lower limit of 0 and an upper limit of 1, where 0 denotes zero
 300 noise and hence zero intrinsic variability and 1 denotes complete intrinsic variability. Again,
 301 this metric is similar to γ in building upon information theory, but γ takes into account
 302 the variability of the ensemble mean, the correlations between the ensemble mean and
 303 the intrinsic variability, and it is time invariant. A time-dependent version of γ can be

304 made using running time windows instead of the whole time series, but care in quanti-
 305 fying or controlling for lack of stationarity is needed in this interpretation (DelSole &
 306 Tippett, 2018). The Gomez (2020) metric uses the spread of the ensemble members sim-
 307 ilar to measuring Shannon entropy, whereas γ utilizes, in an abstract sense, the set of
 308 instructions required to choose a destination for the particular variable among the pos-
 309 sible model states.

310 **2.1.3 Variance based metric**

A variance based metric as given in (Leroux et al., 2018) has been utilized to com-
 pare with our information-based metric. The variance-based metric measures intrinsic
 and extrinsic variability using the second moment, variance. It involves calculation of
 the following terms σ_g and σ_η given by:

$$\sigma_g^2 = \frac{1}{M} \sum_{t=1}^{t=M} \left(g(t) - \overline{g(t)} \right)^2, \quad (8)$$

$$\sigma_\eta^2(t) = \frac{1}{N} \sum_{n=1}^N \eta(n, t)^2, \quad (9)$$

where the overbar denotes the temporal averaging. Total variability has been estimated
 as $\left(\sigma_g^2 + \overline{\sigma_\eta^2(t)} \right)^{1/2}$. The forced variability σ_g is equivalent to $I(f; g)$, and the total vari-
 ability $\left(\sigma_g^2 + \overline{\sigma_\eta^2(t)} \right)^{1/2}$ is equivalent to $H(f)$. Therefore, γ is compared to γ_{std} given by

$$\gamma_{std} = \frac{\left(\overline{\sigma_\eta^2(t)} \right)^{1/2}}{\left(\sigma_g^2 + \overline{\sigma_\eta^2(t)} \right)^{1/2}} \quad (10)$$

311

312 **2.2 Part B: Information Entropy and Boundary Forcing**

313 **2.2.1 Impact of changes in boundary forcings in coastal models**

314 Here instead of using the new metric γ , we use its components– Shannon entropy
 315 and mutual information–individually to compare variability between different simula-
 316 tions. Quantifying differences because of modifications in the extrinsic forcings may be
 317 required for coastal applications where systems vary predominantly due to external forc-
 318 ings. For these forcing significance experiments, OSOM was run after modifying the ex-
 319 ternal forcings (Table 1). OSOM is forced by tides, river runoff, atmospheric winds, air-
 320 sea fluxes, etc. All model details can be found in Sane et al. (2021). For this compar-

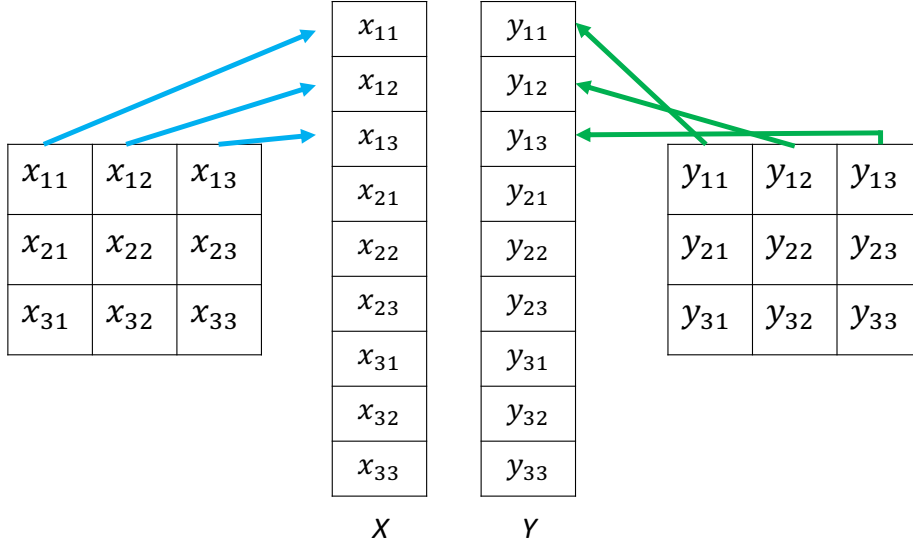


Figure 2. Flattening process for comparing two-dimensional fields using Shannon entropy and mutual information. As the flattened arrays x_1, x_2, \dots and y_1, y_2, \dots may not have linear dependence on each other, using linear dependence measures such as Pearson’s correlation might produce incorrect results. Mutual information measures nonlinear correlations and hence captures all linear and nonlinear dependence.

321 ison, we quantify the effects of altering forcing on 4 modeled variables: sea surface tem-
 322 perature and salinity, and bottom temperature and salinity. One control and four altered
 323 forcing sets were utilized,

- 324 1. (Control) Full atmospheric forcing using the North American Mesoscale (NAM)
 325 analyses, a data-assimilating, high resolution (12 km) meteorological simulation
 326 ([https://www.ncei.noaa.gov/data/north-american-mesoscale-model/access/](https://www.ncei.noaa.gov/data/north-american-mesoscale-model/access/historical/analysis)
 327 [historical/analysis](https://www.ncei.noaa.gov/data/north-american-mesoscale-model/access/historical/analysis)) denoted FF. FF stands for full forcing.
- 328 2. Full set of atmospheric forcing, but using the winds of the Northeast Coastal Ocean
 329 Forecast System (NECOFS) winds (Beardsley & Chen, 2014) instead of NAM, de-
 330 noted as NECOFS.
- 331 3. River flows are replaced with their monthly averaged flow, other forcing as in FF
- 332 4. River flows set to zero, other forcing as in FF.
- 333 5. Wind forcing set to zero, other forcing as in FF.

334 These forcing sets have been tabulated in Table 1. The aim is to quantify the effect on
 335 total variability by removing or altering one of many processes that might contribute.

Forcing Set	Wind forcing	River forcing
FF	NAM	As Observed
NECOFS	NECOFS	As Observed
MR	NAM	Time-averaged rivers
ZR	NAM	Zero river input
ZW	Zero winds	As Observed

Table 1. Different types of forcing combinations were used to test their effect on variability.

FF stands for full forcing: winds, tides, rivers, etc. For more details, see Sane et al. (2021). MR: mean rives; ZR: zero rivers; ZW: zero wind.

336 To evaluate spatial Shannon entropy, the spatial output at a particular instant in
 337 time was rearranged into a row vector by a process called flattening, as shown in Fig-
 338 ure 2. Land mask points were removed. A variable x , which is a two-dimensional vari-
 339 able, was converted to one-dimensional array (flattened) by concatenation. Shannon en-
 340 tropy was found for the flattened variable at each time step to obtain a time-varying en-
 341 tropy of each surface or bottom variable.

342 Mutual information was applied between the flattened row vectors. Our focus is
 343 on a pragmatic approach to using information theory for relative comparisons among sim-
 344 ulations, rather than an equation for the evolution of Shannon entropy and mutual in-
 345 formation with respect to time (see X. S. Liang and Kleeman (2005)). For example, if
 346 mutual information on surface salinity between FF and MR is higher than between FF
 347 and ZR, this implies that the penalty for using time-averaged river runoff is not as se-
 348 vere as using zero river runoff. The replacement of FF with MR will give more similar
 349 results to FF than replacing FF with ZR will. We can interpret this to indicate that small
 350 errors in river runoff flow rates will not cause appreciable changes to surface salinity while
 351 using zero rivers will strongly impact the solution.

3 Results

3.1 Part A: Intrinsic and Extrinsic Variability Results for Ensemble Data

3.1.1 Idealized Gaussian Arrays

We test our metric γ , equation (7) on synthetic data consisting of idealized arrays of Gaussian data: $\mathcal{N}(0, 1)$. For a normal Gaussian distribution Shannon entropy depends¹ only on the standard deviation σ . The variability in a Gaussian distribution can be increased or decreased by changing its standard deviation. Our goal is to compare γ and γ_{std} . We set out our numerical experiment as follows: we create 10 arrays, each having 10,000 elements drawn from a Gaussian distribution. Any two arrays from those 10 have a prescribed correlation coefficient between 0 and 1.

Thus, the 10 arrays are linearly correlated with a specified correlation coefficient. These 10 arrays represent ensemble members from climate simulations. The mean of 10 members gives us the synthetic forced variability signal as would be determined from the model output; averaging over the 10 ensemble members reduces the contribution from uncorrelated variability and reaffirms the covarying component into the forced variability. We apply γ and γ_{std} on this synthetic ensemble by varying the prescribed correlation coefficient from 0 to 1. Figure 3 shows that, as expected, both metrics increase as the correlation decreases, that is, as internal variability dominates forced. Both metrics behave similarly when correlation decreases, i.e., noise increases, but γ is more sensitive as correlation tends to 1. This distinction is due to the logarithmic nature of Shannon entropy for Gaussian distributions—in essence, information measured in bits is not proportional to distance measured between distributions in terms of summed variance—in the examples following the consequences of this distinction will become clearer. Critically, both functions are monotonic with correlation; however, relative comparisons (more intrinsic fraction in one region vs. a different region) are preserved.

¹ $H = \log_2 2\pi e\sigma^2$ is the Shannon entropy of a Gaussian distribution when probability density is continuous with σ as standard deviation. The Shannon entropy of a discrete probability distribution differs, which is inconsequential here, but the reader is encouraged to read Jaynes (1962). Throughout this article, discretely sampled and binned probability distributions are obtained directly from the data without any further parameterization

377 A second related experiment was derived from the first and is also shown in Fig-
 378 ure 3: adding outliers outside of the Gaussian distribution. 50 out of 10000 elements of
 379 each individual member were artificially corrupted (values were set to a constant value
 380 of 5) to test the sensitivity of both metrics. Figure 3 shows that γ is insensitive to out-
 381 liers while γ_{std} is not. γ is not sensitive because outliers occur less frequently and there-
 382 fore do not greatly affect the probability distribution, especially with the prefactor in
 383 (1) and (2). Hence, information theory metrics are robust in comparison to using stan-
 384 dard deviation (or variance). If the outliers (extreme events) occur at higher frequen-
 385 cies, information metrics will naturally start sensing them even if they are discontinu-
 386 ous from the typical conditions (e.g., multimodal distributions). The above process was
 387 repeated for 100 ensemble members, each sampled from Gaussian distributions. Increas-
 388 ing the number of ensemble members does not change the result qualitatively for both
 389 experiments. The results for a Gaussian ensemble of 10 members are shown in Figure 3
 390 a and 100 members in Figure 3 b.

391 Additionally, a set of experiments was carried out using uniformly distributed data
 392 $U(-1, 1)$. The prescribed correlated vectors were created using the procedure described
 393 in Demirtas (2014). 10 and 100 ensemble members were created and γ and γ_{std} were found
 394 between the members and their mean. The results are shown in Figure 3 c, d, respec-
 395 tively. The outlier had a value of 1.5. In all cases, γ was less sensitive to outliers than
 396 γ_{std} .

397 **3.1.2 Regional coastal model output**

398 In this section we show the results of applying γ and γ_{std} on realistic simulation
 399 data from the Ocean State Ocean Model, hereafter OSOM (Sane et al., 2021). OSOM
 400 uses the Regional Ocean Modeling System (ROMS) (Shchepetkin & McWilliams, 2005)
 401 to model Narragansett Bay and the surrounding coastal oceanic regions and waterways.
 402 OSOM's primary purpose is to understand and predictive modeling and forecasting of
 403 the estuarine state and climate of this Rhode Island body. Sane et al. (2021) gives more
 404 details about the model.

405 Using OSOM, an ensemble of simulations has been performed using perturbed ini-
 406 tial (ocean) conditions under the same atmospheric and tidal forcing for the months July
 407 and August of 2006. This ensemble consists of 10 members. Data during the first pre-

408 dictability window (20 days) where results are still linked directly to the initial condi-
 409 tions have been ignored and the remaining simulation has been used to examine vari-
 410 ability within the “climate projection” of the model beyond when forecasts or predic-
 411 tions that rely deterministically on initial conditions are possible. During this phase the
 412 different ensemble members visit different possible futures within the envelope of the pro-
 413 jected “climate” (see the related application of information theory to assess predictabil-
 414 ity in Sane et al. (2021)). The modeled temperature and salinity at each grid point typ-
 415 ically do not follow Gaussian distributions as the skewness and kurtosis each grid point
 416 shown in Figure 4 for salinity and temperature of the sea surface and bottom for the Nar-
 417 ragansett Bay region. The horizontal axis shows skewness and excess kurtosis, which are
 418 the third and fourth statistical moments, respectively, normalized by powers of the stan-
 419 dard deviation to dimensionless ratio, and in the case of excess kurtosis a constant value
 420 of 3 is subtracted. For Gaussian distributions, both skewness and excess kurtosis should
 421 be close to zero. The vertical axis denotes the number of occurrences at a grid point. Ob-
 422 serve that the majority of grid point values are away from zero and thus these variables
 423 are considerably non-Gaussian in OSOM. Therefore, the variance method in Equation (10)
 424 is at a disadvantage because the prevalence of higher statistical moments implies that
 425 the variance does not contain a complete description of the variability. The information
 426 theory metric (7) is suitable for such data as it takes into account higher moments and
 427 does not rely on Gaussian distributions.

428 Figure 7 shows the ratio of intrinsic variability to total variability applied at ev-
 429 ery point in the OSOM grid. γ_{std} is displayed on left whereas γ is shown in the center
 430 for comparison. The uncertainty in γ has been plotted in the third column in Figure 7.
 431 The features highlighted by both metrics are qualitatively different. The contribution
 432 of intrinsic chaos to total variability is more uniform using the γ metric than using γ_{std} .
 433 The intrinsic chaos displayed using γ_{std} might be misleading because the probability dis-
 434 tributions are non-Gaussian. Furthermore, where the γ metric highlights internal vari-
 435 ability, it tends to agree in similar dynamical locations—all river mouths show high sur-
 436 face salinity intrinsic variability. While surface temperature intrinsic variability is higher
 437 in more open regions of the Bay, where eddies form intermittently due to varying topog-
 438 raphy. Also note that the ranges are quite different between γ and γ_{std} , but this is to
 439 be expected from the different rate of increase with correlation seen in Figure 3.

440 **3.1.3 Earth System Model Large Ensemble**

441 A complementary experiment was performed using γ to evaluate internal versus
 442 forced variability in global climate simulation output for the RCP8.5 climate change sce-
 443 nario using the GFDL-LE model (randomly selected among the models compared). The
 444 40 members of the ensemble were utilized. The variability of sea surface temperature (Fig-
 445 ures 5) and sea surface salinity (Figures 6) were estimated using both γ and γ_{std} .

446 Note in particular the Arctic sea surface temperatures in Figure 5, which have a
 447 highly skewed and excessive kurtosis distribution due to the freezing point of seawater.
 448 The standard metric (γ_{std}) considers this region to be among the most intrinsically vari-
 449 able in the world, while the information theory metric considers it as a region of mid-
 450 dling intrinsic variability—much lower than the equatorial regions where El Nino variabil-
 451 ity is profound. This region is also subject to intermittent and drastic swings in salin-
 452 ity as sea ice forms and melts, but note that the standard metric indicates low salinity
 453 variability while the information theory metric ranks it as high in Figure 6. It is clear
 454 that a Gaussian metric should not be applied to the Arctic due to the skewness and ex-
 455 cess kurtosis, and in this case the inference is opposite using the standard and informa-
 456 tion theory metrics. In the equatorial Pacific, where Gaussian statistics are more reli-
 457 able, the two metrics agree that internal variability is high.

458 A less drastic failure occurs from the modest excess kurtosis in extratropical tem-
 459 peratures and in a few isolated regions in surface salinity. These regions are also non-
 460 Gaussian but are also not heavily skewed (i.e., they are more long-tailed and intermit-
 461 tent than Gaussian). These regions differ in the relative estimation of intrinsic versus
 462 total variability. It is also the case that the γ metric is closer to one in most regions than
 463 γ_{std} , which is expected when the correlation coefficients are low in Figure 3.

464 **3.2 Part B: Information Entropy and Boundary Forcing Results**

465 **3.2.1 Impact due to changes in boundary conditions in coastal models:**

466 We show the results of the coastal model analysis under different forcing in Fig-
 467 ures 8 and 9, under the same region as shown in Figure 7. The entropy has been plot-
 468 ted with respect to time as some variability occurs. In Figure 8, Shannon entropy is plot-
 469 ted for spatial quantities. For example, for surface salinity, all surface values have been

470 considered to find the Shannon entropy using the flattening approach. If Shannon en-
 471 tropy is more or less equal for two forcings, it implies that they similarly affect variabil-
 472 ity. Both winds and rivers seem to have similar effects in this regard. However, Figure 9
 473 displays mutual information which should be compared for two pairs of forcings. Greater
 474 mutual information implies that the two pairs share more *bits* of information, suggest-
 475 ing that one of the forcing in that pair can be replaced with the other without signif-
 476 icantly affecting variability. For temperature dependence on wind in Figure 8, we see that
 477 only NAM and NECOFS, our two realistic forcing conditions, share much mutual infor-
 478 mation. Figure 9 shows zeroing the rivers strongly reduces the salinity variability. Futher-
 479 more, in terms of salinity impact, full rivers and mean rivers share information as do NAM
 480 and NECOFS wind forcing.

481 **4 Discussion**

482 Our numerical experiments performed using γ on idealized Gaussian arrays show
 483 that γ is monotonic and decreases as the linear correlation coefficient increases. Thus,
 484 apart from the qualitative differences the new metric finds when the data are non-Gaussian,
 485 the ranges of intrinsic versus total variability are quite different between γ and γ_{std} . This
 486 is to be expected from the different rates of increase with correlation seen in Figure 3.
 487 The traditional metric (γ_{std}) falls approximately linearly as the correlation coefficient
 488 increases, so that a correlation coefficient of 0.5 gives a γ_{std} just above 0.5. The new met-
 489 ric γ agrees with γ_{std} that a correlation of 0 implies $\gamma = 1$, and a correlation of 1 im-
 490 plies $\gamma = 0$, but for a correlation of 0.5 it is closer to $\gamma = 0.9$. Only very near the cor-
 491 relation coefficients of 1 does γ fall below 0.5. If a roughly linear dependence on the cor-
 492 relation coefficient is desired, γ can be raised to a power— γ^3 resembles γ_{std} and γ^6 re-
 493 sembles the correlation coefficient. These higher powers do not lose the ability to apply
 494 to non-Gaussian data nor become non-monotonic, but they will lose their interpretation
 495 as a ratio of bits of information entropy, and instead reflect ratios of bits cubed of in-
 496 formation entropy, etc. An alternative is to take γ_{std} raised to a different power: $\gamma_{std}^{1/3}$
 497 is roughly similar to γ .

498 The uncertainty associated with binning is small—typically much less than the vari-
 499 ability across the domain and the metrics are thus not overly sensitive to the binning
 500 procedure. The exploration of alternative strategies to evaluate entropies will remain a
 501 topic of future investigation and may further improve precision.

502 As can be seen in Figures 7, 5, and 6, information theory metrics show different
 503 patterns compared to variance metrics. Information theory metrics, especially mutual
 504 information, account for *all* non-linear shared information between the ensemble mem-
 505 bers and the mean including linear correlation, and this is one reason for the differences.
 506 We have argued that non-Gaussian statistics are another (which is not wholly indepen-
 507 dent of non-linear shared relationships). There are likely other aspects of differences be-
 508 tween these metrics, but the management of these two expected aspects of geophysical
 509 fluids–nonlinear relationships and non-Gaussian distributions—justifies analyzing the
 510 data with nonparametric metrics in addition to second moment statistics.

511 For the regional coastal model OSOM, forcings differ in shared information and as
 512 to how they affect different variables. As might be expected, river runoff is more impor-
 513 tant for salinity than for temperature. However, for July to August, replacing rivers with
 514 the monthly mean river flow gives nearly the same result (in terms of variability) as fully
 515 time-varying rivers. Similarly, averaging the river runoff gives a similar effect for salin-
 516 ity compared to giving the observed river runoff in the simulations; see Figure 8. This
 517 might be due to lower river runoff during summer leading to lower variability in the flow
 518 rate hence averaging river runoff might be appropriate. We cannot conclude if there will
 519 be a similar effect in winter because the higher river runoff lead to larger variability and
 520 replacing river runoff with its mean might be unfruitful. Temperature is less sensitive
 521 to any of the forcing alterations, because although temperature and salinity are passive
 522 tracers, they have different sources and sinks. Switching the wind product from NAM
 523 to NECOFS does not have a significant effect on the sources or sinks of temperature or
 524 salinity, but switching the wind off definitely affects the parameters by eliminating wind-
 525 driven mixing altogether. Figure 9 shows that zero-wind (ZW) simulations are markedly
 526 different from the rest in terms of *mutual information* (i.e. they do not covary), although
 527 very similar in terms of amount of spatial variability (Shannon entropy, Figure 8), be-
 528 cause even without winds tides, fluxes, and rivers still vary. The zero-river case tends
 529 to eliminate both variability and mutual information (ZR).

530 If we were to prioritize improvements based on Shannon entropy and mutual in-
 531 formation, note that the two highest mutual information cases are where NAM is sub-
 532 stituted with NECOFS and where mean rivers are substituted for varying rivers. The
 533 first observation is important from a forecast perspective, because it means that we can-
 534 not easily tell the difference between different wind products, although something rather

535 than zero winds should be used if the estuary needs to be forecast for the full 20 day pre-
 536 dictability range (weather forecasts are reliable for only about 7 days in this region). Sim-
 537 ilarly, knowing that substituting the mean of the rivers for fully varying rivers has lit-
 538 tle impact implies that rivers can be fixed in time for forecasts beyond where they might
 539 be predicted based on expected weather and precipitation. Finally, despite the fact that
 540 Narragansett Bay is a dominantly tidally mixed estuary, among the sources of overall
 541 variability (i.e. sources of information entropy) considered, preserving an inflow of fresh
 542 water is key, even though that inflow can be steady. Winds do not appreciably increase
 543 information entropy of the Bay, but they are an important source of forced co-variation,
 544 and so are important for predictions but do not raise the overall level of variability.

545 **5 Conclusion**

546 We showed usage of information theory metrics to determine contribution of in-
 547 trinsic chaos and external variability to total variability in ensemble model simulations.
 548 The metric consists of Shannon entropy and mutual information and is non-parametric
 549 compared to variance. We have applied metrics on idealized Gaussian arrays, as well as
 550 realistic coastal ocean and global climate models. We conclude that:

- 551 1. The information theory metric is more reliable when outliers are present, because
 552 outliers get assigned less probability and because Gaussian distributions have a
 553 difficult time approximating long-tailed (i.e., outlier-prone) distributions.
- 554 2. The information theory metric is more reliable when variability is non-Gaussian
 555 because it is based on nonparametric measures of the probability distributions and
 556 captures nonlinear correlations.
- 557 3. The new information theory metric varies monotonically with ensemble member
 558 to ensemble mean correlation, but is quantified in fractions of bits required to cap-
 559 ture internal variability versus bits required to capture total variability.
- 560 4. The use of the information theory ratio metric in a coastal ocean modelensem-
 561 ble and a climate model ensemble qualitatively changes the focus to regions that
 562 were previously erroneously labeled as having high or low internal variability.
- 563 5. The use of Shannon entropy and mutual information can quickly focus attention
 564 on which forcing choices cause the most variability and need attention as their sub-
 565 stitutions significantly affect the outcomes. These conclusions can be drawn re-

566 regardless of the fact that the dimensions of wind, rivers, salinity, and temperature
567 have no specified unit conversion factors.

568 6. In these ensemble simulations, the coastal ensemble had a much smaller intrinsic
569 (chaotic) proportion of its total variability in comparison to the climate ensemble
570 which had more intrinsic variability (weather, climate oscillations, etc.) as a
571 proportion of its total. Importantly, the resolution of the models helps determine
572 the proportion of intrinsic variability, so such comparisons are model-specific: a
573 higher-resolution coastal model might well have a larger intrinsic fraction than a
574 coarser climate model.

575 Other applications of these and similar information theory metrics are likely to be re-
576 vealing of new behavior and sensitivity of models.

577 **Appendix A Open Research**

578 We have made the code and data available at [https://doi.org/10.5281/zenodo](https://doi.org/10.5281/zenodo.7992844)
579 [.7992844](https://doi.org/10.5281/zenodo.7992844)

580 **Acknowledgments**

581 The Rhode Island Coastal Ecology Assessment Innovation & Modeling grant (NSF
582 1655221) supported this work. B. Fox-Kemper was also supported by ONR N00014-17-
583 1-2963. This material is based upon work conducted at a Rhode Island NSF EPSCoR
584 research facility Center for Computation and Visualization (Brown University), supported
585 in part by the National Science Foundation EPSCoR Cooperative Agreement OIA 1655221.

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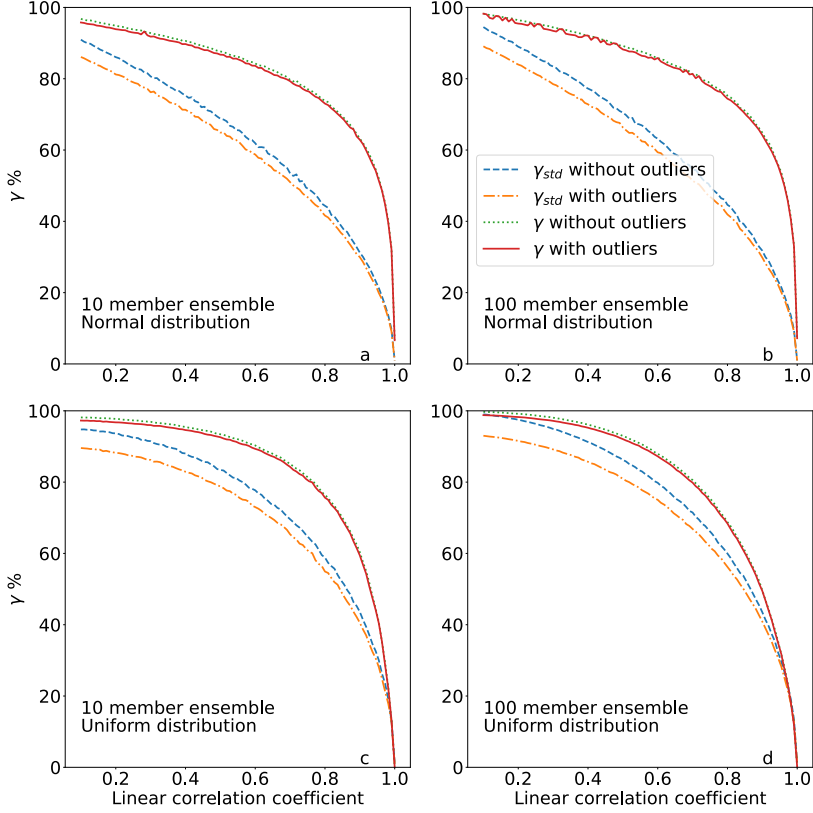


Figure 3. Information theory metric of intrinsic vs. extrinsic variability γ as a function of the correlation coefficient in idealized Gaussian correlated arrays (a and b) and idealized uniformly distributed arrays (c and d). The horizontal axis is the correlation coefficient between the mean member and ensemble members. The vertical axis shows the information theory metric γ from (7) and the traditional metric γ_{std} from Equation (10). A second related experiment is also shown adding (50 out of 10,000) “corrupted” outliers to each individual member. The information theory metric γ does not change for these outliers, which shows its robustness, while γ_{std} is highly sensitive. The results are similar for Gaussian distribution members and uniformly distributed members. γ is more sensitive around linear correlation of 1. This is due to the logarithmic nature of γ .

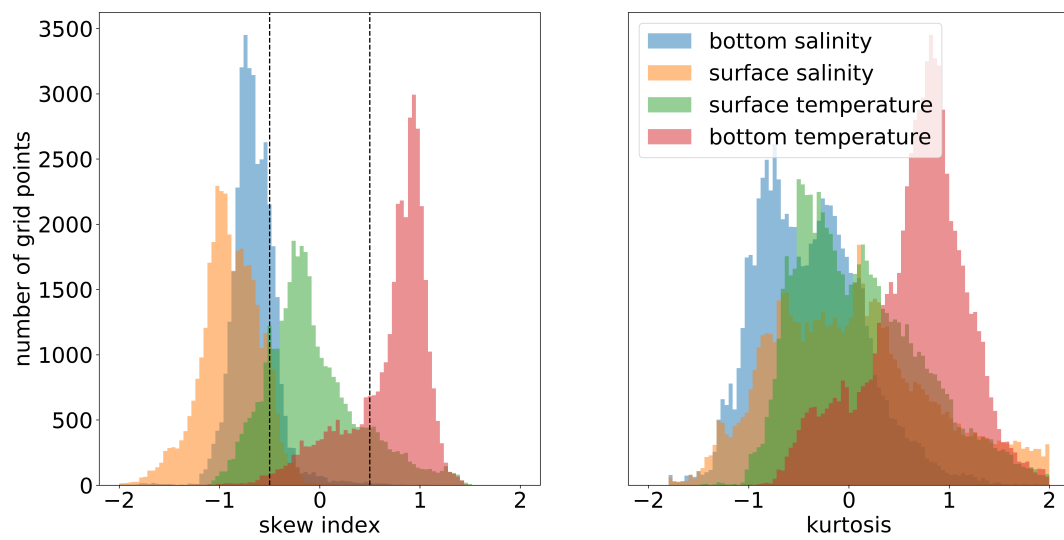


Figure 4. Grid point-wise skewness and excess kurtosis for OSOM output. Neither are close to zero, e.g., within $(-0.5, 0.5)$, suggesting that the temperature and salinity data distribution is non-Gaussian.

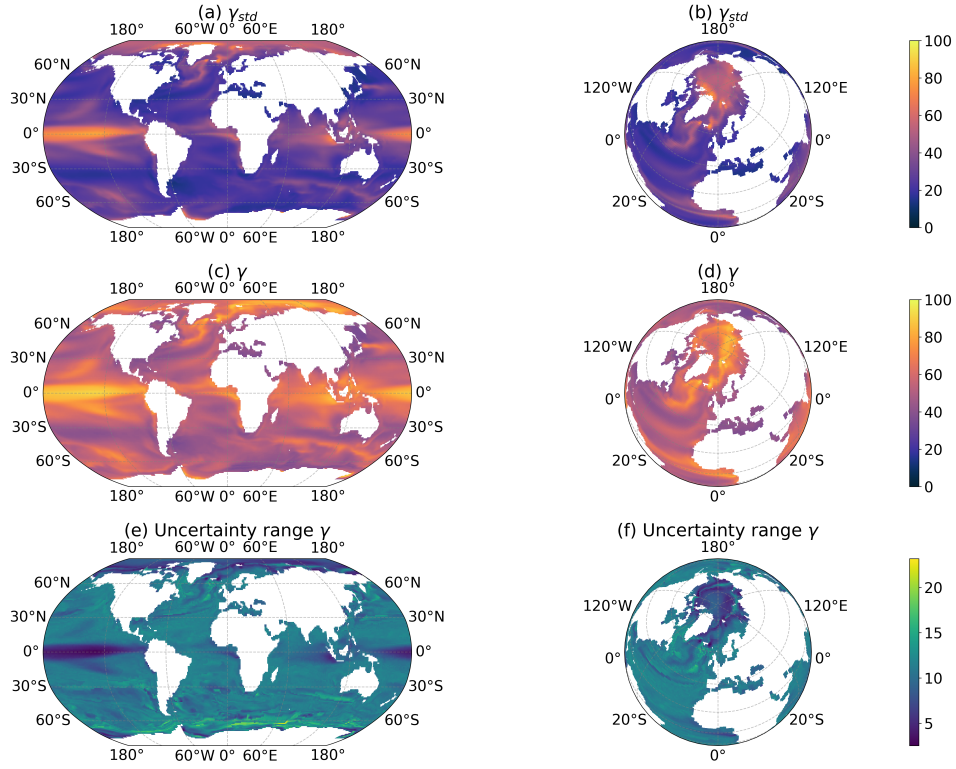


Figure 5. Intrinsic to total variability for sea surface temperature using (a, b) γ_{std} and (c, d) γ . (e, f) Uncertainty range in γ found by sweeping across the bin width as explained in the text. We can see a difference in the magnitude and pattern of the intrinsic to total variability around the Arctic region. Difference in other regions such as Mediterranean sea and Pacific equator is also visible.

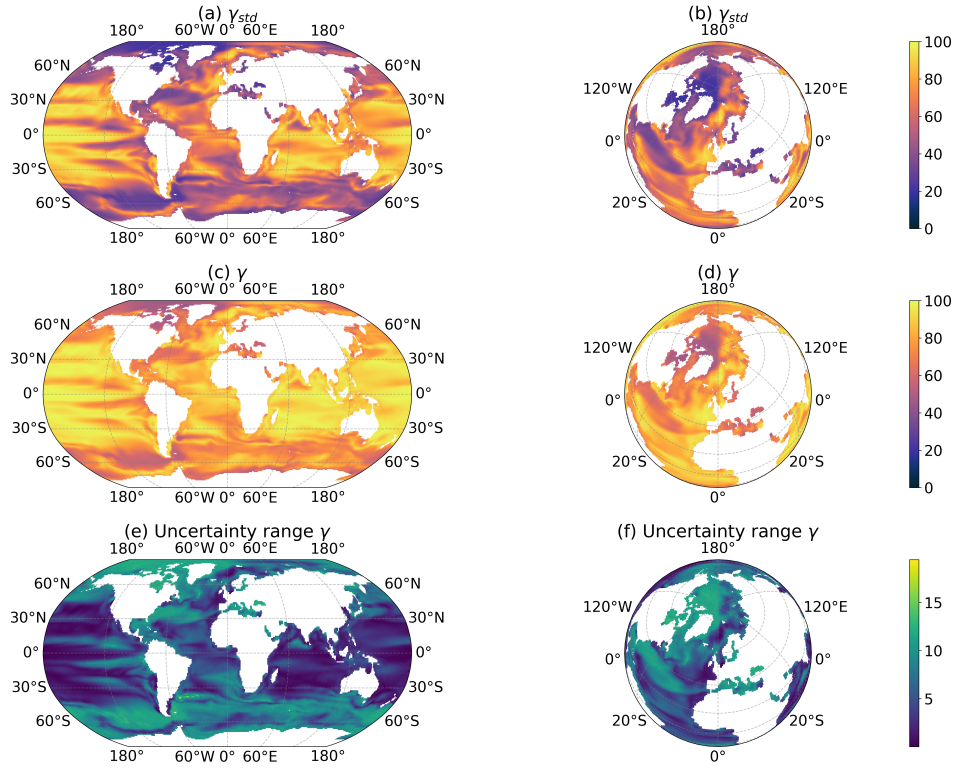


Figure 6. Intrinsic to total variability for sea surface salinity using (a, b) γ_{std} and (c, d) γ . (e, f) Uncertainty range in γ by sweeping across the bin width as explained in the text. We can see a difference in the magnitude and pattern of the intrinsic to total variability around the Arctic region. Difference in other regions such as Mediterranean sea and Pacific equator is also visible

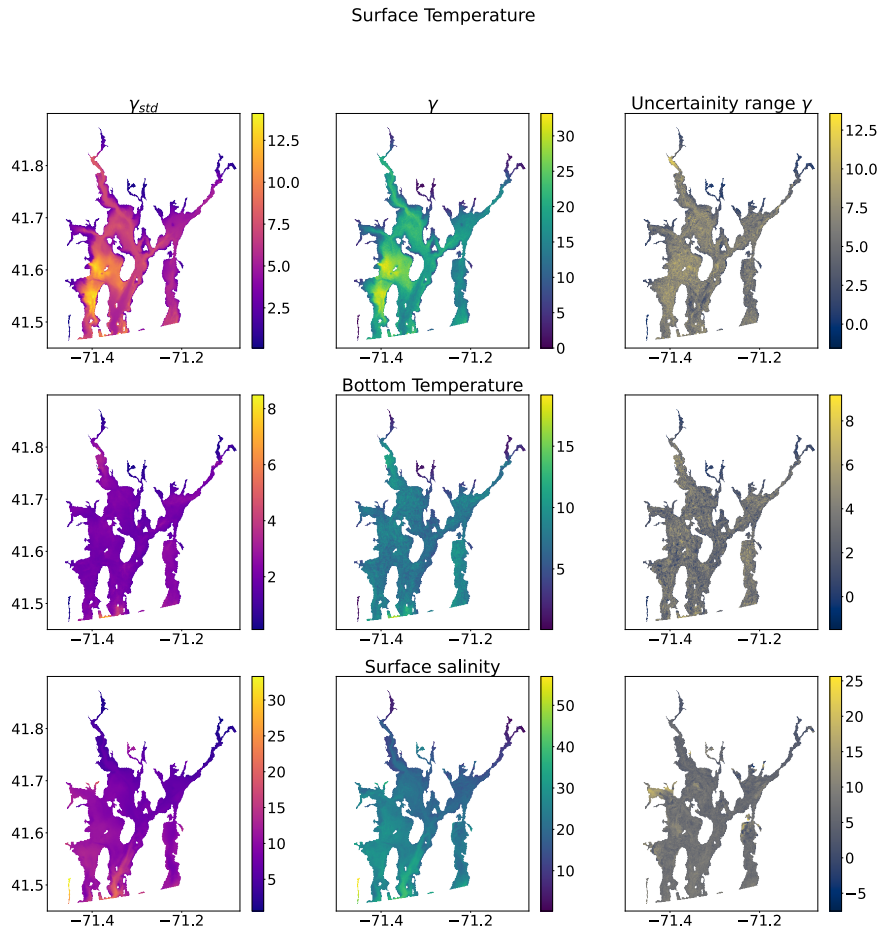


Figure 7. Metrics γ vs γ_{std} for the OSOM output. Both metrics show different contributions of intrinsic variability to total variability. γ is more uniform in the domain than γ_{std} . Right panels show the uncertainty in γ due to binning choices. The color maps for γ and γ_{std} are different to highlight their different ranges. γ_{std} for bottom temperature (not shown) has a maximum value of 5%, and the pattern is almost uniform except at the river sources where the values are on the lower side (less than 1%).

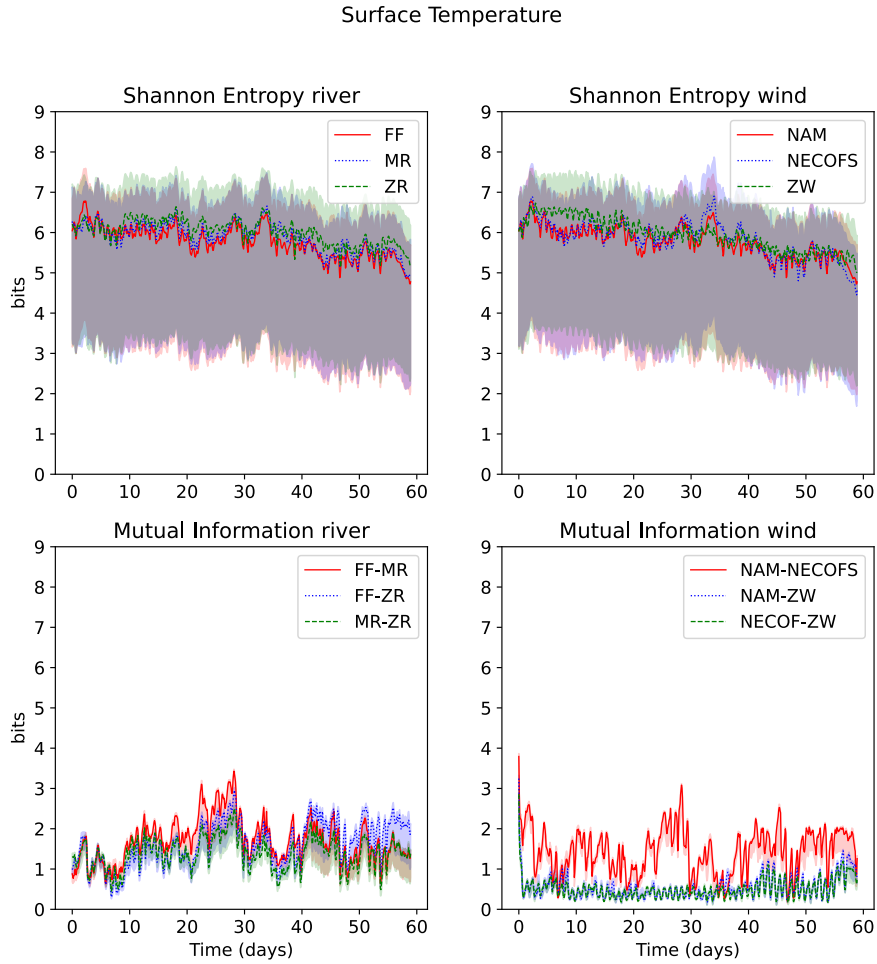


Figure 8. Shannon entropy applied to temperature and salinity. Replacing fully time-varying rivers with monthly mean river flow gives almost the same result for salinity. The same is true by replacing the wind product with a different one. Setting the river to zero affects salinity, but not temperature. Winds are important in terms of variability, but different wind products do not noticeably alter variability.

Surface Salinity

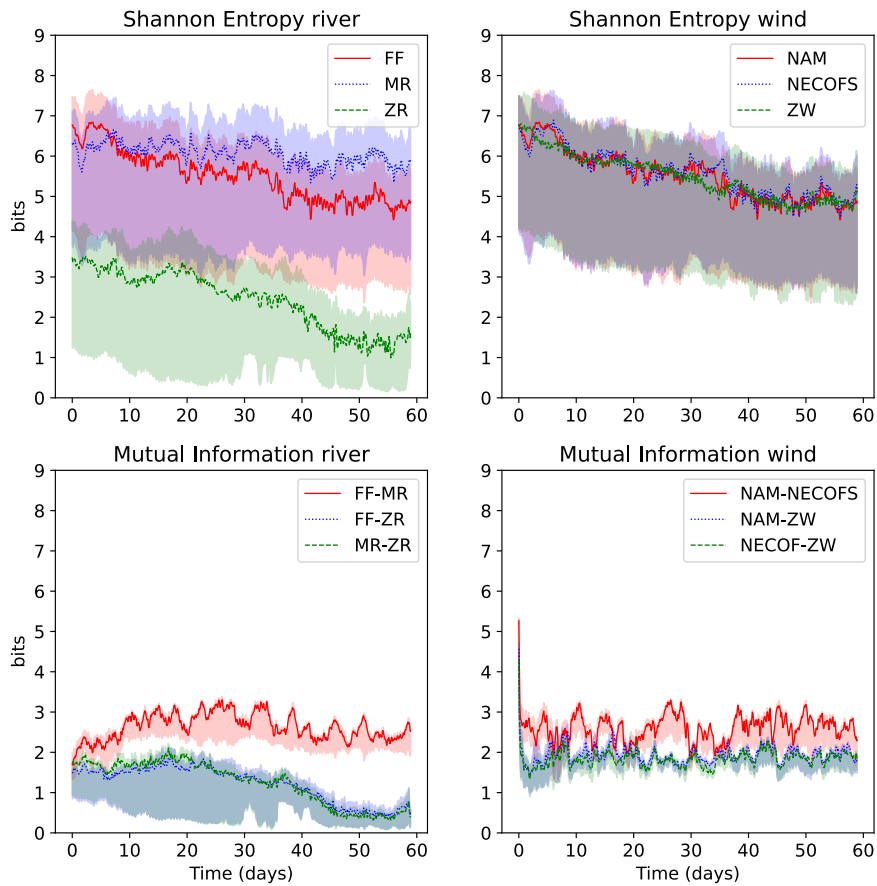


Figure 9. Mutual information applied to simulations from different forcings. Higher mutual information implies higher similarity in terms of variability. For example, NAM-NECOFS values are higher than NAM-ZW implying that NAM and NECOFS are significantly different than having no wind.

Supporting Information for "Insert Title"

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Contents of this file

1. Text S1 to Sx
2. Figures S1 to Sx
3. Tables S1 to Sx

Additional Supporting Information (Files uploaded separately)

1. Captions for Datasets S1 to Sx
2. Captions for large Tables S1 to Sx (if larger than 1 page, upload as separate excel file)
3. Captions for Movies S1 to Sx
4. Captions for Audio S1 to Sx

Introduction

Text S1.

Data Set S1.

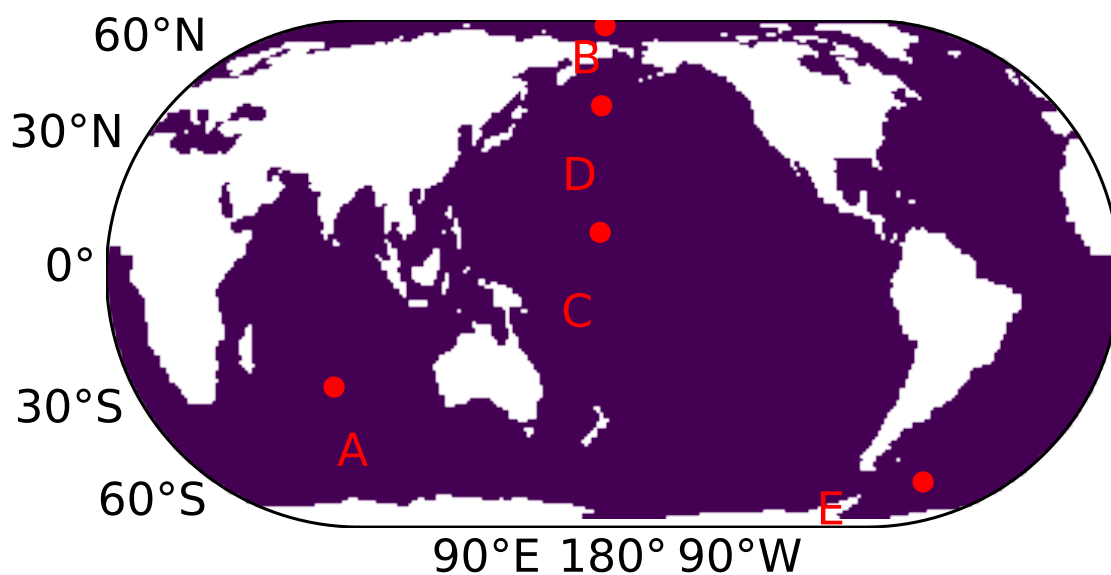
Movie S1.

Audio S1.

Table S1. Location of points for comparing uncertainty due to binning and bootstrapping

Point	Longitude	Latitude
A	75.5 ° E	29.5 ° S
B	175.5 E	80.5 ° N
C	175.5 E	10.5 ° N
D	175.5 E	45.5 ° N
E	39.5 W	59.5 ° S

^a Footnote text here.

**Figure S1.** Points where uncertainty due to binning and bootstrapping is compared.

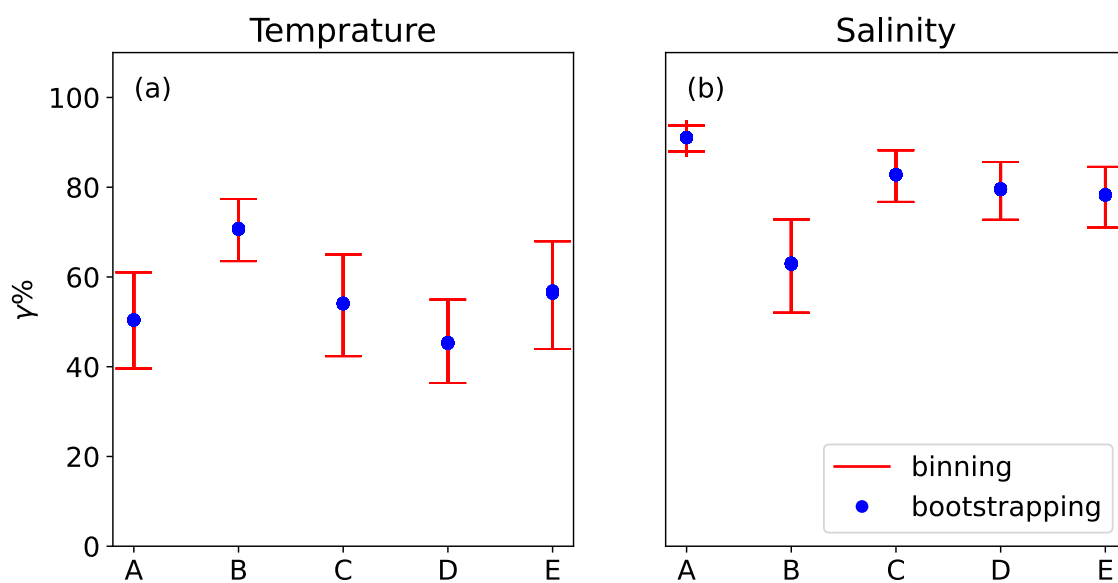


Figure S2. Uncertainty ranges due to binning and bootstrapping. Binning (red) has more uncertainty than bootstrapping (blue).

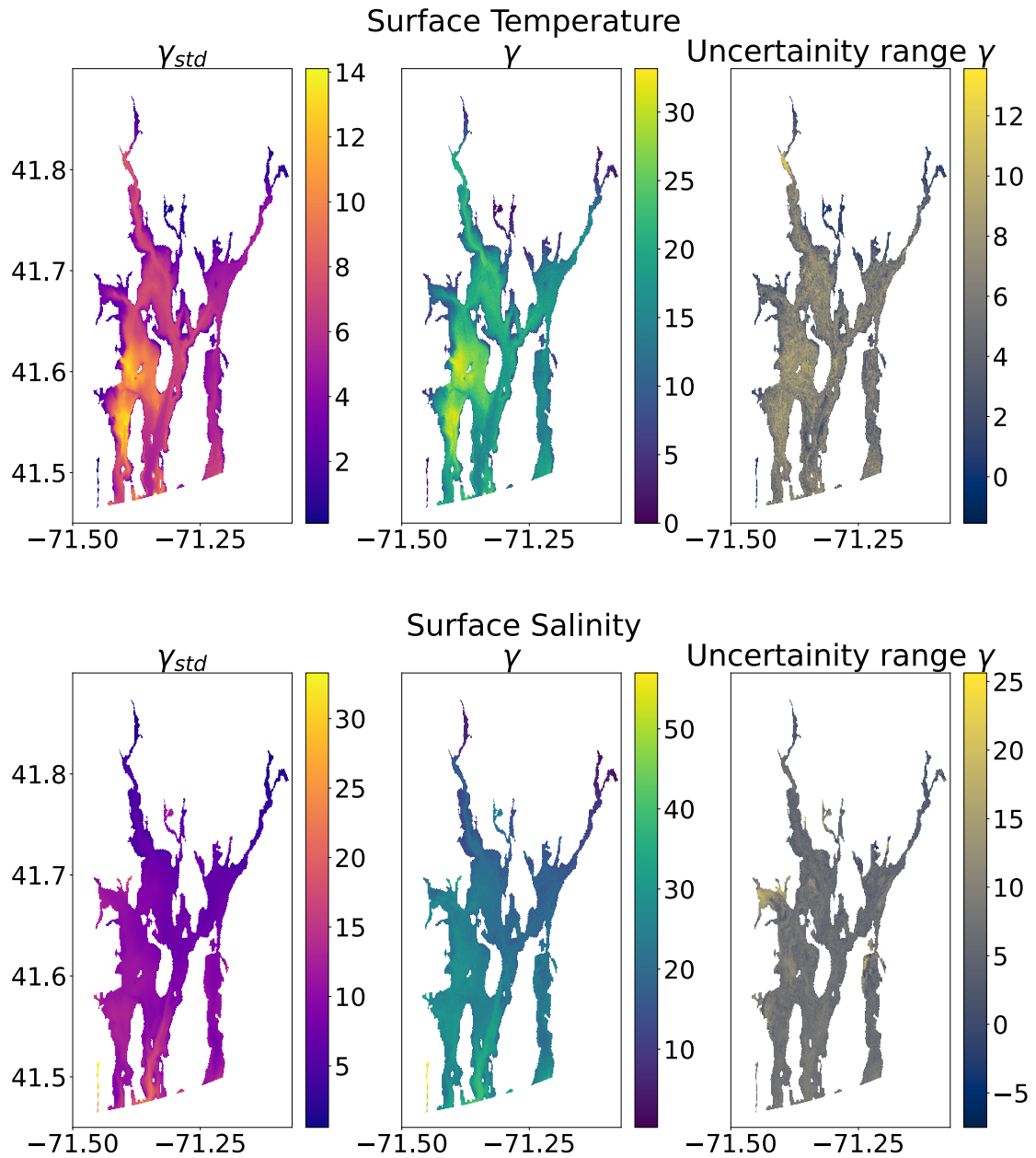


Figure S3. Metrics γ vs γ_{std} for OSOM output. Both metrics show different contribution of intrinsic variability to total variability. γ is more uniform throughout the domain than γ_{std} . Colormaps for γ and γ_{std} are different to highlight the different ranges each of them have. Uncertainty range is calculated using sweeping the bin width from 50% to 150% of the bin width as estimated using F-D rule.

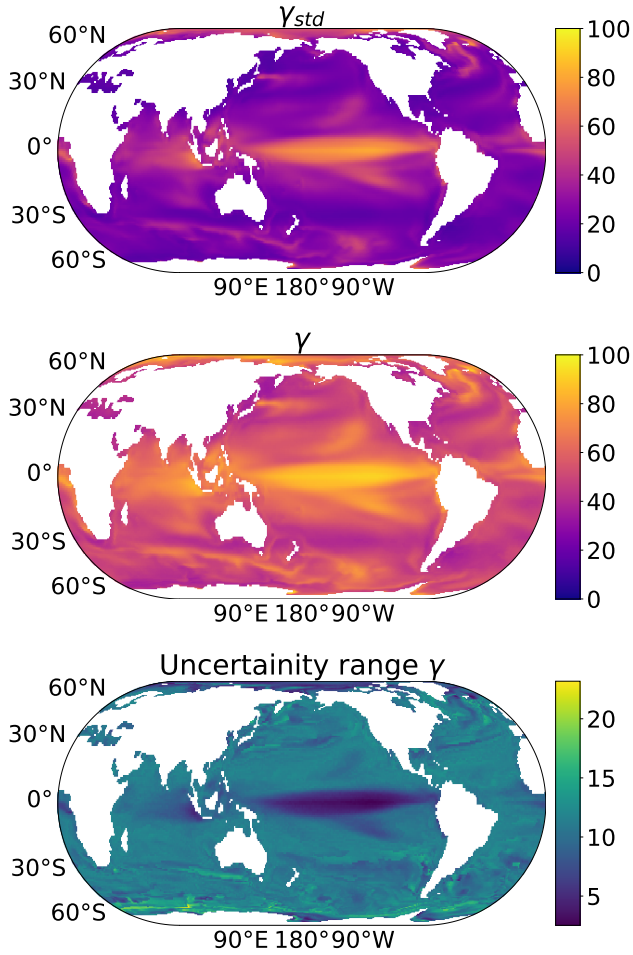


Figure S4. Top: Intrinsic to total variability percentage for sea surface temperature. Uncertainty range is calculated using sweeping the bin width from 50% to 150% of the bin width as estimated using F-D rule.

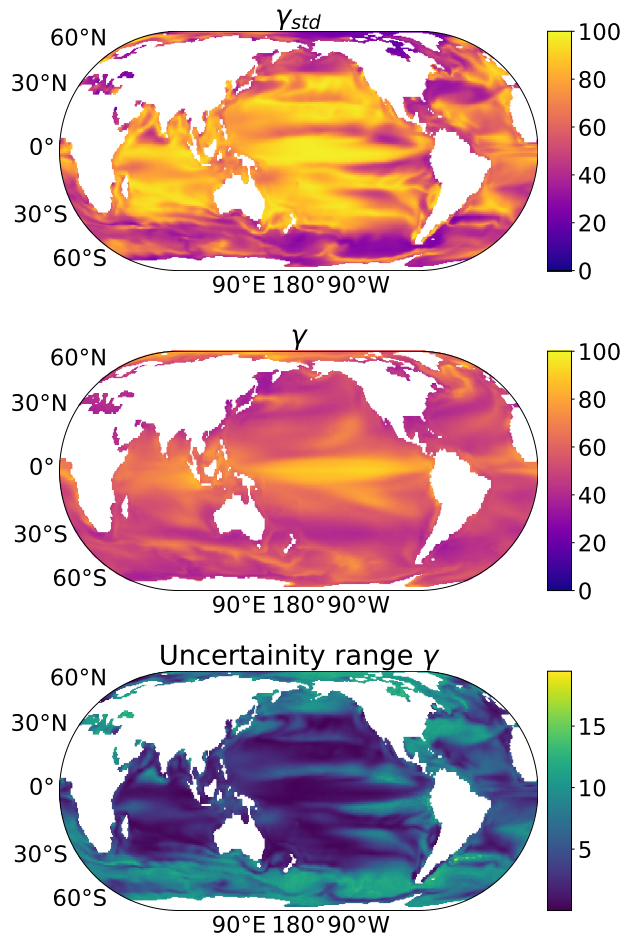


Figure S5. Top: Intrinsic to total variability percentage for sea surface salinity. Uncertainty range is calculated using sweeping the bin width from 50% to 150% of the bin width as estimated using F-D rule.

Surface Temperature

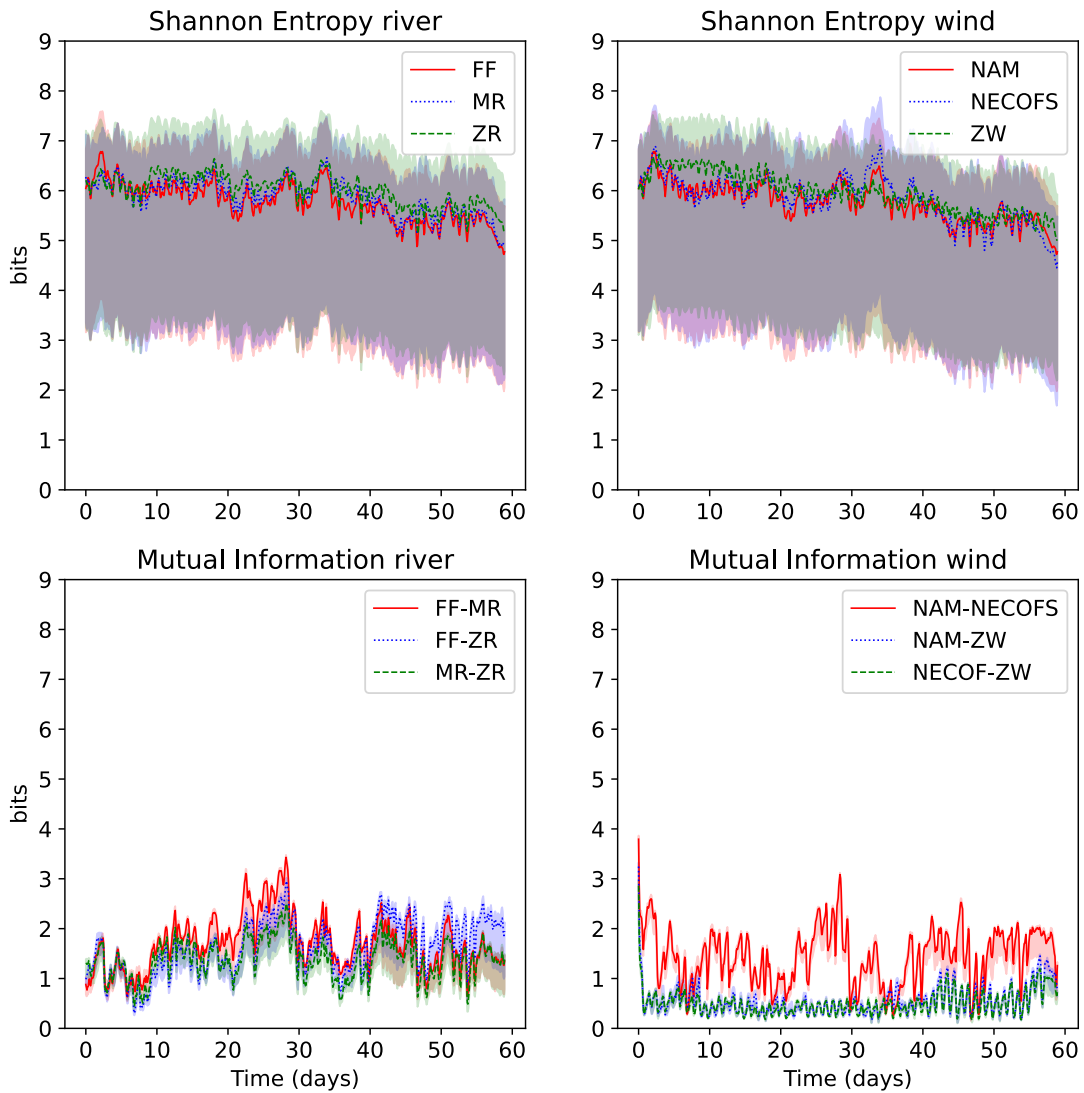


Figure S6. Shannon entropy applied to temperature and salinity. Replacing fully time varying rivers with monthly-mean river flow gives almost the same result for salinity. Same is true by replacing wind product with a different one. Rivers set to zero affects salinity but not temperature. Winds are important in terms of variability but different wind products do not noticeably alter variability. Uncertainty range is calculated using sweeping the bins from 50 to 800. The median is assumed to be the best estimate.

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Surface Salinity

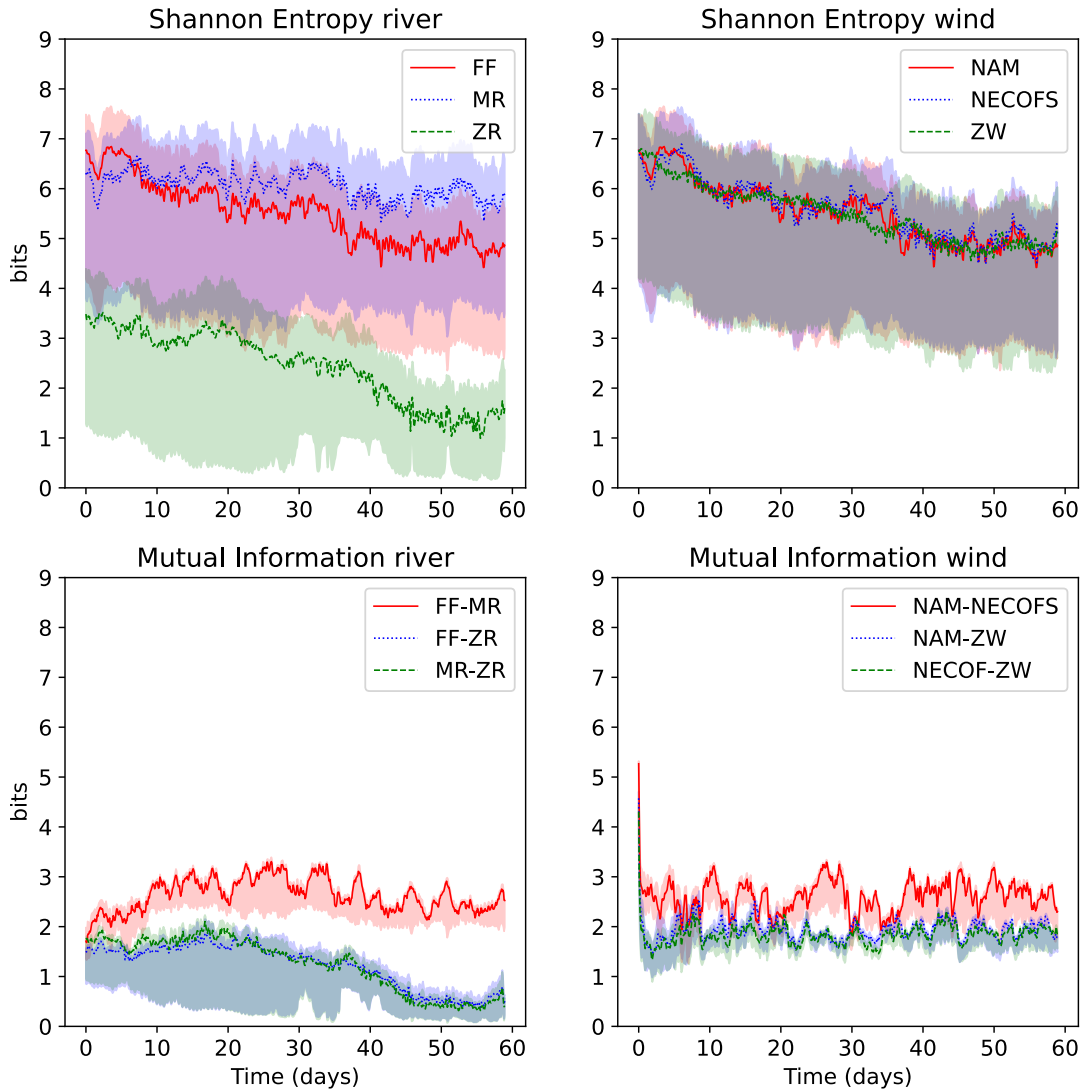


Figure S7. Mutual information applied to simulations from different forcings. Higher mutual information implies higher similarity in terms of variability. For example NAM-NECOFS values are higher than NAM-ZW implying that NAM and NECOFS are significantly different than having no wind. Uncertainty range is calculated using sweeping the bins from 50 to 800. The median is assumed to be the best estimate.